

QUATERNIONIC BERTRAND CURVES IN EUCLIDEAN 4-SPACE

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ABSTRACT. In this paper, by using the similar idea of *Matsuda and Yorozu* [12], we prove that if bitorsion of a quaternionic curve α is not vanish, then there is no quaternionic curve in E^4 is a Bertrand curve. Then we define $(1, 3)$ type Bertrand curves for quaternionic curve in Euclidean 4-space. We give some characterizations for a $(1, 3)$ type quaternionic Bertrand curves in Euclidean 4-space by means of the curvature functions of the curve.

1. INTRODUCTION

Characterization of a regular curve is one of the important and interesting problems in the theory of curves in Euclidean space. There are two ways widely used to solve these problems: to figure out the relationship between the Frenet vectors of the curves (see [11]), and to determine the shape and size of a regular curve by using its curvatures. k_1 (or κ) and k_2 (or τ), the curvature functions of a regular curve, have an effective role. For example: if $k_1 = \text{constant} \neq 0$ and $k_2 = 0$, the curve is a circle with radius $(1/k_1)$, etc.

In 1845, *Saint Venant* (see [14]) proposed the question whether the principal normal of a curve is the principal normal of another's on the surface generated by the principal normal of the given one. Bertrand answered this question in ([3]) published in 1850. He proved that a necessary and sufficient condition for the existence of such a second curve is required in fact a linear relationship calculated with constant coefficients should exist between the first and second curvatures of the given original curve. In other words, if we denote first and second curvatures of a given curve by k_1 and k_2 respectively, we have $\lambda k_1 + \mu k_2 = 1$, $\lambda, \mu \in \mathbb{R}$. Since 1850, after the paper of Bertrand, the pairs of curves like this have been called *Conjugate Bertrand Curves*, or more commonly *Bertrand Curves* (see [11]).

There are many important papers on Bertrand curves in Euclidean space (see: [4],[5],[13]).

When we investigate the properties of Bertrand curves in Euclidean n -space, it is easy to see that either k_2 or k_3 is zero which means that Bertrand curves in

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\mathbb{E}^n ($n > 3$) are degenerate curves (see [13]). This result is restated by *Matsuda and Yorozu* [12]. They proved that *there was not any special Bertrand curves in \mathbb{E}^n ($n > 3$)* and defined a new kind, which is called (1, 3)-type Bertrand curves in 4-dimensional Euclidean space. Bertrand curves and their characterizations were studied by many researchers in Minkowski 3- space and Minkowski space-time (see [1, 7, 10]) as well as in Euclidean space.

K. Bharathi and M. Nagaraj (in [2]) studied a quaternionic curve in Euclidean 3-space \mathbb{E}^3 and Euclidean 4-space \mathbb{E}^4 and gave the Frenet formula for quaternionic curve. For the newest results for quaternionic curves, we refer the papers [8, 9, 16].

In this paper, by using the similar idea of *Matsuda and Yorozu* [12], we prove that if bitorsion of a quaternionic curve α is not vanish, then there is no quaternionic curves in \mathbb{E}^4 is a Bertrand curve. Then we define (1, 3) – type Bertrand curves for quaternionic curves in Euclidean 4-space. We give some characterizations for a (1, 3) – type quaternionic Bertrand curves in Euclidean 4-space by means of the curvature functions of the curve

2. PRELIMINARIES

Let Q_H denote a four dimensional vector space over a field H whose characteristic greater than 2. Let e_i ($1 \leq i \leq 4$) denote a basis for the vector space. Let the rule of multiplication on Q_H be defined on e_i ($1 \leq i \leq 4$) and extended to the whole of the vector space by distributivity as follows:

A real quaternion is defined by $q = ae_1 + be_2 + ce_3 + de_4$ where a, b, c, d are ordinary numbers. Such that

$$\begin{aligned} e_4 &= 1 & e_1^2 &= e_2^2 = e_3^2 = -1 \\ e_1e_2 &= -e_2e_1 = e_3, & e_2e_3 &= -e_3e_2 = e_1, \\ e_3e_1 &= -e_1e_3 = e_2 \end{aligned}$$

If we denote $S_q = d$ and $V_q = ae_1 + be_2 + ce_3$, we can rewrite a real quaternion whose basic algebraic form is $q = S_q + V_q$ where S_q is scalar part and V_q is vectorial part of q . Using these basic products we can now expand the product of two quaternions as

$$p \times q = S_p S_q - \langle V_p, V_q \rangle + S_p V_q + S_q V_p + V_p \wedge V_q \text{ for every } p, q \in Q_H$$

where we have used the inner and cross products in Euclidean space E^3 . There is a unique involutory antiautomorphism of the quaternion algebra, denoted by the symbol γ and defined as follows:

$$\gamma q = -ae_1 - be_2 - ce_3 + de_4 \text{ for every } ae_1 + be_2 + ce_3 + de_4 \in Q_H$$

which is called the "Hamilton conjugation". This defines the symmetric, real valued, non-degenerate, bilinear form h as follows:

$$h(p, q) = \frac{1}{2} [p \times \gamma q + q \times \gamma p] \text{ for every } p, q \in Q_H.$$

And then, the norm of any q real quaternion denoted

$$\|q\|^2 = h(q, q) = q \times \gamma q.$$

q is called a spatial quaternion whenever $q + \gamma q = 0$ ([6]).

The Serret-Frenet formulae for quaternionic curves in E^3 and E^4 are as follows ([2]):

Theorem 2.1. *The three-dimensional Euclidean space E^3 is identified with the space of spatial quaternions $\{p \in Q_H \mid p + \gamma p = 0\}$ in an obvious manner. Let $I = [0, 1]$ denote the unit interval in the real line \mathbb{R} . Let*

$$\begin{aligned} \alpha : I \subset \mathbb{R} &\longrightarrow Q_H \\ s &\longrightarrow \alpha(s) = \sum_{i=1}^3 \alpha_i(s) e_i \quad (1 \leq i \leq 3) \end{aligned}$$

be an arc-lengthed curve with nonzero curvatures $\{k, r\}$ and $\{t(s), n(s), b(s)\}$ denote the Frenet frame of the curve α . Then Frenet formulas are given by

$$\begin{bmatrix} t' \\ n' \\ b' \end{bmatrix} = \begin{bmatrix} 0 & k & 0 \\ -k & 0 & r \\ 0 & -r & 0 \end{bmatrix} \begin{bmatrix} t \\ n \\ b \end{bmatrix}$$

where k is principal curvature, r is torsion of α .

Theorem 2.2. *The four-dimensional Euclidean space E^4 are identified with the space of unique quaternions. Let $I = [0, 1]$ denote the unit interval in the real line \mathbb{R} and*

$$\begin{aligned} \alpha : I \subset \mathbb{R} &\longrightarrow Q_H \\ s &\longrightarrow \alpha(s) = \sum_{i=1}^4 \alpha_i(s) e_i, \end{aligned}$$

be a smooth curve in E^4 with nonzero curvatures $\{K, k, r - K\}$ and $\{T(s), N(s), B_1(s), B_2(s)\}$ denotes the Frenet frame of the curve α . Then the frenet formulas are given by

$$\begin{bmatrix} T' \\ N' \\ B_1' \\ B_2' \end{bmatrix} = \begin{bmatrix} 0 & K & 0 & 0 \\ -K & 0 & k & 0 \\ 0 & -k & 0 & (r - K) \\ 0 & 0 & -(r - K) & 0 \end{bmatrix} \begin{bmatrix} T \\ N \\ B_1 \\ B_2 \end{bmatrix}$$

where K is the principal curvature, k is torsion of β and $(r - K)$ is bitorsion of α .

Definition 2.3. Let $\alpha(s)$ and $\beta(s^*)$ be two quaternionic curves in E^4 . $\{T(s), T(s), N(s), B_1(s), B_2(s)\}$ and $\{T^*(s^*), N^*(s^*), B_1^*(s^*), B_2^*(s^*)\}$ are Frenet frames, respectively, on these curves. α and β are called Bertrand curves if there exist a bijection

$$\begin{aligned} \varphi : I &\longrightarrow I^* \\ s &\longrightarrow \varphi(s) = s^*, \quad \frac{ds^*}{ds} \neq 0 \end{aligned}$$

and the principal normal lines of α and β at corresponding points coincide.

3. MAIN RESULTS

Theorem 3.1. *Let α be a quaternionic curve in E^4 . If bitorsion of α is no vanish, then there is no quaternionic curve in E^4 is a Bertrand curve.*

Proof. Let α be a Bertrand curve in E^4 and β a Bertrand mate of β . β distinct from β . Let the pair of $\alpha(s)$ and $\beta(s^*) = \beta(\varphi(s))$ be of corresponding points of α and β . Then the curve β is given by

$$\beta(s^*) = \beta(\varphi(s)) = \alpha(s) + \lambda(s)N(s) \quad (1)$$

where λ is a C^∞ function on I . Differentiating (1) with respect to s , we obtain

$$\varphi'(s) \frac{d\beta(s^*)}{ds^*} = \alpha'(s) + \lambda'(s) N(s) + \lambda(s) N'(s).$$

Here and hereafter, the prime denotes the derivative with respect to s . By the Frenet equations, it holds that

$$\varphi'(s) T^*(\varphi(s)) = (1 - \lambda(s) K(s)) T(s) + \lambda'(s) N(s) + \lambda(s) k(s) B_1(s).$$

Since $\langle T^*(\varphi(s)), N^*(\varphi(s)) \rangle = 0$ and $N^*(\varphi(s)) = \mp N(s)$, we obtain, for all $s \in I$,

$$\lambda'(s) = 0,$$

that is, λ is a constant function on I . Thus (1) are rewritten as

$$\beta(s^*) = \beta(\varphi(s)) = \alpha(s) + \lambda N(s) \quad (2)$$

and we obtain

$$\varphi'(s) T^*(\varphi(s)) = (1 - \lambda K(s)) T(s) + \lambda k(s) B_1(s) \quad (3)$$

for all $s \in I$. By (3), we can set

$$T^*(\varphi(s)) = (\cos \theta(s)) T(s) + (\sin \theta(s)) B_1(s), \quad (4)$$

where θ is a C^∞ -function on I and

$$\cos \theta(s) = \frac{1 - \lambda K(s)}{\varphi'(s)} \quad (5)$$

$$\sin \theta(s) = \frac{\lambda k(s)}{\varphi'(s)}. \quad (6)$$

Differentiating (4) and using the Frenet equations, we obtain

$$\begin{aligned} \bar{K}(\varphi(s)) \varphi'(s) N^*(\varphi(s)) &= \frac{d \cos \theta(s)}{ds} T(s) \\ &+ (K(s) \cos \theta(s) - k(s) \sin \theta(s)) N(s) \\ &+ \frac{d \sin \theta(s)}{ds} B_1(s) \\ &+ (r - K)(s) \sin \theta(s) B_2(s). \end{aligned}$$

Since $N^*(\varphi(s)) = \mp N(s)$ for all $s \in I$, we obtain

$$(r - K)(s) \sin \theta(s) = 0. \quad (7)$$

By $(r - K)(s) \neq 0$ and (7), we obtain that $\sin \theta(s) = 0$. Thus, by $k(s) \neq 0$ and (6), we obtain that $\lambda = 0$. This completes the proof of theorem. \square

Definition 3.2. Let $\alpha(s)$ and $\beta(s^*)$ be two quaternionic curves in E^4 . $\{T(s), N(s), B_1(s), B_2(s)\}$ and $\{T^*(s^*), N^*(s^*), B_1^*(s^*), B_2^*(s^*)\}$ are Frenet frames, respectively, on these curves. α and β are called quaternionic (1,3)-Bertrand curves if there exist a bijection

$$\begin{aligned} \varphi: I &\longrightarrow I^* \\ s &\longrightarrow \varphi(s) = s^*, \frac{ds^*}{ds} \neq 0 \end{aligned}$$

and the plane spanned by $N(s), B_2(s)$ at the each point $\alpha(s)$ of α coincides with the plane spanned by $N^*(s^*), B_2^*(s^*)$ corresponding point $\beta(s^*) = \beta(\varphi(s))$ of β .

Theorem 3.3. *Let α be a quaternionic curve in E^4 with curvature functions $K, k, r - K$ and $r - K \neq 0$. Then α is a $(1, 3)$ -Bertrand curve if and only if there exist constant real numbers $\lambda, \mu, \gamma, \delta$ satisfying*

- (i) $\lambda k(s) - \mu(r - K)(s) \neq 0$
- (ii) $\lambda K(s) + \gamma(\lambda k(s) - \mu(r - K)(s)) = 1$
- (iii) $\gamma K(s) - k(s) = \delta(r - K)(s)$
- (iv) $\left\{ (\gamma^2 - 1)k(s)K(s) + \gamma \left[(K(s))^2 - (k(s))^2 - (r - K)^2(s) \right] \right\} \neq 0$
for all $s \in I$.

Proof. We assume that α is a $(1, 3)$ -Bertrand curve parametrized by arclength s . The $(1, 3)$ -Bertrand mate β is given by

$$\beta(s^*) = \beta(\varphi(s)) = \alpha(s) + \lambda(s)N(s) + \mu(s)B_2(s) \quad (8)$$

for all $s \in I$. Here λ and μ are C^∞ -function on I , and s^* is the arclength parameter of β . Differentiating (8) with respect to s , and using the Frenet equations, we obtain

$$\begin{aligned} \varphi'(s)T^*(\varphi(s)) &= (1 - \lambda(s)K(s))T(s) + \lambda'(s)N(s) \\ &\quad + (\lambda(s)k(s) - \mu(s)(r - K)(s))B_1(s) + \mu'(s)B_2(s) \end{aligned}$$

for all $s \in I$.

Since the plane spanned by $N(s)$ and $B_2(s)$ coincides with the plane spanned by $N^*(\varphi(s))$ and $B_2^*(\varphi(s))$, we can put

$$N^*(\varphi(s)) = (\cos \theta(s))N(s) + (\sin \theta(s))B_2(s) \quad (9)$$

$$B_2^*(\varphi(s)) = (-\sin \theta(s))N(s) + (\cos \theta(s))B_2(s) \quad (10)$$

and we notice that $\sin \theta(s) \neq 0$ for all $s \in I$. By the following facts

$$0 = \langle \varphi'(s)T^*(\varphi(s)), N^*(\varphi(s)) \rangle = \lambda'(s)\cos \theta(s) + \mu'(s)\sin \theta(s)$$

$$0 = \langle \varphi'(s)T^*(\varphi(s)), B_2^*(\varphi(s)) \rangle = -\lambda'(s)\sin \theta(s) + \mu'(s)\cos \theta(s)$$

we obtain

$$\lambda'(s) = 0, \mu'(s) = 0$$

that is, λ and μ are constant function on I with values λ and μ , respectively. Therefore, for all $s \in I$, (8) is rewritten as

$$\beta(s^*) = \beta(\varphi(s)) = \alpha(s) + \lambda N(s) + \mu B_2(s) \quad (11)$$

and we obtain

$$\varphi'(s)T^*(\varphi(s)) = (1 - \lambda K(s))T(s) + (\lambda k(s) - \mu(r - K)(s))B_1(s). \quad (12)$$

Here we notice that

$$(\varphi'(s))^2 = (1 - \lambda K(s))^2 + (\lambda k(s) - \mu(r - K)(s))^2 \neq 0 \quad (13)$$

for all $s \in I$. Thus we can set

$$T^*(\varphi(s)) = (\cos \tau(s))T(s) + (\sin \tau(s))B_1(s) \quad (14)$$

and

$$\cos \tau(s) = \frac{1 - \lambda K(s)}{\varphi'(s)} \quad (15)$$

$$\sin \tau(s) = \frac{\lambda k(s) - \mu(r - K)(s)}{\varphi'(s)} \quad (16)$$

where τ is a C^∞ -function on I . Differentiating (14) with respect to s and using the Frenet equations, we obtain

$$\begin{aligned} \bar{K}(\varphi(s))\varphi'(s)N^*(\varphi(s)) &= \frac{d\cos\tau(s)}{ds}T(s) \\ &+ (K(s)\cos\tau(s) - k(s)\sin\tau(s))N(s) \\ &+ \frac{d\sin\tau(s)}{ds}B_1(s) + (r - K(s))\sin\tau(s)B_2(s). \end{aligned}$$

Since $N^*(\varphi(s))$ is expressed by linear combination of $N(s)$ and $B_2(s)$, it holds that

$$\frac{d\cos\tau(s)}{ds} = 0, \quad \frac{d\sin\tau(s)}{ds} = 0,$$

that is, τ is a constant function on I with value τ_0 . Thus we obtain

$$T^*(\varphi(s)) = (\cos\tau_0)T(s) + (\sin\tau_0)B_1(s) \quad (17)$$

$$\varphi'(s)\cos\tau_0 = 1 - \lambda K(s) \quad (18)$$

$$\varphi'(s)\sin\tau_0 = \lambda k(s) - \mu(r - K(s)) \quad (19)$$

for all $s \in I$. There fore we obtain

$$(1 - \lambda K(s))\sin\tau_0 = (\lambda k(s) - \mu(r - K(s)))\cos\tau_0 \quad (20)$$

for all $s \in I$.

If $\sin\tau_0 = 0$, then it holds $\cos\tau_0 = \mp 1$. Thus (17) implies that $T^*(\varphi(s)) = \mp T(s)$. Differentiating this equality, we obtain

$$\bar{K}(\varphi(s))\varphi'(s)N^*(\varphi(s)) = \mp K(s)N(s),$$

that is,

$$N^*(\varphi(s)) = \mp N(s),$$

for all $s \in I$. By Theorem 3.1, this fact is a contradiction. Thus we must consider only the case of $\sin\tau_0 \neq 0$. Then (19) implies

$$\lambda k(s) - \mu(r - K(s)) \neq 0$$

that is, we obtain the relation (i).

The fact $\sin\tau_0 \neq 0$ and (20) imply

$$\lambda K(s) + \sin^{-1}\tau_0 \cos\tau_0 (\lambda k(s) - \mu(r - K(s))) = 1.$$

From this, we obtain

$$\lambda K(s) + \gamma (\lambda k(s) - \mu(r - K(s))) = 1$$

for all $s \in I$, where $\gamma = \sin^{-1}\tau_0 \cos\tau_0$ is a constant number. Thus we obtain the relation (ii).

Differentiating (17) with respect to s and using the Frenet equations, we obtain

$$\begin{aligned} \bar{K}(\varphi(s))\varphi'(s)N^*(\varphi(s)) &= (K(s)\cos\tau_0 - k(s)\sin\tau_0)N(s) \\ &+ (r - K(s))\sin\tau_0 B_2(s) \end{aligned}$$

for all $s \in I$. From the above equality, (18), (19) and (b), we obtain

$$\begin{aligned} [\bar{K}(\varphi(s))\varphi'(s)]^2 &= [K(s)\cos\tau_0 - k(s)\sin\tau_0]^2 + [(r - K(s))\sin\tau_0]^2 \\ &= (\lambda k(s) - \mu(r - K(s)))^2 \\ &\quad \times \left[(\gamma K(s) - k(s))^2 + ((r - K(s)))^2 \right] (\varphi'(s))^{-2} \end{aligned}$$

for all $s \in I$. From (13) and (ii), it holds

$$(\varphi'(s))^2 = (\gamma^2 + 1) (\lambda k(s) - \mu(r - K)(s))^2.$$

Thus we obtain

$$[\bar{K}(\varphi(s))\varphi'(s)]^2 = \frac{1}{\gamma^2 + 1} [(\gamma K(s) - k(s))^2 + ((r - K)(s))^2]. \quad (21)$$

By (18), (19) and (ii), we can set

$$N^*(\varphi(s)) = (\cos \eta(s))N(s) + (\sin \eta(s))B_2(s), \quad (22)$$

where

$$\cos \eta(s) = \frac{(\lambda k(s) - \mu(r - K)(s))(\gamma K(s) - k(s))}{\bar{K}(\varphi(s))(\varphi'(s))^2} \quad (23)$$

$$\sin \eta(s) = \frac{(r - K)(s)(\lambda k(s) - \mu(r - K)(s))}{\bar{K}(\varphi(s))(\varphi'(s))^2} \quad (24)$$

for all $s \in I$. Here, η is a C^∞ -function on I .

Differentiating (22) with respect to s and using Frenet equations, we get

$$\begin{aligned} & -\varphi'(s)\bar{K}(\varphi(s))T^*(\varphi(s)) + \varphi'(s)\bar{k}(\varphi(s))B_1^*(\varphi(s)) \\ = & -K(s)\cos \eta(s)T(s) + \frac{d\cos \eta(s)}{ds}N(s) \\ & + \{k(s)\cos \eta(s) - (r - K)(s)\sin \eta(s)\}B_1(s) \\ & + \frac{d\sin \eta(s)}{ds}B_2(s) \end{aligned}$$

for all $s \in I$. From the above fact, it holds

$$\frac{d\cos \eta(s)}{ds} = 0, \quad \frac{d\sin \eta(s)}{ds} = 0$$

that is, η is a constant function on I with value η_0 . Let $\delta = (\cos \eta_0)(\sin \eta_0)^{-1}$ be constant number. Then (23) and (24) imply

$$\gamma K(s) - k(s) = \delta(r - K)(s)$$

that is we obtain the relation (iii).

Moreover, we obtain

$$\begin{aligned} & -\varphi'(s)\bar{K}(\varphi(s))T^*(\varphi(s)) + \varphi'(s)\bar{k}(\varphi(s))B_1^*(\varphi(s)) \\ = & -K(s)\cos \eta(s)T(s) \\ & + \{k(s)\cos \eta(s) - (r - K)(s)\sin \eta(s)\}B_1(s). \end{aligned}$$

By the above equality and (12), we obtain

$$\begin{aligned} \varphi'(s)\bar{k}(\varphi(s))B_1^*(\varphi(s)) & = \varphi'(s)\bar{K}(\varphi(s))T^*(\varphi(s)) \\ & \quad - K(s)\cos \eta_0 T(s) \\ & \quad + \{k(s)\cos \eta_0 - (r - K)(s)\sin \eta_0\}B_1(s) \\ & = \left\{(\varphi'(s))^2\bar{K}(\varphi(s))\right\}^{-1} \\ & \quad \times \{A(s)T(s) + B(s)B_1(s)\}, \end{aligned}$$

where

$$\begin{aligned} A(s) &= \{\varphi'(s) \bar{K}(\varphi(s))\}^2 (1 - \lambda K(s)) \\ &\quad - K(s) (\lambda k(s) - \mu(r - K)(s)) (\gamma K(s) - k(s)) \\ B(s) &= \{\varphi'(s) \bar{K}(\varphi(s))\}^2 (\lambda k(s) - \mu(r - K)(s)) \\ &\quad + \left\{ \begin{array}{l} k(s) (\lambda k(s) - \mu(r - K)(s)) (\gamma K(s) - k(s)) \\ - (r - K)^2(s) (\lambda k(s) - \mu(r - K)(s)) \end{array} \right\} \end{aligned}$$

From (ii) and (21), $A(s)$ and $B(s)$ are rewritten as:

$$\begin{aligned} A(s) &= (\lambda k(s) - \mu(r - K)(s)) (\gamma^2 + 1)^{-1} \\ &\quad \times \left\{ (1 - \gamma^2) K(s) k(s) - \gamma \left[(K(s))^2 - (k(s))^2 - ((r - K)(s))^2 \right] \right\} \\ B(s) &= \gamma (\gamma^2 + 1)^{-1} (\lambda k(s) - \mu(r - K)(s)) \\ &\quad \times \left\{ (\gamma^2 - 1) k(s) K(s) + \gamma \left[(K(s))^2 - (k(s))^2 - (r - K)^2(s) \right] \right\}. \end{aligned}$$

Since $\bar{K}(\varphi(s)) \varphi'(s) N^*(\varphi(s)) \neq 0$ for all $s \in I$, it holds

$$\left\{ (\gamma^2 - 1) k(s) K(s) + \gamma \left[(K(s))^2 - (k(s))^2 - (r - K)^2(s) \right] \right\} \neq 0$$

for all $s \in I$. Thus we obtain the relation (iv).

We assume that α ($\alpha : I \rightarrow Q_H$) is a C^∞ special Frenet curve in Q_H with curvature functions K, k and $(r - K)$ satisfying the relation (i), (ii), (iii) and (iv) for constant numvers λ, μ, γ and δ . Then we define a C^∞ -curve β by

$$\beta(s) = \alpha(s) + \lambda N(s) + \mu B_2(s) \quad (25)$$

for all $s \in I$, where s is arclenght parameter of α . Differentiating (25) with respect to s and using the Frenet equations, we obtain

$$\frac{d\beta(s)}{ds} = (1 - \lambda K(s)) T(s) + (\lambda k(s) - \mu(r - K)) B_1(s)$$

for all $s \in I$. Thus, by the relation (ii), we obtain

$$\frac{d\beta(s)}{ds} = (\lambda k(s) - \mu(r - K)(s)) [\gamma T(s) + B_1(s)]$$

for all $s \in I$. Since the relation (i) holds, the curve β is a regular curve. Then there exists a regular map $\varphi : I \rightarrow \bar{I}$ defined by

$$s^* = \varphi(s) = \int_0^s \left\| \frac{d\beta(t)}{dt} \right\| dt$$

where s^* denotes the arclenght parameter of β , and we obtain

$$\varphi'(s) = \varepsilon \sqrt{\gamma^2 + 1} (\lambda k(s) - \mu(r - K)(s)) > 0, \quad (26)$$

where $\varepsilon = 1$ if $\lambda k(s) - \mu(r - K)(s) > 0$, and $\varepsilon = -1$ if $\lambda k(s) - \mu(r - K)(s) < 0$. Thus the curve β is rewritten as

$$\begin{aligned} \beta(s^*) &= \beta(\varphi(s)) \\ &= \alpha(s) + \lambda N(s) + \mu B_2(s) \end{aligned}$$

for all $s \in I$. Differentiating the above equality with respect to s , we obtain

$$\varphi'(s) \frac{d\beta(s^*)}{ds^*} = (\lambda k(s) - \mu(r - K)(s)) [\gamma T(s) + B_1(s)]. \quad (27)$$

We can define the unit vector field T^* along β by $T^*(s^*) = \frac{d\beta(s^*)}{ds^*}$ for all $s^* \in \bar{I}$. By (26) and (27), we obtain

$$T^*(\varphi(s)) = \varepsilon (\gamma^2 + 1)^{-\frac{1}{2}} [\gamma T(s) + B_1(s)] \quad (28)$$

for all $s \in I$. Differentiating (28) with respect to s and using the Frenet equations, we obtain

$$\varphi'(s) \frac{dT^*(s^*)}{ds^*} \Big|_{s^*=\varphi(s)} = \varepsilon (\gamma^2 + 1)^{-\frac{1}{2}} \{[\gamma K(s) - k(s)] N(s) + (r - K)(s) B_2(s)\}$$

and

$$\left\| \frac{dT^*(s^*)}{ds^*} \Big|_{s^*=\varphi(s)} \right\| = \frac{\sqrt{[\gamma K(s) - k(s)]^2 + ((r - K)(s))^2}}{\varphi'(s) \sqrt{\gamma^2 + 1}}.$$

By the fact that $(r - K)(s) \neq 0$ for all $s \in I$, we obtain

$$\bar{K}(\varphi(s)) = \left\| \frac{dT^*(s^*)}{ds^*} \Big|_{s^*=\varphi(s)} \right\| > 0 \quad (29)$$

for all $s \in I$. Then we can define a unit vector field N^* along β by

$$\begin{aligned} N^*(s^*) &= N^*(\varphi(s)) \\ &= \frac{1}{\bar{K}(\varphi(s))} \frac{dT^*(s^*)}{ds^*} \Big|_{s^*=\varphi(s)} \\ &= \frac{1}{\sqrt{[\gamma K(s) - k(s)]^2 + ((r - K)(s))^2}} \\ &\quad \times \{[\gamma K(s) - k(s)] N(s) + (r - K)(s) B_2(s)\} \end{aligned}$$

for all $s \in I$. Thus we can put

$$N^*(\varphi(s)) = \cos \xi(s) N(s) + \sin \xi(s) B_2(s) \quad (30)$$

where

$$\cos \xi(s) = \frac{\gamma K(s) - k(s)}{\sqrt{[\gamma K(s) - k(s)]^2 + ((r - K)(s))^2}} \quad (31)$$

$$\sin \xi(s) = \frac{(r - K)(s)}{\sqrt{[\gamma K(s) - k(s)]^2 + ((r - K)(s))^2}} \quad (32)$$

for all $s \in I$. Here ξ is a C^∞ function on I . Differentiating (30) with respect to s and using the Frenet equations, we obtain

$$\begin{aligned} \varphi'(s) \frac{dN^*(s^*)}{ds^*} \Big|_{s^*=\varphi(s)} &= -\cos \xi(s) K(s) T(s) \\ &\quad + \frac{d \cos \xi(s)}{ds} N(s) \\ &\quad + (k(s) \cos \xi(s) - (r - K)(s) \sin \xi(s)) B_1(s) \\ &\quad + \frac{d \sin \xi(s)}{ds} B_2(s). \end{aligned}$$

Differentiatin (iii) with respect to s , we obtain

$$(\gamma K'(s) - k'(s))(r - K(s)) - (\gamma K(s) - k(s))(r - K)'(s) = 0. \quad (33)$$

Differentiating (31) and (32) with respect to s and using (33), we obtain

$$\frac{d \cos \xi(s)}{ds} = 0, \quad \frac{d \sin \xi(s)}{ds} = 0$$

that is, ξ is a constant function on I with value ξ_0 . Thus we obtain

$$\cos \xi_0 = \frac{\gamma K(s) - k(s)}{\varepsilon \sqrt{[\gamma K(s) - k(s)]^2 + ((r - K)(s))^2}} \quad (34)$$

$$\sin \xi_0 = \frac{(r - K)(s)}{\varepsilon \sqrt{[\gamma K(s) - k(s)]^2 + ((r - K)(s))^2}} \neq 0. \quad (35)$$

From (30), it holds

$$N^*(\varphi(s)) = \cos \xi_0 N(s) + \sin \xi_0 B_2(s). \quad (36)$$

Thus we obtain, by (28) and (29),

$$\begin{aligned} \bar{K}(\varphi(s)) T^*(\varphi(s)) &= \frac{[\gamma K(s) - k(s)]^2 + ((r - K)(s))^2}{\varepsilon \varphi'(s) (\gamma^2 + 1) \sqrt{[\gamma K(s) - k(s)]^2 + ((r - K)(s))^2}} \\ &\quad \times [\gamma T(s) + B_1(s)] \end{aligned}$$

and by (34),(35) and (36)

$$\begin{aligned} \frac{dN^*(s^*)}{ds^*} \Big|_{s^*=\varphi(s)} &= - \frac{K(s) (\gamma K(s) - k(s))}{\varepsilon \varphi'(s) \sqrt{[\gamma K(s) - k(s)]^2 + ((r - K)(s))^2}} T(s) \\ &\quad + \left(\frac{k(s) (\gamma K(s) - k(s)) - ((r - K)(s))^2}{\varepsilon \varphi'(s) \sqrt{[\gamma K(s) - k(s)]^2 + ((r - K)(s))^2}} \right) B_1(s) \end{aligned}$$

for all $s \in I$. By the above equalities, we obtain

$$\frac{dN^*(s^*)}{ds^*} \Big|_{s^*=\varphi(s)} + \bar{K}(\varphi(s)) T^*(\varphi(s)) = \frac{P(s)}{R(s)} T(s) + \frac{Q(s)}{R(s)} B_1(s),$$

where

$$P(s) = - \left[\gamma \left\{ (K(s))^2 - (k(s))^2 - ((r - K)(s))^2 \right\} + (\gamma^2 - 1) K(s) k(s) \right]$$

$$Q(s) = \gamma \left[\gamma \left\{ (K(s))^2 - (k(s))^2 - ((r - K)(s))^2 \right\} + (\gamma^2 - 1) K(s) k(s) \right]$$

$$R(s) = \varepsilon \varphi'(s) (\gamma^2 + 1) \sqrt{[\gamma K(s) - k(s)]^2 + ((r - K)(s))^2} \neq 0$$

for all $s \in I$. We notice that, by (iii), $P(s) \neq 0$ for all $s \in I$. Thus we obtain

$$\begin{aligned} \bar{k}(\varphi(s)) &= \left\| \frac{dN^*(s^*)}{ds^*} \Big|_{s^*=\varphi(s)} + \bar{K}(\varphi(s)) T^*(\varphi(s)) \right\| \\ &= \frac{\left| \gamma \left\{ (K(s))^2 - (k(s))^2 - ((r - K)(s))^2 \right\} + (\gamma^2 - 1) K(s) k(s) \right|}{\varphi'(s) \sqrt{\gamma^2 + 1} \sqrt{[\gamma K(s) - k(s)]^2 + ((r - K)(s))^2}} \end{aligned}$$

for all $s \in I$. Thus we can define a unit vector field $B_1^*(s^*)$ along β by

$$\begin{aligned} B_1^*(s^*) &= B_1^*(\varphi(s)) \\ &= \frac{1}{k(\varphi(s))} \left(\frac{dN^*(s^*)}{ds^*} \Big|_{s^*=\varphi(s)} + \bar{K}(\varphi(s)) T^*(\varphi(s)) \right) \end{aligned}$$

that is,

$$B_1^*(\varphi(s)) = \frac{1}{\varepsilon\sqrt{\gamma^2+1}} (-T(s) + \gamma B_2(s)) \quad (37)$$

for all $s \in I$. Next we can define a unit vector field B_2^* along β by

$$\begin{aligned} B_2^*(s^*) &= B_2^*(\varphi(s)) \\ &= \frac{1}{\varepsilon\sqrt{[\gamma K(s) - k(s)]^2 + ((r - K)(s))^2}} \\ &\quad \{- (r - K)(s) N(s) + (\gamma K(s) - k(s)) B_2(s)\} \end{aligned}$$

that is

$$B_2^*(\varphi(s)) = -\sin \xi_0 N(s) + \cos \xi_0 B_2(s) \quad (38)$$

for all $s \in I$. Now we obtain, by (28), (36), (37) and (38),

$$\det [T^*(\varphi(s)), N^*(\varphi(s)), B_1^*(\varphi(s)), B_2^*(\varphi(s))] = 1$$

and $\{T^*(\varphi(s)), N^*(\varphi(s)), B_1^*(\varphi(s)), B_2^*(\varphi(s))\}$ is orthonormal for all $s \in I$. Thus the frame $\{T^*, N^*, B_1^*, B_2^*\}$ along β is of orthonormal and of positive. And we obtain

$$\begin{aligned} \overline{(r - K)}(s) &= \left\langle \frac{dB_1^*(s^*)}{ds^*} \Big|_{s^*=\varphi(s)}, B_2^*(\varphi(s)) \right\rangle \\ &= \frac{\sqrt{\gamma^2+1} K(s) (r - K)(s)}{\varphi'(s) \sqrt{[\gamma K(s) - k(s)]^2 + ((r - K)(s))^2}} \neq 0 \end{aligned}$$

for all $s \in I$. Thus the curve β is a Frenet curve in Q_H . And it is trivial that the plane spanned by $N(s), B_2(s)$ at the each point $\alpha(s)$ of α coincides with the plane spanned by $N^*(s^*), B_2^*(s^*)$ corresponding point $\beta(s^*) = \beta(\varphi(s))$ of β . Therefore α is a (1, 3)-Bertrand curve in Q_H . \square

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