

ROBUST CONTROL PROBLEMS OF VORTEX DYNAMICS IN SUPERCONDUCTING FILMS WITH GINZBURG-LANDAU COMPLEX SYSTEMS

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We formulate and study robust control problems for a two-dimensional time-dependent Ginzburg-Landau model with Robin boundary conditions on phase-field parameter, which describes the phase transitions taking place in superconductor films with variable thickness. The objective of such study is to control the motion of vortices in the superconductor films by taking into account the influence of noises in data. Firstly, we introduce the perturbation problem of the nonlinear governing coupled system of equations (the deviation from the desired target). The existence and the uniqueness of the solution of the perturbation are proved as well as stability under mild assumptions. Afterwards, the robust control problems are formulated in the case when the control is in the external magnetic field and in the case when the control is in the initial condition of the vector potential. We show the existence of an optimal solution, and we also find necessary conditions for a saddle point optimality.

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1. Introduction

The aim of this contribution is the study of robust control problems to describe the phenomenon of vortex structure in the superconducting phase transitions, using the time-dependent Ginzburg-Landau (TDGL) complex superconductivity model. Such model was derived by Gor'kov and Eliashberg in [28] from the microscopic BCS (Bardeen-Cooper-Schrieffer) theory [4] for a superconductor with paramagnetic impurities. It involves the real vector potential U for the total magnetic field and a complex phase-field variable ϕ so that $|\phi|^2 = \phi\bar{\phi}$ ($\bar{\phi}$ is the complex conjugate of ϕ) gives the relative density of the superconducting charge carriers (Cooper pairs of electrons), which varies between 0 in the normal phase and 1 in the superconducting phase. The need for ϕ to be complex is associated with the macroscopic quantum nature of superconductivity. Here we will be connected with the response of a superconducting material to an applied magnetic field under isothermal conditions below its critical temperature T_c (the transition from

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normally conducting to superconducting is usually associated with critical temperature). The time evolution of (U, ϕ) is then governed by the following system:

$$\begin{aligned} \eta\rho \frac{\partial\phi}{\partial t} - i\eta\kappa \operatorname{div}(\rho U)\phi + b(U) \cdot (\rho b(U))(\phi) + \rho G(\phi) &= 0 \quad \text{in } \mathfrak{Q} = \Omega \times (0, T), \\ \rho \frac{\partial U}{\partial t} + \operatorname{curl}(\rho \operatorname{curl}(U)) - \nabla(\operatorname{div}(\rho U)) + \rho \mathcal{R}(b(U)(\phi)\bar{\phi}) &= \operatorname{curl}(\rho H) \quad \text{in } \mathfrak{Q}, \end{aligned} \quad (1.1a)$$

subjected to the Robin-type boundary conditions

$$\frac{1}{\kappa^2} \frac{\partial\phi}{\partial n} = \mu\phi, \quad U \cdot \mathbf{n} = 0, \quad \operatorname{curl}(U) = H \quad \text{in } \Sigma = \partial\Omega \times (0, T), \quad (1.1b)$$

and the initial conditions

$$\phi(0) = \phi_0, \quad U(0) = U_0 \quad \text{in } \Omega, \quad (1.1c)$$

where the operator b (the covariant derivative) and the function G are defined by: $b(U) = ((i/\kappa)\nabla + U)$, $\bar{b}(U) = ((-i/\kappa)\nabla + U)$ and $G(z) = (|z|^2 - 1)z$. The domain Ω is an open bounded domain in \mathbb{R}^2 with Lipschitz boundary $\partial\Omega$, \mathbf{n} is the unit normal to the surface of the superconductor $\Gamma = \partial\Omega$ and μ is an arbitrary real number (the boundary condition is appropriate for the superconductor interface with vacuum or an insulator if $\mu = 0$ and for the superconductor interface with normal metal if $\mu \neq 0$). $\mathcal{R}(\cdot)$ (resp., $\mathcal{I}(\cdot)$) denotes the real part (resp., the imaginary part) of the quantity in (\cdot) and curl denote the curl operators defined by (on the (x, y) -plane)

$$\begin{aligned} \operatorname{curl}(\phi) &= \left(\frac{\partial\phi}{\partial y}, -\frac{\partial\phi}{\partial x} \right)^T \quad (\phi \text{ is a scalar}) \\ \operatorname{curl}(U) &= \frac{\partial u_2}{\partial x} - \frac{\partial u_1}{\partial y} \quad (U = (u_1, u_2) \text{ is a vector}). \end{aligned} \quad (1.2)$$

H is the applied magnetic field, η is the nondimensional diffusivity and $\rho > 0$ is a smooth function characterizing the vertical shape of the superconducting films and satisfying the following hypothesis

$$\rho \in C^1(\bar{\Omega}) \quad \text{such that } \rho_0 \leq \rho \leq \rho_1 \text{ and } |\nabla\rho| \leq \rho_d, \quad (1.3)$$

where (ρ_0, ρ_1, ρ_d) are nonnegative constants.

The positive constant κ is the Ginzburg-Landau parameter with $\kappa = \lambda/\xi$, where ξ is a coherence length describing the size of thermodynamic fluctuations in the superconducting phase, and λ is the London penetration depth describing the depth to which an external magnetic field can penetrate the superconductor. The parameter κ determine the type of superconducting material: $\kappa < 1/\sqrt{2}$ describes type-I superconductors, $\kappa > 1/\sqrt{2}$ describes type-II superconductors.

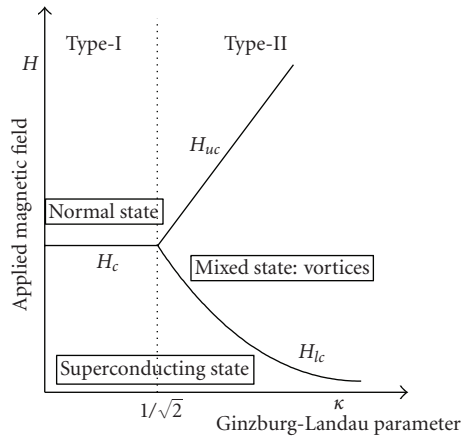


Figure 1.1. The reaction of a superconducting material.

Most applications of superconducting materials involve type-II superconductors in high magnetic fields. It is known that for the type-II superconductors, there is a critical magnetic field which splits into a lower critical field H_{lc} and an upper critical field H_{uc} (cf. Figure 1.1). For the magnetic fields below H_{lc} the material is in the superconducting state and for magnetic field above H_{uc} the material is in the normal state. For the magnetic fields between H_{lc} and H_{uc} the material is in the mixed state. This mixed state is described by physicists as follows: around some isolated points (called vortices, which are most commonly arranged in a hexagonal arrangement, see Abrikosov [1]) inside the material, the superconducting property is destroyed and the magnetic field become stronger in the nearby regions surrounding these vortices. While elsewhere, the superconducting property is still dominant and the field magnetic is excluded. Moreover, the motion of the vortices depends highly on the magnetization processes of the material (this is the result of the “Lorentz force”—that is more than anchorage forces of vortices and causes the displacement of the vortex—on the magnetic flux line carried by the vortex due to the transport current, e.g.). The motion of the vortices is undesirable, because this motion dissipate energy and leads to an electric field. So it is very interesting to study the applied magnetic in order to prevent their motion.

The objective of a robust control theory, which generalizes optimal control theory, is to compensate the undesirable effects of system disturbances through control actions such that a cost function achieves its minimum for the worst disturbances, that is, to find the best control which takes into account the worst-case (maximal) disturbance. More recently robust control frameworks have been the object of numerous studies either from a theoretical or from a numerical point of view to some classes of infinite (or finite)-dimensional linear or nonlinear systems (see, e.g., [5–8, 10, 18, 24–26, 29, 31, 36, 37, 40] and the references therein).

Various problems associated with the Ginzburg-Landau models in superconductivity have been studied these last years (the literature on such model is vast, see, for example, [9, 11–13, 15, 17, 19–22, 30, 38, 39] and the references therein). For the optimal control

problems associated with the TDGL models, we can mention [14] in which the authors studied the control of the vortices in superconducting films through the external magnetic field. Here, we consider a robust control problem, for the TDGL models with a Robin boundary condition on phase-field variable, which describes the phase transitions taking place in superconductors films, in order to take account the influence of noises in data. Indeed, such perturbations (noises) have the effect of impeding the ability of the material to become superconducting.

1.1. Assumptions and notations. We denote by $V_n = \{U \in H^1(\Omega); U \cdot \mathbf{n} = 0 \text{ on } \Omega\}$ and V'_n the dual of V_n . We denote by $\langle \cdot, \cdot \rangle_{V'_n, V_n}$ the duality product between V'_n and V_n . For any pair of real number $r, s \geq 0$, we introduce the Sobolev space $H^{r,s}(\mathcal{Q})$ defined by $H^{r,s}(\mathcal{Q}) = L^2(0, T, H^r(\Omega)) \cap H^s(0, T, L^2(\Omega))$, which is a Hilbert space normed by

$$\left(\int_0^T \|v\|_{H^r(\Omega)}^2 dt + \|v\|_{H^s(0, T, L^2(\Omega))}^2 \right)^{1/2}, \quad (1.4)$$

where $H^s(0, T, L^2(\Omega))$ denotes the Sobolev space of order s of functions defined on $(0, T)$ and taking values in $L^2(\Omega)$ and defined by $H^s(0, T, L^2(\Omega)) = [H^m(0, T, L^2(\Omega)), L^2(\mathcal{Q})]_\theta$, where $\theta \in (0, 1)$, $s = (1 - \theta)m$, m is an integer and $H^m(0, T, L^2(\Omega)) = \{v \in L^2(\mathcal{Q}) \mid \partial^j v / \partial t^j \in L^2(\mathcal{Q}), \forall j = 1, m\}$.

Remark 1.1. (i) According to [27], we have the following embedding inequality on V_n :

$$\|U\|_{H^1}^2 \leq C \left(\|U\|_{L^2}^2 + \|\operatorname{div}(U)\|_{L^2}^2 + \|\operatorname{curl}(U)\|_{L^2}^2 \right), \quad \forall U \in V_n. \quad (1.5)$$

(ii) For $v \in H^{r,s}(\mathcal{Q})$, the trace functions of v : $\partial^j v / \partial n^j$ on $\Sigma = \partial \Omega \times (0, T)$ for an integer j such that $j \in [0, r - (1/2)]$ exist and satisfy $\partial^j v / \partial n^j \in H^{r_j, s_j}(\Sigma)$, where $r_j = r - j - 1/2$ and $s_j = s(r - j - 1/2)/r$. Moreover the functions $v \rightarrow \partial^j v / \partial n^j$ are continuous linear mappings from $H^{r,s}(\mathcal{Q})$ into $H^{r_j, s_j}(\Sigma)$ (see, e.g., [34, 35]).

If X denotes some Banach space of real-valued functions, the corresponding space of complex-valued functions will be denoted by \mathcal{X} and the corresponding space of vector-valued functions, each of components belonging to X , will be denoted by \mathbf{X} , and we use $\|\cdot\|_X$ to denote the norms of spaces X , \mathbf{X} or \mathcal{X} .

We can now introduce the following spaces: $\mathbf{W}_n = L^2(0, T, V_n) \cap H^1(0, T, V'_n)$, $\mathbf{E}_n = L^2(0, T, V_n) \cap L^\infty(0, T, \mathbf{L}^2(\Omega))$, $\mathcal{W} = L^2(0, T, \mathcal{H}^1(\Omega)) \cap H^1(0, T, (\mathcal{H}^1)'(\Omega))$, $\mathcal{E} = L^2(0, T, \mathcal{H}^1(\Omega)) \cap L^\infty(0, T, \mathcal{L}^2(\Omega))$ and $\mathcal{L}_\infty^2(\Omega) = \{\phi \in \mathcal{L}^2(\Omega); |\phi| \text{ is bounded a.e. in } \Omega\}$.

Remark 1.2. (i) \mathcal{W} and \mathbf{W}_n are continuously embedded into $C^0([0, T], \mathcal{L}^2(\Omega))$ and $C^0([0, T], \mathbf{L}^2(\Omega))$, respectively (see, e.g., [34, 35]).

(ii) Although $\mathcal{L}_\infty^2(\Omega)$ is a subset of $\mathcal{L}^\infty(\Omega)$, we equipped this space by the standard norm of the space $\mathcal{L}^2(\Omega)$.

The weak formulation associated to problem (1.1) is then to find $(\phi, U) \in \mathcal{W} \times \mathbf{W}_n$ such that

$$\begin{aligned} & \eta \int_{\Omega} \rho \frac{\partial \phi}{\partial t} q dx - i\eta\kappa \int_{\Omega} \operatorname{div}(\rho U) \phi q dx + \int_{\Omega} \rho b(U)(\phi) \bar{b}(U)(q) dx \\ & - i\mu \int_{\Gamma} \rho \phi q d\Gamma + \int_{\Omega} \rho G(\phi) q dx = 0, \quad \text{a.e. } t \in (0, T), \quad \forall q \in \mathcal{H}^1(\Omega), \\ & \int_{\Omega} \rho \frac{\partial U}{\partial t} v dx + \int_{\Omega} \rho \operatorname{curl}(U) \operatorname{curl}(v) dx + \int_{\Omega} \operatorname{div}(\rho U) \operatorname{div}(v) dx \\ & + \int_{\Omega} \rho \mathcal{R}(b(U)(\phi) \bar{\phi}) v dx = \int_{\Omega} \rho H \operatorname{curl}(v) dx, \quad \text{a.e. } t \in (0, T), \quad \forall v \in V_n, \\ & \phi(0) = \phi_0, \quad U(0) = U_0 \quad \text{in } \Omega. \end{aligned} \tag{1.6}$$

1.2. Preliminary results

Definition 1.3. Let \mathcal{H}_1 and \mathcal{H}_2 be two arbitrary sets (whose nature will be stated precisely at each situation) and denote $\mathcal{H} = \mathcal{H}_1 \times \mathcal{H}_2$.

A pair $(f^*, g^*) \in \mathcal{H}$ is a saddle point of the cost function J on \mathcal{H} if

$$J(f^*, g) \leq J(f^*, g^*) \leq J(f, g^*), \quad \forall (f, g) \in \mathcal{H}. \tag{1.7}$$

LEMMA 1.4. Let $\Omega \subset \mathbb{R}^m$, $m \geq 1$, be an open and bounded set with a smooth boundary and let q be a nonnegative integer. There exist the following results.

(i) $H^q(\Omega) \subset L^p(\Omega)$, for all $p \in [1, 2m/(m-2q)]$, with continuous embedding (with the exception that if $2q = m$, then $p \in [1, +\infty[$ and if $2q > m$, then $p \in [1, +\infty]$).

(ii) (Gagliardo-Nirenberg inequalities). There exists $C > 0$ such that $\|v\|_{L^p} \leq C \|v\|_{H^q}^\theta \|v\|_{L^2}^{1-\theta}$, for all $v \in H^q(\Omega)$, where $0 \leq \theta < 1$ and $p = 2m/(m-2\theta q)$ (with the exception that if $q - m/2$ is a nonnegative integer, then θ is restricted to 0).

For the proof of this lemma, see for example Adams [2].

LEMMA 1.5. (i) For all $z \in \mathbb{C}$, $G(z)\bar{z} \in \mathbb{R}$.

(ii) For all $(z_1, z_2) \in \mathbb{C}^2$, $G(z_1) - G(z_2) = (|z_1|^2 + |z_2|^2 - 1)(z_1 - z_2) + z_1 z_2 \overline{(z_1 - z_2)}$.

LEMMA 1.6. For (φ, u) and (ψ, v) sufficiently regular,

(i) $b(u)(\varphi) = b(v)(\varphi) + (u - v)\varphi$;

(ii) $b(u)(\varphi) - b(u)(\psi) = b(u)(\varphi - \psi)$;

(iii) $b(u)(\varphi) \cdot \overline{b(u)(\varphi)} = (1/\kappa^2) |\nabla \varphi|^2 + |u|^2 |\varphi|^2 - (2/\kappa) \mathcal{F}(\overline{\varphi} \nabla \varphi) \cdot u$.

The proof of the previous lemmas are immediate.

LEMMA 1.7. For (u, v, w, X) sufficiently regular,

(i) $\|u\|_{H^1} \|v\|_{L^4} \|X\|_{L^4} \leq C_1 \|u\|_{H^1}^2 \|v\|_{L^2} + \delta \|\nabla X\|_{L^2}^2 + C_2 (\|v\|_{H^1} + \|v\|_{H^1}^2) \|X\|_{L^2}^2$, with δ chooses suitably at each situation,

(ii) $\|u\|_{L^4} \|v\|_{L^4} \|w\|_{L^4} \|X\|_{L^4} \leq C_1 \|u\|_{L^4}^2 \|v\|_{L^4}^2 + \gamma \|\nabla X\|_{L^2}^2 + C_2 (\|w\|_{L^4}^4 + \|w\|_{L^4}^2) \|X\|_{L^2}^2$, with γ chooses suitably at each situation.

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Proof. (i) By using Gagliardo-Nirenberg inequalities (Lemma 1.4, with $m = 2$), we obtain $\|u\|_{H^1} \|v\|_{L^4} \|X\|_{L^4} \leq c_1 \|u\|_{H^1}^2 \|v\|_{L^2} + c_2 \|v\|_{H^1} \|X\|_{L^2}^2 + c_3 \|v\|_{H^1} \|X\|_{L^2} \|\nabla X\|_{L^2}$ and then by using the Young's formula, we can deduce the result (i).

$$(ii) \|u\|_{L^4} \|v\|_{L^4} \|w\|_{L^4} \|X\|_{L^4} \leq c_1 \|u\|_{L^4}^2 \|v\|_{L^4}^2 + c_2 \|w\|_{L^4}^2 \|X\|_{L^4}^2.$$

By using Gagliardo-Nirenberg inequalities, we can deduce that $\|w\|_{L^4}^2 \|X\|_{L^4}^2 \leq c_3 \|w\|_{L^4}^2 \|X\|_{L^2}^2 + c_4 \|w\|_{L^4}^2 \|X\|_{L^2} \|\nabla X\|_{L^2}$.

According to Young's formula, we can deduce the result (ii). \square

1.3. Outline of paper. The rest of the paper is organized as follows. In Section 2, we prove the existence and uniqueness of the problem (1.1). In Section 3, we introduce the initial perturbation problem and prove the existence and the uniqueness of the perturbation and obtain a stability result. In Section 4, we study the Fréchet differentiability of the solution operator of the perturbation problem. This property is necessary to develop the robust control problem. In Section 5, we study the robust control problem corresponding to obtain the saddle point of the cost function J . The functional J is depending on the disturbance (or noise), the control and the perturbation solution in the domain Ω over the time interval under consideration $[0, T]$. The robust problem is formulated in two cases of control: firstly the control is in the external magnetic field and the disturbance is in the external magnetic field or in the initial condition of the order parameter variable, secondly the control is in the initial condition of the vector potential and the disturbance is in the external magnetic field or in the initial condition of the order parameter variable. We prove the existence of an optimal solution (saddle point), to the robust control problem under consideration, and give necessary optimality conditions. The optimality system is corresponding to identify the gradient of the cost function that is necessary to develop a numerical computation in order to solve the robust control problem.

2. Existence and uniqueness of the solution of the TDGL model

The following results concern the existence and the uniqueness of the solution of the Ginzburg-Landau model with Robin-type boundary conditions on phase-field parameter (1.1).

THEOREM 2.1. *For any $(\phi_0, U_0) \in \mathcal{L}^2(\Omega) \times \mathbf{L}^2(\Omega)$ satisfying $|\phi_0| \leq 1$ a.e. in Ω and $H \in L^2(\mathcal{Q})$, there exists a unique solution $(\phi, U) \in \mathcal{W} \times \mathbf{W}_n$ of (1.6) satisfying $|\phi| \leq 1$ a.e. in \mathcal{Q} . Moreover, the following estimation exists:*

$$\|\phi\|_{\mathcal{W}}^2 + \|U\|_{\mathbf{W}_n}^2 \leq C \left(\|H\|_{L^2(\mathcal{Q})}^2 + \|\phi_0\|_{L^2(\Omega)}^2 + \|U_0\|_{L^2(\Omega)}^2 \right). \quad (2.1)$$

Proof. The proof of this result is obtained by using the same technique as in, for example, [14, 15]. Here, we sketch only the proof of the existence by using the method of lines. The uniqueness is obtained by using the same technique as to prove the stability result in Theorem 3.1.

To begin with, we introduce the following semidiscretized approximation problem: let $N \geq 1$ be an integer, let $\tau = T/N$ be the time size, $t_j = j\tau$, for $j = 0, \dots, N$. The

approximation problem is then to find $(\phi_j, u_j) \in \mathcal{H}^1(\Omega) \times V_n$, $j = 1, N$, such that

$$\begin{aligned} & \eta \int_{\Omega} \rho \frac{\phi_j - \phi_{j-1}}{\tau} q dx - i\eta\kappa \int_{\Omega} \operatorname{div}(\rho U_j) \phi_j q dx + \int_{\Omega} \rho b(U_j)(\phi_j) \bar{b}(U_j)(q) dx \\ & \quad - i\mu \int_{\Gamma} \rho \phi_j q d\Gamma + \int_{\Omega} \rho G(\phi_j) q dx = 0, \\ & \int_{\Omega} \rho \frac{U_j - U_{j-1}}{\tau} v dx + \int_{\Omega} \rho \mathcal{R}(\overline{\phi_{j-1}} b(U_{j-1})(\phi_{j-1})) v dx + \int_{\Omega} \rho \operatorname{curl}(U_j) \operatorname{curl}(v) dx \\ & \quad + \int_{\Omega} \operatorname{div}(\rho U_j) \operatorname{div}(v) dx = \int_{\Omega} \rho H_j \operatorname{curl}(v) dx, \end{aligned} \quad (2.2)$$

where (ϕ_0, U_0) is given by (1.6), and $H_j(\cdot) = (1/\tau) \int_{t_{j-1}}^{t_j} H(\cdot, t) dt$, $j = 1, N$.

We first notice that the first part and the second part of system (2.2) are independent of each other. The second part of (2.2) is a linear elliptic problem for U_j with (U_{j-1}, ϕ_{j-1}) given by the previous step. By using the standard argument we obtain the existence and the uniqueness of $U_j \in V_n$. The first part of (2.2) is a semilinear elliptic problem with respect to ϕ_j . By using the regularity of (U_j, ϕ_{j-1}) and a standard argument, we obtain the existence and the uniqueness of ϕ_j .

Prove now, by using the maximum principal, that $|\phi_j| \leq 1$, $\forall j \geq 1$.

Let us consider the following notation: $r^+ = \max(r, 0)$ and $r^- = (-r)^+$ and then $r = r^+ - r^-$. Prove now that if $|\phi_0| \leq 1$ a.e. in Ω then $|\phi_j| \leq 1$, $\forall j \geq 1$ a.e. in Ω . Suppose now that $|\phi_{j-1}| \leq 1$ a.e. in Ω , and prove that $|\phi_j| \leq 1$ a.e. in Ω . By choosing $q = (|\phi_j|^2 - 1)^+ \bar{\phi}_j$ in the first part of (2.2) and by taking the real part we have then

$$\begin{aligned} & \eta \int_{\Omega} \rho \frac{((|\phi_j|^2 - 1)^+)^2}{\tau} dx + \eta \int_{\Omega} \rho \frac{(|\phi_j|^2 - 1)^+}{\tau} dx - \eta \int_{\Omega} \rho \mathcal{R}(\bar{\phi}_j \phi_{j-1}) \frac{(|\phi_j|^2 - 1)^+}{\tau} dx \\ & \quad + \int_{\Omega} \rho |b(U_j)(\phi_j)|^2 (|\phi_j|^2 - 1)^+ dx \\ & \quad + \int_{\Omega^+} \frac{\rho}{2\kappa^2} |\nabla |\phi_j|^2|^2 dx + \int_{\Omega} \rho |\phi_j|^2 ((|\phi_j|^2 - 1)^+)^2 dx = 0, \end{aligned} \quad (2.3)$$

where $\Omega^+ = \{x \in \Omega \mid |\phi_j|^2 - 1 > 0\}$. This implies (since $|\phi_{j-1}| \leq 1$)

$$\begin{aligned} & \eta \int_{\Omega} \rho \frac{((|\phi_j|^2 - 1)^+)^2}{2\tau} dx + \eta \int_{\Omega} \rho \frac{(|\phi_j|^2 - 1)^+}{2\tau} dx + \int_{\Omega} \rho |b(U_j)(\phi_j)|^2 (|\phi_j|^2 - 1)^+ dx \\ & \quad + \int_{\Omega^+} \frac{\rho}{2\kappa^2} |\nabla |\phi_j|^2|^2 dx + \int_{\Omega} \rho |\phi_j|^2 ((|\phi_j|^2 - 1)^+)^2 dx \leq 0. \end{aligned} \quad (2.4)$$

Therefore, we have that $|\phi_j| \leq 1$ a.e. in Ω .

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Prove now some estimates. First, we take $(q, \nu) = (\overline{\phi_j}, U_j)$ in (2.2), by using the uniform boundness of the sequence $(\phi_j)_{j \geq 1}$ and the discrete Gronwall's formula we obtain the following estimate

$$\begin{aligned} & \max_{j=1,n} \left(\|\phi_j\|_{L^2}^2 + \|U_j\|_{L^2}^2 \right) + \tau \sum_{j=1,N} \left(\|\nabla \phi_j\|_{L^2}^2 + \|\operatorname{div}(U_j)\|_{L^2}^2 + \|\operatorname{curl}(U_j)\|_{L^2}^2 \right) \\ & \leq C \left(\|\phi_0\|_{L^2}^2 + \|U_0\|_{L^2}^2 + \|H\|_{L^2(\mathcal{Q})}^2 \right). \end{aligned} \quad (2.5)$$

By taking now $(q, \nu) = (1/\tau)(\overline{\phi_j - \phi_{j-1}}, U_j - U_{j-1})$ in (2.2), by using the uniform boundness of the sequence $(\phi_j)_{j \geq 1}$ and the discrete Gronwall's formula, we have

$$\begin{aligned} & \tau \sum_{j=1,N} \left(\left\| \frac{1}{\tau} (\phi_j - \phi_{j-1}) \right\|_{L^2}^2 + \left\| \frac{1}{\tau} (u_j - u_{j-1}) \right\|_{L^2}^2 \right) \\ & \quad + \sum_{j=1,N} \left(\|\nabla \phi_j\|_{L^2}^2 + \|\operatorname{div}(U_j)\|_{L^2}^2 + \|\operatorname{curl}(U_j)\|_{L^2}^2 \right) \\ & \leq C \left(\|\phi_0\|_{L^2}^2 + \|U_0\|_{L^2}^2 + \|H\|_{L^2(\mathcal{Q})}^2 \right). \end{aligned} \quad (2.6)$$

By using (2.2), Green's formula, the estimates (2.5), (2.6), the uniform boundness of $(\phi_j)_{j \geq 1}$, and Lemma 1.4 we obtain the following estimation:

$$\begin{aligned} & \sum_{j=1,N} \left(\left\| \frac{1}{\tau} (\phi_j - \phi_{j-1}) \right\|_{(H^1(\Omega))'}^2 + \left\| \frac{1}{\tau} (U_j - U_{j-1}) \right\|_{V_n'}^2 \right) \\ & \leq C \left(\|\phi_0\|_{L^2}^2 + \|U_0\|_{L^2}^2 + \|H\|_{L^2(\mathcal{Q})}^2 \right). \end{aligned} \quad (2.7)$$

The proof of theorem can be completed by standard convergence arguments (see, e.g., [32]), by taking advantage of the above estimates and by using the continuous mapping from $\mathcal{H}^{1/2+s}(\Omega)$ into $\mathcal{L}^2(\Gamma)$, $0 < s < 1/2$, see, for example, [34, 35] (to pass to the limit in term $\int_{\Gamma} \rho \phi_j q d\Gamma$). \square

Remark 2.2. Throughout the paper, we suppose that the hypotheses of Theorem 2.1 are satisfied, to ensure that the solution of problem (1.1), is in $({}^{\mathcal{Q}}W \cap \mathcal{L}^{\infty}(\mathcal{Q})) \times \mathbf{W}_n$.

3. Studying the perturbation problem

In the following, the solution (ϕ, U) of problem (1.1) will be treated as the target function. We are then interested in the robust regulation of the deviation of the problem from the desired target (ϕ, U) . We analyze the full nonlinear equation which models large perturbations (φ, u) to the target (ϕ, U) . Hence we consider the following system (for (ϕ, U))

given satisfying the regularity of Theorem 2.1):

$$\begin{aligned} \eta\rho \frac{\partial\varphi}{\partial t} - i\eta\kappa \operatorname{div}(\rho(u+U))\varphi - i\eta\kappa \operatorname{div}(\rho u)\phi + b(u+U) \cdot (\rho b(u+U))(\varphi) \\ + (b(u+U) \cdot (\rho b(u+U)) - b(U) \cdot (\rho b(U)))(\phi) \\ + \rho(G(\phi+\varphi) - G(\phi)) = 0 \quad \text{in } \mathfrak{D}, \end{aligned} \quad (3.1a)$$

$$\begin{aligned} \rho \frac{\partial v}{\partial t} + \operatorname{curl}(\rho \operatorname{curl}(u)) - \nabla(\operatorname{div}(\rho u)) \\ + \rho\mathcal{R}(b(u+U)(\varphi+\phi)(\overline{\varphi}+\overline{\phi}) - b(U)(\phi)\overline{\phi}) = \operatorname{curl}(\rho h) \quad \text{in } \mathfrak{D}, \end{aligned}$$

subjected to the Robin-type boundary conditions

$$\frac{1}{\kappa^2} \frac{\partial\varphi}{\partial n} = \mu\varphi, \quad u \cdot \mathbf{n} = 0, \quad \operatorname{curl}(u) = h \quad \text{in } \Sigma, \quad (3.1b)$$

and the initial conditions

$$\varphi(0) = \varphi_0, \quad u(0) = u_0 \quad \text{in } \Omega. \quad (3.1c)$$

If we set $F(\varphi) = G(\varphi+\phi) - G(\phi)$, $B(u) = b(U+u)$ then (3.1) is reduced to

$$\begin{aligned} \eta\rho \frac{\partial\varphi}{\partial t} - i\eta\kappa \operatorname{div}(\rho(u+U))\varphi - i\eta\kappa \operatorname{div}(\rho u)\phi \\ + B(u) \cdot (\rho B(u))(\varphi+\phi) + \rho F(\varphi) = B(0) \cdot (\rho B(0))(\phi) \quad \text{in } \mathfrak{D}, \end{aligned} \quad (3.2a)$$

$$\begin{aligned} \rho \frac{\partial u}{\partial t} + \operatorname{curl}(\rho \operatorname{curl}(u)) - \nabla(\operatorname{div}(\rho u)) + \rho\mathcal{R}(B(u)(\varphi+\phi)(\overline{\varphi}+\overline{\phi})) \\ = \rho\mathcal{R}(B(0)(\phi)\overline{\phi}) + \operatorname{curl}(\rho h) \quad \text{in } \mathfrak{D}, \end{aligned}$$

subjected to the Robin-type boundary conditions

$$\frac{1}{\kappa^2} \frac{\partial\varphi}{\partial n} = \mu\varphi, \quad u \cdot \mathbf{n} = 0, \quad \operatorname{curl}(u) = h \quad \text{in } \Sigma, \quad (3.2b)$$

and the initial conditions

$$\varphi(0) = \varphi_0, \quad u(0) = u_0 \quad \text{in } \Omega. \quad (3.2c)$$

Now we give the weak formulation associated to problem (3.2).

Multiplying the first part of (3.2) by $q \in \mathcal{H}^1(\Omega)$ and the second part by $v \in V_n$ and integrating over Ω gives (according to the third part of (3.2)) the weak formulation

(a.e. $t \in (0, T)$)

$$\begin{aligned}
& \eta \int_{\Omega} \rho \frac{\partial \varphi}{\partial t} q \, dx - i\eta\kappa \int_{\Omega} \operatorname{div}(\rho(u+U))\varphi q \, dx - i\eta\kappa \int_{\Omega} \operatorname{div}(\rho u)\varphi q \, dx - i\mu \int_{\Gamma} \rho\varphi q \, d\Gamma \\
& \quad + \int_{\Omega} \rho B(u)(\varphi + \bar{\varphi})\bar{B}(u)(q) \, dx + \int_{\Omega} \rho F(\varphi)q \, dx = \int_{\Omega} \rho B(0)(\varphi)\bar{B}(0)(q) \, dx, \\
& \int_{\Omega} \rho \frac{\partial v}{\partial t} \, dx + \int_{\Omega} \rho \operatorname{curl}(u) \operatorname{curl}(v) \, dx + \int_{\Omega} \operatorname{div}(\rho u) \operatorname{div}(v) \, dx \\
& \quad + \int_{\Omega} \rho \mathcal{R}((\bar{\varphi} + \varphi)B(u)(\varphi + \bar{\varphi}))v \, dx = \int_{\Omega} \rho \mathcal{R}(\bar{\varphi}B(0)(\varphi))v \, dx + \int_{\Omega} \rho h \operatorname{curl}(v) \, dx, \\
& (\varphi(0), u(0)) = (\varphi_0, u_0).
\end{aligned} \tag{3.3}$$

3.1. The existence and the stability results. Now we show the existence of the solution to the problem (3.3).

THEOREM 3.1. *For any $(\varphi_0, u_0) \in \mathcal{L}^2_{\infty}(\Omega) \times \mathbf{L}^2(\Omega)$ and $h \in L^2(\mathcal{Q})$, there exists $(\varphi, u) \in (\mathcal{W} \cap \mathcal{L}^{\infty}(\mathcal{Q})) \times \mathbf{W}_n$ solution of (3.3). Moreover, the following estimation exists:*

$$\|\varphi\|_{\mathcal{W}}^2 + \|u\|_{\mathbf{W}_n}^2 \leq C\left(\|h\|_{L^2(\mathcal{Q})}^2 + \|\varphi_0\|_{L^2(\Omega)}^2 + \|u_0\|_{L^2(\Omega)}^2\right). \tag{3.4}$$

Proof. The proof of this result is a consequence of the result of Theorem 2.1, so we omit the tedious details. \square

Next we will establish a stability result which gives us uniqueness of solution of (3.2) as a corollary.

THEOREM 3.2. *Let $(u_{01}, \varphi_{01}, h_1)$ and $(u_{02}, \varphi_{02}, h_2)$ be functions of $\mathbf{L}^2(\Omega) \times \mathcal{L}^2_{\infty}(\Omega) \times L^2(\mathcal{Q})$. If $(u_1, \varphi_1) \in \mathbf{W}_n \times (\mathcal{W} \cap \mathcal{L}^{\infty}(\mathcal{Q}))$ (resp., $(u_2, \varphi_2) \in \mathbf{W}_n \times (\mathcal{W} \cap \mathcal{L}^{\infty}(\mathcal{Q}))$) is solution of (3.3) with $(\varphi_{01}, u_{01}, h_1)$ (resp., $(\varphi_{02}, u_{02}, h_2)$) the given data, then the following estimation exists:*

$$\|\varphi\|_{\mathcal{W}}^2 + \|u\|_{\mathbf{W}_n}^2 \leq C\left(\|\varphi_0\|_{L^2}^2 + \|u_0\|_{L^2}^2 + \|h\|_{L^2(\mathcal{Q})}^2\right), \tag{3.5}$$

where $\varphi = \varphi_1 - \varphi_2$, $u = u_1 - u_2$, $\varphi_0 = \varphi_{01} - \varphi_{02}$, $u_0 = u_{01} - u_{02}$, and $h = h_1 - h_2$.

Proof. Let $(\varphi_i, u_i, \varphi_{0i}, u_{0i}, h_i)_{i=1,2}$ be two solutions of (3.3) with $(\varphi_{0i}, u_{0i})_{i=1,2}$ the initial data respectively. We denote by $\varphi = \varphi_1 - \varphi_2$, $u = u_1 - u_2$, $\varphi_0 = \varphi_{01} - \varphi_{02}$, $u_0 = u_{01} - u_{02}$, and

$h = h_1 - h_2$. Then the couple (φ, u) is solution of

$$\begin{aligned}
& \eta \int_{\Omega} \rho \frac{\partial \varphi}{\partial t} q \, dx - i\eta \kappa \left(\int_{\Omega} \operatorname{div}(\rho U_2) \varphi q \, dx + \int_{\Omega} \operatorname{div}(\rho u) \phi_1 q \, dx \right) - i\mu \int_{\Gamma} \rho \varphi q \, d\Gamma \\
& \quad + \int_{\Omega} \rho B(u_1)(\varphi) \overline{B}(u_1)(q) \, dx + \int_{\Omega} \rho u \phi_2 \overline{B}(u_2)(q) \, dx \\
& \quad + \int_{\Omega} \rho B(u_1)(\phi_2) u q \, dx + \int_{\Omega} \rho (F(\varphi_1) - F(\varphi_2)) q \, dx = 0, \\
& \int_{\Omega} \rho \frac{\partial u}{\partial t} v \, dx + \int_{\Omega} \rho \operatorname{curl}(u) \operatorname{curl}(v) \, dx + \int_{\Omega} \operatorname{div}(\rho u) \operatorname{div}(v) \, dx + \int_{\Omega} \rho \mathcal{R}(\overline{\varphi} B(u_1)(\phi_1)) v \, dx \\
& \quad + \int_{\Omega} \rho \mathcal{R}(\overline{\phi}_2 B(u_1)(\varphi)) v \, dx + \int_{\Omega} u |\phi_2|^2 v \, dx = \int_{\Omega} \rho h \operatorname{curl}(v) \, dx, \\
& (\varphi(0), u(0)) = (\varphi_0, u_0),
\end{aligned} \tag{3.6}$$

where $U_i = u_i + U$ and $\phi_i = \varphi_i + \phi$, for $i = 1, 2$ (according to the regularity of (φ_i, u_i) and (ϕ, U) , the couple (ϕ_i, U_i) is in $(\mathcal{W} \cap \mathcal{L}^\infty(\mathcal{Q})) \times \mathbf{W}_n$).

By choosing $(q, v) = (\overline{\varphi}, u)$ in (3.6) and taking the real part of the first part of the system, we have (according to Lemma 1.5 and to the expression of F)

$$\begin{aligned}
& \eta \int_{\Omega} \rho \frac{\partial |\varphi|^2}{2 \partial t} \, dx + \eta \kappa \int_{\Omega} \operatorname{div}(\rho u) \mathcal{I}(\phi_1 \overline{\varphi}) \, dx + \int_{\Omega} \rho B(u_1)(\varphi) \overline{B}(u_1)(\varphi) \, dx \\
& \quad + \int_{\Omega} \rho u \mathcal{R}(\overline{\phi_2 B}(u_2)(\varphi)) \, dx + \int_{\Omega} \rho u \mathcal{R}(B(u_1)(\phi_2) \overline{\varphi}) \, dx \\
& \quad + \int_{\Omega} \rho (|\phi_1|^2 + |\phi_2|^2 - 1) |\varphi|^2 \, dx + \int_{\Omega} \rho \mathcal{R}(\phi_1 \phi_2 (\overline{\varphi})^2) \, dx = 0, \\
& \int_{\Omega} \rho \frac{\partial |u|^2}{2 \partial t} v \, dx + \int_{\Omega} \rho |\operatorname{curl}(u)|^2 \, dx + \int_{\Omega} \operatorname{div}(\rho u) \operatorname{div}(u) \, dx + \int_{\Omega} \rho \mathcal{R}(\overline{\varphi} B(u_1)(\phi_1)) u \, dx \\
& \quad + \int_{\Omega} \rho \mathcal{R}(\overline{\phi}_2 B(u_1)(\varphi)) u + \int_{\Omega} |u|^2 |\phi_2|^2 \, dx = \int_{\Omega} \rho h \operatorname{curl}(u) \, dx, \\
& (\varphi(0), u(0)) = (\varphi_0, u_0)
\end{aligned} \tag{3.7}$$

and then (according to hypothesis (1.3), Lemma 1.6 and to the regularity of ϕ , $(\phi_i)_{i=1,2}$, that is, in $\mathcal{L}^\infty(\mathcal{Q})$)

$$\begin{aligned}
& \eta \rho_0 \frac{d \|\varphi\|_{L^2}^2}{2 dt} + \frac{\rho_0}{\kappa^2} \|\nabla \varphi\|_{L^2}^2 \\
& \leq c_1 \int_{\Omega} |\operatorname{div}(\rho u)| |\varphi| \, dx + c_2 \int_{\Omega} |U_1| |\varphi| |\nabla \varphi| \, dx + c_3 \int_{\Omega} (1 + |U_1|^2) |\varphi|^2 \, dx \\
& \quad + c_4 \int_{\Omega} (|\nabla \phi_2| + |U_2| + |U_1|) |\varphi| |u| \, dx + c_5 \int_{\Omega} |\nabla \varphi| |u| \, dx,
\end{aligned}$$

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$$\begin{aligned}
& \rho_0 \frac{d\|u\|_{L^2}^2}{2dt} + \rho_0 \left(\|\operatorname{curl}(u)\|_{L^2}^2 + \|\operatorname{div}(u)\|_{L^2}^2 \right) \\
& \leq c_6 \left(\int_{\Omega} |u| |\operatorname{div}(u)| dx + \int_{\Omega} |h| |\operatorname{curl}(u)| dx \right) + c_7 \int_{\Omega} |u| |\varphi| (|\nabla \phi_1| + |U_1|) dx \\
& \quad + c_8 \int_{\Omega} |u| |\nabla \varphi| dx + c_9 \int_{\Omega} |u|^2 dx, \\
& (\varphi(0), u(0)) = (\varphi_0, u_0).
\end{aligned} \tag{3.8}$$

According to the regularity of U_i , ϕ_i , and by using Young's inequality, we have then for all $\delta > 0$ and $\gamma > 0$ (since $H^1 \subset L^4$),

$$\begin{aligned}
& \eta \rho_0 \frac{d\|\varphi\|_{L^2}^2}{2dt} + \frac{\rho_0}{\kappa^2} \|\nabla \varphi\|_{L^2}^2 \\
& \leq \delta \left(\|\operatorname{curl}(u)\|_{L^2}^2 + \|\operatorname{div}(u)\|_{L^2}^2 \right) + \frac{\rho_0}{4\kappa^2} \|\nabla \varphi\|_{L^2}^2 + c_{10} \|U_1\|_{L^4} \|\varphi\|_{L^4} \|\nabla \varphi\|_{L^2} \\
& \quad + c_{11} \left(\|U_1\|_{L^4}^2 + \|U_2\|_{L^4}^2 \right) \|\varphi\|_{L^4}^2 + c_{12} \|u\|_{L^4} \|\varphi\|_{L^4} \|\nabla \phi_2\|_{L^2} + c_{13} \left(\|\varphi\|_{L^2}^2 + \|u\|_{L^2}^2 \right), \\
& \rho_0 \frac{d\|u\|_{L^2}^2}{2dt} + \rho_0 \left(\|\operatorname{curl}(u)\|_{L^2}^2 + \|\operatorname{div}(u)\|_{L^2}^2 \right) \\
& \leq \frac{\rho_0}{4} \left(\|\operatorname{curl}(u)\|_{L^2}^2 + \|\operatorname{div}(u)\|_{L^2}^2 \right) + \gamma \|\nabla \varphi\|_{L^2}^2 + c_{14} \|h\|_{L^2}^2 \\
& \quad + c_{15} \|u\|_{L^4} \|\varphi\|_{L^4} \left(\|\nabla \phi_1\|_{L^2} + \|U_1\|_{L^2} \right) + c_{16} \left(\|\varphi\|_{L^2}^2 + \|u\|_{L^2}^2 \right), \\
& (\varphi(0), u(0)) = (\varphi_0, u_0).
\end{aligned} \tag{3.9}$$

By using Gagliardo-Nirenberg inequalities (Lemma 1.4), we have

$$\begin{aligned}
& \frac{d}{dt} \left(\|\varphi\|_{L^2}^2 + \|u\|_{L^2}^2 \right) + \|\nabla \varphi\|_{L^2}^2 + \left(\|\operatorname{curl}(u)\|_{L^2}^2 + \|\operatorname{div}(u)\|_{L^2}^2 \right) \\
& \leq c_{17} \|h\|_{L^2}^2 + c_{18} \left(1 + \|\nabla \phi_1\|_{L^2}^2 + \|\nabla \phi_2\|_{L^2}^2 \right) \left(\|\varphi\|_{L^2}^2 + \|u\|_{L^2}^2 \right) \\
& \quad + c_{19} \left(\|U_1\|_{L^4}^4 + \|U_2\|_{L^4}^4 + \|U_1\|_{L^4}^2 + \|U_2\|_{L^4}^2 \right) \|\varphi\|_{L^2}^2, \\
& (\varphi(0), u(0)) = (\varphi_0, u_0).
\end{aligned} \tag{3.10}$$

Since U_i , $i = 1, 2$, are in $L^\infty(0, T, L^2(\Omega))$, by using again Gagliardo-Nirenberg inequalities, we can deduce that (by integrating over $(0, t)$ for $t \in (0, T)$)

$$\begin{aligned}
& \|\varphi\|_{L^2}^2 + \|u\|_{L^2}^2 + \int_0^t \|\nabla \varphi\|_{L^2}^2 ds + \int_0^t \left(\|\operatorname{curl}(u)\|_{L^2}^2 + \|\operatorname{div}(u)\|_{L^2}^2 \right) ds \\
& \leq c_{20} \int_0^t \left(1 + \|\phi_1\|_{H^1}^2 + \|\phi_2\|_{H^1}^2 + \|U_1\|_{H^1}^2 + \|U_2\|_{H^1}^2 \right) \left(\|\varphi\|_{L^2}^2 + \|u\|_{L^2}^2 \right) ds \\
& \quad + c_{21} \left(\|\varphi_0\|_{L^2}^2 + \|u_0\|_{L^2}^2 + \|h\|_{L^2(\mathcal{Q})}^2 \right).
\end{aligned} \tag{3.11}$$

Gronwall's formula implies

$$\|\varphi\|_{\mathcal{E}}^2 + \|\mathbf{u}\|_{\mathbf{E}_n}^2 \leq C \left(\|\varphi_0\|_{L^2}^2 + \|\mathbf{u}_0\|_{L^2}^2 + \|\mathbf{h}\|_{L^2(\mathcal{Q})}^2 \right). \quad (3.12)$$

By using (3.6), Green's formula, the estimate (3.12), and Lemma 1.4, we can deduce the result of the theorem. \square

According to Theorem 3.2, we have the following result.

COROLLARY 3.3. *The solution of problem (3.3) is unique.*

We are now going to study the differentiability of the operator solution of problem (3.3).

4. Differentiability of the operator solution

Before proceeding with investigation of the Fréchet differentiability of the function $\mathcal{F} : (\varphi_0, \mathbf{u}_0, \mathbf{h}) \rightarrow (\varphi, \mathbf{u})$, which maps the source term $(\varphi_0, \mathbf{u}_0, \mathbf{h}) \in \mathcal{L}_\infty^2(\Omega) \times \mathbf{L}^2(\Omega) \times L^2(\mathcal{Q})$ of problem (3.3) into the corresponding solution $(\varphi, \mathbf{u}) \in \mathcal{E} \times \mathbf{E}_n$, we study the following problem (\mathcal{P}_I) : find $(\psi, \mathbf{w}) \in \mathcal{E} \times \mathbf{E}_n$ such that $(\forall (q, \mathbf{v}) \in \mathcal{H}^1(\Omega) \times \mathcal{V}_n$ and a.e. $t \in (0, T)$):

$$\begin{aligned} & \eta \int_{\Omega} \rho \frac{\partial \psi}{\partial t} q \, dx - i\eta\kappa \int_{\Omega} \operatorname{div}(\rho U_1) \psi q \, dx - i\eta\kappa \int_{\Omega} \operatorname{div}(\rho \mathbf{w}) \phi_1 q \, dx - i\mu \int_{\Gamma} \rho \psi q \, d\Gamma \\ & \quad + \int_{\Omega} \rho (B(u)(\psi) + \phi_1 \mathbf{w}) \bar{B}(u)(q) \, dx + \int_{\Omega} \rho B(u)(\phi_1) \mathbf{w} q \, dx \\ & \quad + \int_{\Omega} \rho ((2|\phi_1|^2 - 1)\psi + \phi_1^2 \bar{\psi}) q \, dx = 0, \\ & \int_{\Omega} \rho \frac{\partial \mathbf{w}}{\partial t} \mathbf{v} \, dx + \int_{\Omega} \rho \operatorname{curl}(\mathbf{w}) \operatorname{curl}(\mathbf{v}) \, dx + \int_{\Omega} \operatorname{div}(\rho \mathbf{w}) \operatorname{div}(\mathbf{v}) \, dx + \int_{\Omega} \rho \mathcal{R}(\bar{\psi} B(u)(\phi_1)) \mathbf{v} \, dx \\ & \quad + \int_{\Omega} \rho (\mathcal{R}(\bar{\phi}_1 B(u)(\psi)) + \mathbf{w} |\phi_1|^2) \mathbf{v} \, dx = \int_{\Omega} \rho k \operatorname{curl}(\mathbf{v}) \, dx, \\ & (\psi(0), \mathbf{w}(0)) = (\psi_0, \mathbf{w}_0), \end{aligned} \quad (4.1)$$

where $U_1 = U + \mathbf{u}$ and $\phi_1 = \phi + \varphi$.

Remark 4.1. The problem (P_I) is the weak formulation of the following problem:

$$\begin{aligned} & \eta \rho \frac{\partial \psi}{\partial t} - i\eta\kappa (\operatorname{div}(\rho U_1) \psi + \operatorname{div}(\rho \mathbf{w}) \phi_1) + B(u) \cdot (\rho (B(u)(\psi) + \phi_1 \mathbf{w})) + \rho B(u)(\phi_1) \mathbf{w} \\ & \quad + \rho ((2|\phi_1|^2 - 1)\psi + \phi_1^2 \bar{\psi}) = 0 \quad \text{in } \mathcal{Q}, \\ & \rho \frac{\partial \mathbf{w}}{\partial t} + \operatorname{curl}(\rho \operatorname{curl}(\mathbf{w})) + \nabla(\operatorname{div}(\rho \mathbf{w})) + \rho \mathcal{R}(\bar{\psi} B(u)(\phi_1)) \\ & \quad + \rho (\mathcal{R}(\bar{\phi}_1 B(u)(\psi)) + \mathbf{w} |\phi_1|^2) = \operatorname{curl}(\rho k) \quad \text{in } \mathcal{Q}, \end{aligned} \quad (4.2a)$$

subjected to the Robin-type boundary conditions

$$\frac{1}{\kappa^2} \frac{\partial \psi}{\partial n} = \mu \psi, \quad w \cdot \mathbf{n} = 0, \quad \text{curl}(w) = k \quad \text{in } \Sigma, \quad (4.2b)$$

and the initial conditions

$$(\psi(0), w(0)) = (\psi_0, w_0) \quad \text{in } \Omega. \quad (4.2c)$$

THEOREM 4.2. *If (u, ϕ) and (U, Φ) are in $\mathbf{W}_n \times (\mathcal{W} \cap \mathcal{L}^\infty(\mathcal{Q}))$, then*

(i) *for any $(\psi_0, w_0, k) \in \mathcal{L}^2_\infty(\Omega) \times \mathbf{L}^2(\Omega) \times L^2(\mathcal{Q})$, there exists a unique couple of functions $(\psi, w) \in \mathcal{E} \times \mathbf{E}_n$, solution of problem (\mathcal{P}_I) , such that*

$$\|\psi\|_{\mathcal{E}}^2 + \|w\|_{\mathbf{E}_n}^2 \leq C_e \left(\|\psi_0\|_{L^2}^2 + \|w_0\|_{L^2}^2 + \|k\|_{L^2(\mathcal{Q})}^2 \right), \quad (4.3)$$

(ii) *let (ψ_{0i}, w_{0i}, k_i) , $i = 1, 2$, be two couples of $\in \mathcal{L}^2_\infty(\Omega) \times \mathbf{L}^2(\Omega) \times L^2(\mathcal{Q})$. If (ψ_i, w_i) is the solution of (\mathcal{P}_I) , where the initial condition is (ψ_{0i}, w_{0i}) , $i = 1, 2$, then*

$$\|\psi_1 - \psi_2\|_{\mathcal{E}}^2 + \|w_1 - w_2\|_{\mathbf{E}_n}^2 \leq C_e \left(\|\psi_{01} - \psi_{02}\|_{L^2}^2 + \|w_{01} - w_{02}\|_{L^2}^2 + \|k_1 - k_2\|_{L^2(\mathcal{Q})}^2 \right). \quad (4.4)$$

Proof. (i) The existence of the solution of (\mathcal{P}_I) is obtained in the same way as to prove the Theorem 2.1 and by using the regularity of (U_1, ϕ_1) . The uniqueness is a consequence of the estimate (4.3) (since the problem (\mathcal{P}_I) is linear).

To prove the estimate (4.3), we put $(q, v) = (\bar{\psi}, w)$ in (P_I) and we obtain

$$\begin{aligned} & \eta \int_{\Omega} \rho \frac{\partial \psi}{\partial t} \bar{\psi} dx - i\eta\kappa \int_{\Omega} \text{div}(\rho U_1) |\psi|^2 dx - i\eta\kappa \int_{\Omega} \text{div}(\rho w) \phi_1 \bar{\psi} dx \\ & - i\mu \int_{\Gamma} \rho |\psi|^2 d\Gamma + \int_{\Omega} \rho |B(u)(\psi)|^2 dx + \int_{\Omega} \rho (B(u)(\phi_1) \bar{\psi} + \phi_1 \overline{B(u)(\bar{\psi})}) w dx \\ & + \int_{\Omega} \rho ((2|\phi_1|^2 - 1) |\psi|^2 + \phi_1^2 (\bar{\psi})^2) dx = 0, \\ & \int_{\Omega} \rho \frac{\partial |w|^2}{2\partial t} dx + \int_{\Omega} \rho |\text{curl}(w)|^2 dx + \int_{\Omega} \text{div}(\rho w) \text{div}(w) dx + \int_{\Omega} \rho \mathcal{R}(\bar{\psi} B(u)(\phi_1)) w dx \\ & + \int_{\Omega} \rho (\mathcal{R}(\overline{\phi_1} B(u)(\psi)) + w |\phi_1|^2) w dx = \int_{\Omega} \rho k \text{curl}(w) dx, \\ & (\psi(0), w(0)) = (\psi_0, w_0). \end{aligned} \quad (4.5)$$

By taking the real part of the first part of previous system, by using the same technique as to prove the stability result and the regularity of (U_1, ϕ_1) (so we omit the details), we obtain the estimate given in (i).

(ii) Since the problem (\mathcal{P}_I) is linear, then the result (ii) is a consequence of the estimate (4.3). \square

We are now going to study the Fréchet differentiability of \mathcal{F} .

THEOREM 4.3. (i) Let $\varphi_0, \varphi_0 + \beta \in \mathcal{L}_\infty^2(\Omega)$ and $(u_0, h) \in \mathbf{L}^2(\Omega) \times \mathbf{L}^2(\mathcal{Q})$ with $\mathcal{F}(\varphi_0, u_0, h)$ and $\mathcal{F}(\varphi_0 + \beta, u_0, h)$ being the corresponding solutions of (3.3). Then

$$\|\mathcal{F}(\varphi_0 + \beta, u_0, h) - \mathcal{F}(\varphi_0, u_0, h) - \mathcal{F}'_\varphi(\varphi_0, u_0, h)\beta\|_{\mathcal{E} \times \mathbf{E}_n} \leq C\|\beta\|_{\mathcal{L}^2}^{3/2}, \quad (4.6)$$

where $\mathcal{F}'_\varphi(\varphi_0, u_0, h) : \mathcal{L}_\infty^2(\Omega) \rightarrow \mathcal{E} \times \mathbf{E}_n$ is a linear operator, and $(\psi, w) = \mathcal{F}'_\varphi(\varphi_0, u_0, h)\beta$ is the solution of the problem (\mathcal{P}_I) satisfying $(\psi, w)(t=0) = (\beta, 0)$ and $k=0$ (denote this problem by (\mathcal{P}_{FP})).

Moreover, for all $X_h^{(i)} = (\varphi_{0i}, u_{0i}, h_i) \in \mathcal{L}_\infty^2(\Omega) \times \mathbf{L}^2(\Omega) \times \mathbf{L}^2(\mathcal{Q})$, $i=1,2$, the following estimate exists: $\|\mathcal{F}'_\varphi(X_h^{(1)})\beta - \mathcal{F}'_\varphi(X_h^{(2)})\beta\|_{\mathcal{E} \times \mathbf{E}_n}^2 \leq C_e(\|\beta\|_{\mathcal{L}^2} \|X_h\|_{\mathcal{L}^2 \times \mathbf{L}^2 \times \mathbf{L}^2(\mathcal{Q})}^2 + \|\beta\|_{\mathcal{L}^2}^2 \|X_h\|_{\mathcal{L}^2 \times \mathbf{L}^2 \times \mathbf{L}^2(\mathcal{Q})})$, where $\varphi_0 = \varphi_{01} - \varphi_{02}$, $u_0 = u_{01} - u_{02}$, $h = h_1 - h_2$ and $X_h = X_h^{(1)} - X_h^{(2)}$.

(ii) Let $u_0, u_0 + z \in \mathbf{L}^2(\Omega)$ and $(\varphi_0, h) \in \mathcal{L}_\infty^2(\Omega) \times \mathbf{L}^2(\mathcal{Q})$ with $\mathcal{F}(\varphi_0, u_0, h)$ and $\mathcal{F}(\varphi_0, u_0 + z, h)$ being the corresponding solutions of (3.3). Then

$$\|\mathcal{F}(\varphi_0, u_0 + z, h) - \mathcal{F}(\varphi_0, u_0, h) - \mathcal{F}'_u(\varphi_0, u_0, h)z\|_{\mathcal{E} \times \mathbf{E}_n} \leq C\|z\|_{\mathbf{L}^2}^{3/2}, \quad (4.7)$$

where $\mathcal{F}'_u(\varphi_0, u_0, h) : \mathbf{L}^2(\Omega) \rightarrow \mathcal{E} \times \mathbf{E}_n$ is a linear operator, and $(\psi, w) = \mathcal{F}'_u(\varphi_0, u_0, h)z$ is the solution of the problem (\mathcal{P}_I) satisfying $(\psi, w)(t=0) = (0, z)$ and $k=0$ (denote this problem by (\mathcal{P}_{Fu})).

Moreover, for all $X_h^{(i)} = (\varphi_{0i}, u_{0i}, h_i) \in \mathcal{L}_\infty^2(\Omega) \times \mathbf{L}^2(\Omega) \times \mathbf{L}^2(\mathcal{Q})$, $i=1,2$, the following estimate exists:

$$\begin{aligned} & \left\| \mathcal{F}'_u(X_h^{(1)})z - \mathcal{F}'_u(X_h^{(2)})z \right\|_{\mathcal{E} \times \mathbf{E}_n}^2 \\ & \leq C_e \left(\|z\|_{\mathbf{L}^2} \|X_h\|_{\mathcal{L}^2 \times \mathbf{L}^2 \times \mathbf{L}^2(\mathcal{Q})}^2 + \|z\|_{\mathbf{L}^2}^2 \|X_h\|_{\mathcal{L}^2 \times \mathbf{L}^2 \times \mathbf{L}^2(\mathcal{Q})} \right), \end{aligned} \quad (4.8)$$

where $\varphi_0 = \varphi_{01} - \varphi_{02}$, $u_0 = u_{01} - u_{02}$, $h = h_1 - h_2$ and $X_h = X_h^{(1)} - X_h^{(2)}$.

(iii) Let $h, h + f \in \mathbf{L}^2(\mathcal{Q})$ and $(\varphi_0, u_0) \in \mathcal{L}_\infty^2(\Omega) \times \mathbf{L}^2(\Omega)$ with $\mathcal{F}(\varphi_0, u_0, h)$ and $\mathcal{F}(\varphi_0, u_0, h + f)$ being the corresponding solutions of (3.3). Then

$$\|\mathcal{F}(\varphi_0, u_0, h + f) - \mathcal{F}(\varphi_0, u_0, h) - \mathcal{F}'_h(\varphi_0, u_0, h)f\|_{\mathcal{E} \times \mathbf{E}_n} \leq C\|f\|_{\mathbf{L}^2}^{3/2}, \quad (4.9)$$

where $\mathcal{F}'_h(\varphi_0, u_0, h) : \mathbf{L}^2(\mathcal{Q}) \rightarrow \mathcal{E} \times \mathbf{E}_n$ is a linear operator, and $(\psi, w) = \mathcal{F}'_h(\varphi_0, u_0, h)f$ is the solution of the problem (\mathcal{P}_I) satisfying $(\psi, w)(t=0) = (0, 0)$ and $k=f$ (denote this problem by (\mathcal{P}_{Fh})).

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Moreover, for all $X_h^{(i)} = (\varphi_{0i}, u_{0i}, h_i) \in \mathcal{L}_\infty^2(\Omega) \times \mathbf{L}^2(\Omega) \times \mathbf{L}^2(\mathcal{Q})$, $i = 1, 2$, the following estimate exists:

$$\begin{aligned} & \left\| \mathcal{F}'_h(X_h^{(1)})f - \mathcal{F}'_h(X_h^{(2)})f \right\|_{\mathcal{E} \times \mathbf{E}_n}^2 \\ & \leq C_e \left(\|f\|_{L^2(\mathcal{Q})} \|X_h\|_{\mathcal{L}^2 \times \mathbf{L}^2 \times \mathbf{L}^2(\mathcal{Q})}^2 + \|f\|_{L^2(\mathcal{Q})}^2 \|X_h\|_{\mathcal{L}^2 \times \mathbf{L}^2 \times \mathbf{L}^2(\mathcal{Q})} \right), \end{aligned} \quad (4.10)$$

where $\varphi_0 = \varphi_{01} - \varphi_{02}$, $u_0 = u_{01} - u_{02}$, $h = h_1 - h_2$ and $X_h = X_h^{(1)} - X_h^{(2)}$.

Proof. According to Theorem 4.2, the problems (\mathcal{P}_{FP}) , (\mathcal{P}_{Fu}) , and (\mathcal{P}_{Fh}) have a unique solution in $\mathcal{E} \times \mathbf{E}_n$.

(i) Let be $(\varphi, u) = \mathcal{F}(\varphi_0, u_0, h)$ and $(\varphi_\beta, u_\beta) = \mathcal{F}(\varphi_0 + \beta, u_0, h)$. From the stability estimate in Theorem 3.2, we know that

$$\|\varphi_\beta - \varphi\|_{\mathcal{E}}^2 + \|u_\beta - u\|_{\mathbf{E}_n}^2 \leq C \|\beta\|_{\mathcal{L}^2}^2. \quad (4.11)$$

Denote $\phi_\beta = \varphi_\beta - \varphi$, $U_\beta = u_\beta - u$, $U_1 = U + u$, $\phi_1 = \phi + \varphi$, $\varphi^* = \phi_\beta - \psi$, $u^* = U_\beta - w$.

It is easy to see that (φ^*, u^*) satisfies the linear problem (a.e. $t \in (0, T)$)

$$\begin{aligned} & \eta \int_\Omega \rho \frac{\partial \varphi^*}{\partial t} q \, dx \\ & - i\eta \kappa \left(\int_\Omega \operatorname{div}(\rho U_1) \varphi^* q \, dx + \int_\Omega \operatorname{div}(\rho u^*) \phi_1 q \, dx + \int_\Omega \operatorname{div}(\rho U_\beta) \phi_\beta q \, dx \right) \\ & - i\mu \int_\Gamma \rho \varphi^* q \, d\Gamma + \int_\Omega \rho (u^* \phi_1 + B(u)(\varphi^*) + U_\beta \phi_\beta) \overline{B}(u)(q) \, dx \\ & + \int_\Omega \rho U_\beta (U_\beta \phi_1 + B(u)(\phi_\beta) + U_\beta \phi_\beta) q \, dx \\ & + \int_\Omega \rho B(u)(\phi_1) u^* q \, dx + \int_\Omega \rho \left((|\phi_\beta|^2 + 2\Re(\overline{\phi_1} \phi_\beta)) \phi_\beta + \phi_1 |\phi_\beta|^2 \right) q \, dx \\ & + \int_\Omega \rho \left((2|\phi_1|^2 - 1) \varphi^* + (\phi_1)^2 \overline{\varphi^*} \right) q \, dx = 0, \\ & \int_\Omega \rho \frac{\partial u^*}{\partial t} v \, dx + \int_\Omega \rho \operatorname{curl}(u^*) \operatorname{curl}(v) \, dx + \int_\Omega \operatorname{div}(\rho u^*) \operatorname{div}(v) \, dx \\ & + \int_\Omega \rho \mathcal{R}(\overline{\phi_\beta} B(u)(\phi_\beta)) v \, dx \\ & + \int_\Omega \rho U_\beta \left(2\Re(\overline{\phi_1} \phi_\beta) + |\phi_\beta|^2 \right) v \, dx + \int_\Omega \rho \mathcal{R}(\overline{\phi_1} B(u)(\varphi^*) + \overline{\varphi^*} B(u)(\phi_1)) v \, dx \\ & + \int_\Omega \rho u^* |\phi_1|^2 v \, dx = 0, \end{aligned}$$

$$(\varphi^*, u^*)(0) = (0, 0) \quad \text{in } \Omega.$$

(4.12)

By choosing $(q, v) = (\overline{\varphi^*}, u^*)$ in (4.12), we obtain (a.e. $t \in (0, T)$)

$$\begin{aligned}
& \eta \int_{\Omega} \rho \frac{\partial \varphi^*}{\partial t} \overline{\varphi^*} dx - i\eta\kappa \left(\int_{\Omega} \operatorname{div}(\rho U_1) |\varphi^*|^2 dx + \int_{\Omega} \operatorname{div}(\rho u^*) \phi_1 \overline{\varphi^*} dx \right. \\
& \quad \left. + \int_{\Omega} \operatorname{div}(\rho U_{\beta}) \phi_{\beta} \overline{\varphi^*} dx \right) \\
& - i\mu \int_{\Gamma} \rho |\varphi^*|^2 d\Gamma + \int_{\Omega} \rho (u^* \phi_1 + B(u)(\varphi^*) + U_{\beta} \phi_{\beta}) \overline{B(u)(\varphi^*)} dx \\
& + \int_{\Omega} \rho U_{\beta} (U_{\beta} \phi_1 + B(u)(\phi_{\beta}) + U_{\beta} \phi_{\beta}) \overline{\varphi^*} dx + \int_{\Omega} \rho B(u)(\phi_1) u^* \overline{\varphi^*} dx \\
& + \int_{\Omega} \rho \left((|\phi_{\beta}|^2 + 2\Re(\overline{\phi_1} \phi_{\beta})) \phi_{\beta} + \phi_1 |\phi_{\beta}|^2 \right) \overline{\varphi^*} dx \\
& + \int_{\Omega} \rho \left((2|\phi_1|^2 - 1) |\varphi^*|^2 + (\phi_1 \overline{\varphi^*})^2 \right) dx = 0, \\
& \int_{\Omega} \rho \frac{\partial |u^*|^2}{2\partial t} dx + \int_{\Omega} \rho |\operatorname{curl}(u^*)|^2 dx + \int_{\Omega} \operatorname{div}(\rho u^*) \operatorname{div}(u^*) dx + \int_{\Omega} \rho \Re(\overline{\phi_{\beta}} B(u)(\phi_{\beta})) u^* dx \\
& + \int_{\Omega} \rho U_{\beta} (2\Re(\overline{\phi_1} \phi_{\beta}) + |\phi_{\beta}|^2) u^* dx + \int_{\Omega} \rho \Re(\overline{\phi_1} B(u)(\varphi^*) + \overline{\varphi^*} B(u)(\phi_1)) u^* dx \\
& + \int_{\Omega} \rho |u^*|^2 |\phi_1|^2 dx = 0.
\end{aligned} \tag{4.13}$$

Since ϕ_1 and ϕ_{β} are in $\mathcal{L}^{\infty}(\Omega)$ and according to hypothesis (1.3) and Lemma 1.6, we have (by taking the real part of the first part of previous system)

$$\begin{aligned}
& \eta \rho_0 \frac{d\|\varphi^*\|_{L^2}^2}{2dt} + \frac{\rho_0}{\kappa^2} \|\nabla \varphi^*\|_{L^2}^2 \\
& \leq C \left(\int_{\Omega} |\operatorname{div}(\rho u^*)| |\varphi^*| dx + \int_{\Omega} |\operatorname{div}(\rho U_{\beta})| |\phi_{\beta}| |\varphi^*| dx \right. \\
& \quad + \int_{\Omega} (|u^*| + |U_{\beta}| |\phi_{\beta}|) |B(u)(\varphi^*)| dx \\
& \quad + \int_{\Omega} |U_1|^2 |\varphi^*|^2 dx + \int_{\Omega} |U_1| |\varphi^*| |\nabla \varphi^*| dx \\
& \quad + \int_{\Omega} |U_{\beta}| (|U_{\beta}| + |B(u)(\phi_{\beta})|) |\varphi^*| dx \\
& \quad \left. + \int_{\Omega} |B(u)(\phi_1)| |u^*| |\varphi^*| dx + \int_{\Omega} |\phi_{\beta}|^2 |\varphi^*| dx + \int_{\Omega} |\varphi^*|^2 dx \right), \\
& \rho_0 \frac{d\|u^*\|_{L^2}^2}{2dt} + \rho_0 \left(\|\operatorname{curl}(u^*)\|_{L^2}^2 + \|\operatorname{div}(u^*)\|_{L^2}^2 \right) \\
& \leq C \left(\int_{\Omega} |u^*| |\operatorname{div}(u^*)| dx + \int_{\Omega} |u^*|^2 dx + \int_{\Omega} |\phi_{\beta}| |B(u)(\phi_{\beta})| |u^*| dx \right. \\
& \quad \left. + \int_{\Omega} |U_{\beta}| |\phi_{\beta}| |u^*| dx + \int_{\Omega} (|B(u)(\varphi^*)| + |\varphi^*| |B(u)(\phi_1)|) |u^*| dx \right),
\end{aligned} \tag{4.14}$$

and then

$$\begin{aligned}
& \eta\rho_0 \frac{d\|\varphi^*\|_{L^2}^2}{2dt} + \frac{\rho_0}{\kappa^2} \|\nabla\varphi^*\|_{L^2}^2 \\
& \leq C \left(\|U_\beta\|_{H^1} \|\phi_\beta\|_{L^4} \|\varphi^*\|_{L^4} + \|u^*\|_{H^1} \|\varphi^*\|_{L^2} + \|\varphi^*\|_{L^2}^2 + \|U_1\|_{L^4} \|\varphi^*\|_{L^4} \|\nabla\varphi^*\|_{L^2} \right. \\
& \quad + \|U_1\|_{L^4}^2 \|\varphi^*\|_{L^4}^2 \\
& \quad + \left((1 + \|U_1\|_{L^4}) (\|u^*\|_{L^2} + \|U_\beta\|_{L^4}^2 + \|\phi_\beta\|_{L^4}^2) + \|U_\beta\|_{L^4} \|\phi_\beta\|_{H^1} \right) \|\varphi^*\|_{L^4} \\
& \quad + (\|u^*\|_{L^2} + \|U_\beta\|_{L^4} \|\phi_\beta\|_{L^4}) \|\varphi^*\|_{H^1} \\
& \quad \left. + (\|U_1\|_{H^1} + \|\phi_1\|_{H^1}) \|\varphi^*\|_{L^4} \|u^*\|_{L^4} + \|\phi_\beta\|_{L^4}^4 \right), \\
& \rho_0 \frac{d\|u^*\|_{L^2}^2}{2dt} + \rho_0 (\|\operatorname{curl}(u^*)\|_{L^2}^2 + \|\operatorname{div}(u^*)\|_{L^2}^2) \\
& \leq C \left(\|u^*\|_{L^2} \|\operatorname{div}(u^*)\|_{L^2} + \|u^*\|_{L^2}^2 + \|\phi_\beta\|_{L^4} (\|\phi_\beta\|_{H^1} + \|U_1\|_{L^4} \|\phi_\beta\|_{L^4}) \|u^*\|_{L^4} \right. \\
& \quad + \|U_\beta\|_{L^4} \|\phi_\beta\|_{L^4} \|u^*\|_{L^4} + \|u^*\|_{L^2} (\|\varphi^*\|_{H^1} + \|U_1\|_{L^4} \|\varphi^*\|_{L^4}) \\
& \quad \left. + \|u^*\|_{L^4} (\|\phi_1\|_{H^1} + \|U_1\|_{L^4} \|\phi_1\|_{L^4}) \|\varphi^*\|_{L^4} \right). \tag{4.15}
\end{aligned}$$

By using Nirenberg-Gagliardo inequality and Lemma 1.7, we have then

$$\begin{aligned}
& \eta\rho_0 \frac{d\|\varphi^*\|_{L^2}^2}{2dt} + \frac{\rho_0}{\kappa^2} \|\nabla\varphi^*\|_{L^2}^2 \leq \frac{\rho_0}{4} (\|\operatorname{curl}(u^*)\|_{L^2}^2 + \|\operatorname{div}(u^*)\|_{L^2}^2) + \frac{\rho_0}{4\kappa^2} \|\nabla\varphi^*\|_{L^2}^2 \\
& \quad + c_1 (1 + \|U_1\|_{H^1}^2 + \|\phi_1\|_{H^1}^2 + \|U_1\|_{L^4}^4) (\|u^*\|_{L^2}^2 + \|\varphi^*\|_{L^2}^2) \\
& \quad + c_2 (1 + \|U_\beta\|_{H^1}^2 + \|\phi_\beta\|_{H^1}^2) (\|u^*\|_{L^2}^2 + \|\varphi^*\|_{L^2}^2) \\
& \quad + c_3 \|U_\beta\|_{H^1}^2 \|\phi_\beta\|_{L^2} + c_4 \|\phi_\beta\|_{H^1}^2 \|U_\beta\|_{L^2} + c_5 (\|\phi_\beta\|_{L^4}^4 + \|U_\beta\|_{L^4}^4), \\
& \rho_0 \frac{d\|u^*\|_{L^2}^2}{2dt} + \rho_0 (\|\operatorname{curl}(u^*)\|_{L^2}^2 + \|\operatorname{div}(u^*)\|_{L^2}^2) \leq \frac{\rho_0}{4} (\|\operatorname{curl}(u^*)\|_{L^2}^2 + \|\operatorname{div}(u^*)\|_{L^2}^2) \\
& \quad + \frac{\rho_0}{4\kappa^2} \|\nabla\varphi^*\|_{L^2}^2 + c_6 (1 + \|U_1\|_{H^1}^2 + \|\phi_1\|_{H^1}^2 + \|U_1\|_{L^4}^4) (\|u^*\|_{L^2}^2 + \|\varphi^*\|_{L^2}^2) \\
& \quad + c_7 \|\phi_\beta\|_{H^1}^2 \|\phi_\beta\|_{L^2} + c_8 (\|\phi_\beta\|_{L^4}^4 + \|U_\beta\|_{L^4}^4). \tag{4.16}
\end{aligned}$$

By adding the first part and the second part of (4.16) and integrating in time, we obtain

$$\begin{aligned}
& \|\varphi^*\|_{L^2}^2 + \|u^*\|_{L^2}^2 + \int_0^t \left(\|\nabla \varphi^*\|_{L^2}^2 + \|\operatorname{curl}(u^*)\|_{L^2}^2 + \|\operatorname{div}(u^*)\|_{L^2}^2 \right) ds \\
& \leq c_9 \int_0^t \left(1 + \|U_1\|_{H^1}^2 + \|\phi_1\|_{H^1}^2 + \|U_1\|_{L^4}^4 + \|U_\beta\|_{H^1}^2 + \|\phi_\beta\|_{H^1}^2 \right) \left(\|u^*\|_{L^2}^2 + \|\varphi^*\|_{L^2}^2 \right) ds \\
& \quad + c_{10} \int_0^t \left(\|U_\beta\|_{H^1}^2 + \|\phi_\beta\|_{H^1}^2 \right) \left(\|\phi_\beta\|_{L^2} + \|U_\beta\|_{L^2} \right) ds + \int_0^t \left(\|\phi_\beta\|_{L^4}^4 + \|U_\beta\|_{L^4}^4 \right) ds.
\end{aligned} \tag{4.17}$$

By using again Nirenberg-Gagliardo inequality, we have (since all functions in previous result are in $L^\infty(0, T, L^2(\Omega))$)

$$\begin{aligned}
& \left(\|\varphi^*\|_{L^2}^2 + \|u^*\|_{L^2}^2 \right) + \int_0^t \left(\|\nabla \varphi^*\|_{L^2}^2 + \|\operatorname{curl}(u^*)\|_{L^2}^2 + \|\operatorname{div}(u^*)\|_{L^2}^2 \right) ds \\
& \leq c_{11} \int_0^t \left(1 + \|U_1\|_{H^1}^2 + \|\phi_1\|_{H^1}^2 + \|U_\beta\|_{H^1}^2 + \|\phi_\beta\|_{H^1}^2 \right) \left(\|u^*\|_{L^2}^2 + \|\varphi^*\|_{L^2}^2 \right) ds \\
& \quad + c_{12} \left(\|(\phi_\beta, U_\beta)\|_{\mathbb{E} \times \mathbf{E}_n}^3 + \|(\phi_\beta, U_\beta)\|_{\mathbb{E} \times \mathbf{E}_n}^4 \right).
\end{aligned} \tag{4.18}$$

Using now Gronwall's formula and the stability result (4.11), we can deduce that (since all functions in previous inequality are bounded in $L^2(0, T, H^1(\Omega))$) $\|(\varphi^*, u^*)\|_{\mathbb{E} \times \mathbf{E}_n} \leq \|\beta\|_{L^2}^{3/2}$. Therefore, we have the first part of (i).

Prove now the second part of (i).

Let $(\varphi_{0i}, u_{0i}, h_i)$ be in $\mathcal{L}^2_\infty(\Omega) \times \mathbf{L}^2(\Omega) \times L^2(\mathcal{Q})$, $i = 1, 2$, (given) and $(\psi_i, w_i) = \mathcal{F}'_\varphi(\varphi_{0i}, u_{0i}, h_i) \cdot \beta$ be solution of (\mathcal{P}_{FP}) (we denote by $(\varphi_i, u_i) = \mathcal{F}(\varphi_{0i}, u_{0i}, h_i)$ and by $(\varphi, u) = (\varphi_1 - \varphi_2, u_1 - u_2)$).

Set $(\psi, w) = (\psi_1 - \psi_2, w_1 - w_2)$, $(\varphi_0, u_0, h) = (\varphi_{01} - \varphi_{02}, u_{01} - u_{02}, h_1 - h_2)$, $U_i = u_i + U$ and $\phi_i = \varphi_i + \phi$, $(i = 1, 2)$. According to the equations satisfied by (ψ_1, w_1) and (ψ_2, w_2) we have $(\forall (q, v) \in \mathcal{H}^1(\Omega) \times V_n$ and a.e. $t \in (0, T)$):

$$\begin{aligned}
& \eta \int_\Omega \rho \frac{\partial \psi}{\partial t} q dx - i\eta\kappa \int_\Omega (\operatorname{div}(\rho U_2)\psi + \operatorname{div}(\rho w)\phi_1) q dx \\
& \quad - i\eta\kappa \int_\Omega (\operatorname{div}(\rho u)\psi_1 + \operatorname{div}(\rho w_2)\varphi) q dx \\
& \quad - i\mu \int_\Gamma \rho \psi q d\Gamma + \int_\Omega \rho (B(u_2)(\psi) + u\psi_1 + \varphi w_1 + \phi_2 w) \overline{B}(u_2)(q) dx \\
& \quad + \int_\Omega \rho (B(u_1)(\psi_1) + w_1 \phi_1) u q dx + \int_\Omega \rho B(u_1)(\phi_1) w q dx \\
& \quad + \int_\Omega \rho (B(u_2)(\varphi) + u\phi_1) w_2 q dx \\
& \quad + \int_\Omega \rho ((2|\phi_1|^2 - 1)\psi + \phi_1^2 \overline{\psi}) q dx \\
& \quad + \int_\Omega \rho (2(|\varphi|^2 + 2\mathcal{R}(\varphi \overline{\varphi_2}))\psi_2 + \varphi(\phi_1 + \phi_2) \overline{\psi_2}) q dx = 0,
\end{aligned}$$

$$\begin{aligned}
& \int_{\Omega} \rho \frac{\partial w}{\partial t} v dx + \int_{\Omega} \rho \operatorname{curl}(w) \operatorname{curl}(v) dx + \int_{\Omega} \operatorname{div}(\rho w) \operatorname{div}(v) dx \\
& \quad + \int_{\Omega} \rho (\mathcal{R}(\overline{\psi} B(u_1)(\phi_1)) + w |\phi_1|^2) v dx \\
& \quad + \int_{\Omega} \rho \mathcal{R}(\overline{\phi_2}(B(u_2)(\psi) + u\psi_1) + \varphi B(u_1)(\psi_1)) v dx \\
& \quad + \int_{\Omega} \rho \mathcal{R}(\overline{\psi_2}(B(u_2)(\varphi) + u\phi_1)) v dx + \int_{\Omega} \rho w_2 (|\varphi|^2 + 2\mathcal{R}(\varphi\phi_2)) v dx = 0, \\
& (\psi(0), w(0)) = (0, 0).
\end{aligned} \tag{4.19}$$

Putting $(q, v) = (\overline{\psi}, w)$ in (4.19), we have then

$$\begin{aligned}
& \eta \int_{\Omega} \rho \frac{\partial \psi}{\partial t} \overline{\psi} dx - i\eta\kappa \int_{\Omega} (\operatorname{div}(\rho U_2) |\psi|^2 + \operatorname{div}(\rho w) \phi_1 \overline{\psi}) dx \\
& \quad - i\eta\kappa \int_{\Omega} (\operatorname{div}(\rho u) \psi_1 + \operatorname{div}(\rho w_2) \varphi) \overline{\psi} dx \\
& \quad - i\mu \int_{\Gamma} \rho |\psi|^2 d\Gamma + \int_{\Omega} \rho (B(u_2)(\psi) + u\psi_1 + \varphi w_1 + \phi_2 w) \overline{B(u_2)(\psi)} dx \\
& \quad + \int_{\Omega} \rho (B(u_1)(\psi_1) + w_1 \phi_1) u \overline{\psi} dx + \int_{\Omega} \rho B(u_1)(\phi_1) w \overline{\psi} dx \\
& \quad + \int_{\Omega} \rho (B(u_2)(\varphi) + u\phi_1) w_2 \overline{\psi} dx \\
& \quad + \int_{\Omega} \rho ((2|\phi_1|^2 - 1)\psi + \phi_1^2 \overline{\psi}) \overline{\psi} dx \\
& \quad + \int_{\Omega} \rho (2(|\varphi|^2 + 2\mathcal{R}(\varphi\phi_2)) \psi_2 + \varphi(\phi_1 + \phi_2) \overline{\psi_2}) \overline{\psi} dx = 0,
\end{aligned} \tag{4.20}$$

$$\begin{aligned}
& \int_{\Omega} \rho \frac{\partial |w|^2}{2\partial t} dx + \int_{\Omega} \rho |\operatorname{curl}(w)|^2 dx + \int_{\Omega} \operatorname{div}(\rho w) \operatorname{div}(w) dx \\
& \quad + \int_{\Omega} \rho (\mathcal{R}(\overline{\psi} B(u_1)(\phi_1)) + w |\phi_1|^2) w dx \\
& \quad + \int_{\Omega} \rho \mathcal{R}(\overline{\phi_2}(B(u_2)(\psi) + u\psi_1) + \varphi B(u_1)(\psi_1)) w dx \\
& \quad + \int_{\Omega} \rho \mathcal{R}(\overline{\psi_2}(B(u_2)(\varphi) + u\phi_1)) w dx + \int_{\Omega} \rho w_2 (|\varphi|^2 + 2\mathcal{R}(\varphi\phi_2)) w dx = 0, \\
& (\psi(0), w(0)) = (0, 0).
\end{aligned}$$

By taking the real part of the first part of previous system, we obtain (since $(\phi_i)_{i=1,2}$, φ are

in $\mathcal{L}^\infty(\mathcal{Q})$ and according to Lemma 1.6)

$$\begin{aligned}
& \eta\rho_0 \frac{d\|\psi\|_{L^2}^2}{2dt} + \frac{\rho_0}{\kappa^2} \|\nabla\psi\|_{L^2}^2 \\
& \leq C \left(\|w\|_{H^1} \|\psi\|_{L^2} + \|u\|_{H^1} \|\psi_1\|_{L^4} \|\psi\|_{L^4} + \|w_2\|_{H^1} \|\varphi\|_{L^4} \|\psi\|_{L^4} \right. \\
& \quad + (\|\psi\|_{H^1} + \|U_2\|_{L^4} \|\psi\|_{L^4}) (\|u\|_{L^4} \|\psi_1\|_{L^4} + \|w_1\|_{L^4} \|\varphi\|_{L^4} + \|w\|_{L^2} + \|U_2\|_{L^4} \|\psi\|_{L^4}) \\
& \quad + \|u\|_{L^4} \|\psi\|_{L^4} (\|w_1\|_{L^2} + \|\psi_1\|_{H^1} + \|U_1\|_{L^4} \|\psi_1\|_{L^4}) \\
& \quad + \|\psi\|_{L^4} \|w\|_{L^4} (\|\phi_1\|_{H^1} + \|U_1\|_{L^2}) \\
& \quad + \|w_2\|_{L^4} \|\psi\|_{L^4} (\|\varphi\|_{H^1} + \|U_2\|_{L^4} \|\varphi\|_{L^4} + \|u\|_{L^2}) + \|\psi\|_{L^2}^2 \\
& \quad \left. + \|\psi_2\|_{L^4} \|\varphi\|_{L^4} \|\psi\|_{L^4} \right), \\
& \rho_0 \frac{d\|w\|_{L^2}^2}{2dt} + \rho_0 (\|\operatorname{curl}(w)\|_{L^2}^2 + \|\operatorname{div}(w)\|_{L^2}^2) \\
& \leq \frac{\rho_0}{4} (\|\operatorname{curl}(w)\|_{L^2}^2 + \|\operatorname{div}(w)\|_{L^2}^2) \\
& \quad + C \|w\|_{L^4} \|\psi\|_{L^4} (\|\phi_1\|_{H^1} + \|U_1\|_{L^2}) \\
& \quad + \|w\|_{L^2} (1 + \|\psi\|_{H^1} + \|U_2\|_{L^4} \|\psi\|_{L^4} + \|u\|_{L^4} \|\psi_1\|_{L^4}) \\
& \quad + \|w\|_{L^4} \|\varphi\|_{L^4} (\|\psi_1\|_{H^1} + \|U_1\|_{L^4} \|\psi_1\|_{L^4} + \|w_2\|_{L^4}) \\
& \quad + \|w\|_{L^4} \|\psi_2\|_{L^4} (\|U_2\|_{L^4} \|\varphi\|_{L^4} + \|\varphi\|_{H^1} + \|u\|_{L^2}).
\end{aligned} \tag{4.21}$$

By using Lemma 1.7 and Young's formula, we can deduce that

$$\begin{aligned}
& \eta\rho_0 \frac{d\|\psi\|_{L^2}^2}{2dt} + \frac{\rho_0}{\kappa^2} \|\nabla\psi\|_{L^2}^2 \\
& \leq \frac{\rho_0}{4\kappa^2} \|\nabla\psi\|_{L^2}^2 + \frac{\rho_0}{4} (\|\operatorname{curl}(w)\|_{L^2}^2 + \|\operatorname{div}(w)\|_{L^2}^2) \\
& \quad + c_0 \|\psi\|_{L^2} (1 + \|U_2\|_{H^1}^2 + \|U_2\|_{L^4}^4) \\
& \quad + c_1 \|\psi\|_{L^2} (\|(\phi_1, U_1)\|_{H^1}^2 + \|(\varphi, u)\|_{H^1}^2 + \|(\psi_1, w_1)\|_{H^1}^2 + \|(\psi_2, w_2)\|_{H^1}^2) \\
& \quad + c_2 (1 + \|(\phi_1, U_1)\|_{H^1}^2) \|w\|_{L^2}^2 + c_3 \|(\varphi, u)\|_{L^4}^2 (\|(\psi_1, w_1)\|_{L^4}^2 + \|(\psi_2, w_2)\|_{L^4}^2) \\
& \quad + c_4 (\|(\varphi, u)\|_{L^2} (\|(\psi_1, w_1)\|_{H^1}^2 + \|(\psi_2, w_2)\|_{H^1}^2) \\
& \quad \quad + \|(\varphi, u)\|_{H^1}^2 (\|(\psi_1, w_1)\|_{L^2} + \|(\psi_2, w_2)\|_{L^2})),
\end{aligned}$$

$$\begin{aligned}
& \rho_0 \frac{d\|w\|_{L^2}^2}{2dt} + \rho_0 \left(\|\operatorname{curl}(w)\|_{L^2}^2 + \|\operatorname{div}(w)\|_{L^2}^2 \right) \\
& \leq \frac{\rho_0}{4} \left(\|\operatorname{curl}(w)\|_{L^2}^2 + \|\operatorname{div}(w)\|_{L^2}^2 \right) + \frac{\rho_0}{4\kappa^2} \|\nabla\psi\|_{L^2}^2 + c_5 \|w\|_{L^2} \left(1 + \|U_2\|_{H^1}^2 + \|U_2\|_{L^4}^4 \right) \\
& \quad + c_6 \|w\|_{L^2} \left(\|(\phi_1, U_1)\|_{H^1}^2 + \|(\varphi, u)\|_{H^1}^2 + \|(\psi_1, w_1)\|_{H^1}^2 + \|(\psi_2, w_2)\|_{H^1}^2 \right) \\
& \quad + c_7 \left(1 + \|(\phi_1, U_1)\|_{H^1}^2 + \|U_2\|_{L^4}^4 \right) \|\psi\|_{L^2}^2 \\
& \quad + c_8 \|(\varphi, u)\|_{L^4}^2 \left(\|(\psi_1, w_1)\|_{L^4}^2 + \|(\psi_2, w_2)\|_{L^4}^2 \right) \\
& \quad + c_9 \left(\|(\varphi, u)\|_{L^2} \left(\|(\psi_1, w_1)\|_{H^1}^2 + \|(\psi_2, w_2)\|_{H^1}^2 \right) \right. \\
& \quad \quad \left. + \|(\varphi, u)\|_{H^1}^2 \left(\|(\psi_1, w_1)\|_{L^2} + \|(\psi_2, w_2)\|_{L^2} \right) \right).
\end{aligned} \tag{4.22}$$

Integrating over $(0, t)$ for $t \in (0, T)$ and using Theorems 3.2 and 4.2, we obtain then (since (ϕ_i, U_i) , (ψ_i, w_i) , $i = 1, 2$, and (φ, u) are in $L^\infty(0, T, L^2(\Omega))$)

$$\begin{aligned}
& \eta\rho_0 \|\psi\|_{L^2}^2 + \frac{2\rho_0}{\kappa^2} \int_0^t \|\nabla\psi\|_{L^2}^2 ds \\
& \leq \frac{\rho_0}{2\kappa^2} \int_0^t \|\nabla\psi\|_{L^2}^2 ds + \frac{\rho_0}{2} \int_0^t \left(\|\operatorname{curl}(w)\|_{L^2}^2 + \|\operatorname{div}(w)\|_{L^2}^2 \right) ds \\
& \quad + c_{10} \int_0^t \|\psi\|_{L^2} \left(1 + \|U_2\|_{H^1}^2 \right) ds \\
& \quad + c_{11} \int_0^t \|\psi\|_{L^2} \left(\|(\phi_1, U_1)\|_{H^1}^2 + \|(\varphi, u)\|_{H^1}^2 + \|(\psi_1, w_1)\|_{H^1}^2 + \|(\psi_2, w_2)\|_{H^1}^2 \right) ds \\
& \quad + c_{12} \int_0^t \left(1 + \|(\phi_1, U_1)\|_{H^1}^2 \right) \|w\|_{L^2}^2 ds + c_{13} \left(\|X\|_{L^2}^2 \|\beta\|_{L^2} + \|\beta\|_{L^2}^2 \|X_h\|_{L^2} \right), \\
& \rho_0 \|w\|_{L^2}^2 + 2\rho_0 \int_0^t \left(\|\operatorname{curl}(w)\|_{L^2}^2 + \|\operatorname{div}(w)\|_{L^2}^2 \right) ds \\
& \leq \frac{\rho_0}{2} \int_0^t \left(\|\operatorname{curl}(w)\|_{L^2}^2 + \|\operatorname{div}(w)\|_{L^2}^2 \right) ds + \frac{\rho_0}{2\kappa^2} \int_0^t \|\nabla\psi\|_{L^2}^2 ds \\
& \quad + c_{14} \int_0^t \|w\|_{L^2} \left(1 + \|U_2\|_{H^1}^2 \right) ds \\
& \quad + c_{15} \int_0^t \|w\|_{L^2} \left(\|(\phi_1, U_1)\|_{H^1}^2 + \|(\varphi, u)\|_{H^1}^2 + \|(\psi_1, w_1)\|_{H^1}^2 + \|(\psi_2, w_2)\|_{H^1}^2 \right) ds \\
& \quad + c_{16} \int_0^t \left(1 + \|(\phi_1, U_1)\|_{H^1}^2 + \|U_2\|_{L^4}^4 \right) \|\psi\|_{L^2}^2 ds + c_{17} \left(\|X\|_{L^2}^2 \|\beta\|_{L^2} + \|\beta\|_{L^2}^2 \|X_h\|_{L^2} \right),
\end{aligned} \tag{4.23}$$

where $X_h = (\varphi_0, u_0, h)$.

The previous statements lead to

$$\begin{aligned}
& \|\psi\|_{L^2}^2 + \|w\|_{L^2}^2 + \int_0^t \left(\|\nabla\psi\|_{L^2}^2 + \|\operatorname{curl}(w)\|_{L^2}^2 + \|\operatorname{div}(w)\|_{L^2}^2 \right) ds \\
& \leq c_{18} \int_0^t \left(\|\psi\|_{L^2}^2 + \|w\|_{L^2}^2 \right) \left(1 + \|(\phi_1, U_1)\|_{H^1}^2 + \|U_2\|_{H^1}^2 \right) ds \\
& \quad + c_{19} \int_0^t \left(\|\psi\|_{L^2}^2 + \|w\|_{L^2}^2 \right) \left(\|(\varphi, u)\|_{H^1}^2 + \|(\psi_1, w_1)\|_{H^1}^2 + \|(\psi_2, w_2)\|_{H^1}^2 \right) ds \\
& \quad + c_{20} \left(\|X\|_{L^2}^2 \|\beta\|_{L^2} + \|\beta\|_{L^2}^2 \|X_h\|_{L^2} \right).
\end{aligned} \tag{4.24}$$

According to Gronwall's formula, we can deduce the second part of (i) (since $(\phi_i, U_i)_{i=1,2}$, $(\psi_i, w_i)_{i=1,2}$ and (φ, u) are in $L^2(0, T, H^1(\Omega))$).

By using the same technique to prove the results of (i), we have the result of (ii) and (iii). So, we omit the tedious details. \square

5. Robust control problem

The objective of the robust control problem is to find the best admissible control in the presence of the worst disturbance which maximally spoils the control objective. We formulated the problem in two situations: firstly the case where the control is in the external magnetic field and secondly the case where the control is in the initial condition u_0 (data assimilation).

5.1. The control in the external magnetic field. In this section, we consider two situations: firstly the case where the worst disturbance is in the external magnetic field h and secondly the case where the disturbance is in the initial condition φ_0 .

5.1.1. Distributed disturbance in the external magnetic field. We suppose now that the value h is decomposed into the disturbance $f \in L^2(\mathcal{Q})$ and the control $g \in L^2(\mathcal{Q})$, that is, $h = f + g$. So the function (φ, u) is assumed to be related to the disturbance f and control g through the problem (3.3) ($\forall (q, v) \in \mathcal{H}^1(\Omega) \times V_n$ and a.e. $t \in (0, T)$)

$$\begin{aligned}
& \eta \int_{\Omega} \rho \frac{\partial \varphi}{\partial t} q dx - i\eta\kappa \int_{\Omega} \operatorname{div}(\rho(u + U)) \varphi q dx - i\eta\kappa \int_{\Omega} \operatorname{div}(\rho u) \varphi q dx - i\mu \int_{\Gamma} \rho \varphi q d\Gamma \\
& \quad + \int_{\Omega} \rho B(u)(\varphi + \phi) \bar{B}(u)(q) dx + \int_{\Omega} \rho F(\varphi) q dx = \int_{\Omega} \rho B(0)(\phi) \bar{B}(0)(q) dx, \\
& \int_{\Omega} \rho \frac{\partial u}{\partial t} v dx + \int_{\Omega} \rho \operatorname{curl}(u) \operatorname{curl}(v) dx + \int_{\Omega} \operatorname{div}(\rho u) \operatorname{div}(v) dx \\
& \quad + \int_{\Omega} \rho \mathcal{R}((\bar{\varphi} + \bar{\phi}) B(u)(\varphi + \phi)) v dx \\
& = \int_{\Omega} \rho \mathcal{R}(\bar{\phi} B(0)(\phi)) v dx + \int_{\Omega} \rho (f + g) \operatorname{curl}(v) dx, \\
& (\varphi(0), u(0)) = (\varphi_0, u_0).
\end{aligned} \tag{5.1}$$

To obtain the regularity of Theorem 4.2, we suppose the following hypotheses: $(\varphi_0, u_0) \in \mathcal{L}_\infty^2(\Omega) \times L^2(\Omega)$. Let $\mathcal{P} : (f, g) \rightarrow (\varphi, u) = \mathcal{P}(f, g)$ be the map: $(L^2(\mathcal{Q}))^2 \rightarrow \mathcal{E} \times \mathbf{E}_n$ defined by (5.1) and introduce the cost function defined by

$$J(f, g) = \frac{a}{4} \|\varphi\|^2 - \Lambda \|u\|_{L^2(\mathcal{Q})}^2 + \frac{b}{2} \|u - u_{\text{obs}}\|_{L^2(\mathcal{Q})}^2 + \frac{\alpha}{2} \|f\|_{L^2(\mathcal{Q})}^2 - \frac{\gamma}{2} \|g\|_{L^2(\mathcal{Q})}^2, \quad (5.2)$$

where a, b, α, γ are fixed such that $\alpha, \gamma > 0$, $a, b \geq 0$, and $a + b > 0$. The functions $u_{\text{obs}} \in L^2(\mathcal{Q})$ and $\Lambda \in L^\infty(\mathcal{Q})$ are given and represent the observation.

Let $\mathcal{H} = \mathcal{H}_1 \times \mathcal{H}_2$ such that \mathcal{H}_1 and \mathcal{H}_2 are (given) nonempty, closed, convex, bounded subsets of $L^2(\mathcal{Q})$. We want to minimize the functional J with respect to f and maximize J with respect to g , that is, to study the following problem (\mathcal{MP}_1) :

find an admissible control $f^* \in \mathcal{H}_1$ and a disturbance $g^* \in \mathcal{H}_2$ such that (f^*, g^*) is a saddle point of the functional J on \mathcal{H} , subject to system (5.1).

Such a pair (f^*, g^*) is called an optimal solution to (\mathcal{MP}_1) .

PROPOSITION 5.1. *The function \mathcal{P} is continuously Fréchet differentiable from $(L^2(\mathcal{Q}))^2$ to $\mathcal{E} \times \mathbf{E}_n$ with the derivative $\mathcal{P}'(f, g) : Y = (\beta_1, \beta_2) \rightarrow (\psi, w)$ given by the linear problem (\mathcal{P}_{F1}) ($\forall (q, v) \in \mathcal{H}^1(\Omega) \times V_n$ and a.e. $t \in (0, T)$)*

$$\begin{aligned} & \eta \int_\Omega \rho \frac{\partial \psi}{\partial t} q \, dx - i\eta\kappa \int_\Omega \operatorname{div}(\rho U_1) \psi q \, dx - i\eta\kappa \int_\Omega \operatorname{div}(\rho w) \phi_1 q \, dx - i\mu \int_\Gamma \rho \psi q \, d\Gamma \\ & + \int_\Omega \rho (B(u)(\psi) + \phi_1 w) \bar{B}(u)(q) \, dx + \int_\Omega \rho B(u)(\phi_1) w q \, dx \\ & + \int_\Omega \rho ((2|\phi_1|^2 - 1)\psi + \phi_1^2 \bar{\psi}) q \, dx = 0, \\ & \int_\Omega \rho \frac{\partial w}{\partial t} v \, dx + \int_\Omega \rho \operatorname{curl}(w) \operatorname{curl}(v) \, dx + \int_\Omega \operatorname{div}(\rho w) \operatorname{div}(v) \, dx + \int_\Omega \rho \mathcal{R}(\bar{\psi} B(u)(\phi_1)) v \, dx \\ & + \int_\Omega \rho (\mathcal{R}(\bar{\phi}_1 B(u)(\psi)) + w |\phi_1|^2) v \, dx = \int_\Omega \rho (\beta_1 + \beta_2) \operatorname{curl}(v) \, dx, \\ & (\psi(0), w(0)) = (0, 0), \end{aligned} \quad (5.3)$$

where $(U_1, \phi_1) = (u + U, \varphi + \phi)$.

Moreover the following estimates exist: ($\forall (f_i, g_i) \in (L^2(\Omega))^2$, $i = 1, 2$)

$$(i) \|\mathcal{P}'(f_1, g_1)\|_{\mathcal{L}((L^2(\mathcal{Q}))^2, \mathcal{E} \times \mathbf{E}_n)} \leq C_e;$$

$$(ii) \|\mathcal{P}'(f_1, g_1)Y - \mathcal{P}'(f_2, g_2)Y\|_{\mathcal{E} \times \mathbf{E}_n}^2 \leq C_e (\|X\|_{L^2(\mathcal{Q})} \|Y\|_{L^2(\mathcal{Q})}^2 + \|X\|_{L^2(\mathcal{Q})}^2 \|Y\|_{L^2(\mathcal{Q})}),$$

where $f = f_1 - f_2$, $g = g_1 - g_2$ and $X = (f, g)$.

Proof. The proof of this proposition is the consequence of the result of Theorem 4.3. Here we omit the tedious details. \square

PROPOSITION 5.2. *The map \mathcal{P} defined by (5.1) is continuous from the weak topology of $(L^2(\mathcal{Q}))^2$ to the strong topology of $\mathcal{L}^2(\mathcal{Q}) \times \mathbf{L}^2(\mathcal{Q})$.*

Proof. Let $\mathbf{f} = (f, g)$ be given in $(L^2(\mathcal{Q}))^2$ and let be a sequence $\mathbf{f}_k = (f_k, g_k)$ such that \mathbf{f}_k is weakly convergent in $(L^2(\mathcal{Q}))^2$ to \mathbf{f} .

Set $(\varphi, u) = \mathcal{P}(f, g)$ and $(\varphi_k, u_k) = \mathcal{P}(f_k, g_k)$. Since $\mathbf{f}_k \rightharpoonup \mathbf{f}$ weakly in $(L^2(\mathcal{Q}))^2$ then \mathbf{f}_k is uniformly bounded in $(L^2(\mathcal{Q}))^2$. In view of Theorem 3.2, we can deduce that the sequence (φ_k, u_k) is uniformly bounded in $\mathcal{E} \times \mathbf{E}_n$. Therefore we can extract from $(\mathbf{f}_k, \varphi_k, u_k)$ a subsequence also denoted by $(\mathbf{f}_k, \varphi_k, u_k)$ and such that

$$\begin{aligned} (f_k, g_k) &\longrightarrow (f, g) \text{ weakly in } (L^2(\mathcal{Q}))^2, \\ (\varphi_k, u_k) &\longrightarrow (\tilde{\varphi}, \tilde{u}) \text{ weakly in } \mathcal{E} \times \mathbf{E}_n, \\ (\varphi_k, u_k) &\longrightarrow (\tilde{\varphi}, \tilde{u}) \text{ strongly in } \mathcal{L}^2(\mathcal{Q}) \times \mathbf{L}^2(\mathcal{Q}), \\ \varphi_k &\longrightarrow \tilde{\varphi} \text{ weakly in } \mathcal{L}^2(\Sigma). \end{aligned} \tag{5.4}$$

We prove easily that $(\tilde{\varphi}, \tilde{u}) = \mathcal{P}(f, g)$ and according to the uniqueness of the solution of (5.1), we have then $\tilde{\varphi} = \varphi$ and $\tilde{u} = u$. \square

THEOREM 5.3. *For α and γ sufficiently large (i.e., there exists (α_1, γ_1) such that $\alpha \geq \alpha_1$ and $\gamma \geq \gamma_1$), there exists $(f^*, g^*) \in \mathcal{K}$ and $(\varphi^*, u^*) \in \mathcal{E} \times \mathbf{E}_n$ such that (f^*, g^*) is defined by $(\mathcal{M}\mathcal{P}_1)$ and $(\varphi^*, u^*) = \mathcal{P}(f^*, g^*)$ is solution of (5.1).*

Proof. Let P_g be the map $f \rightarrow J(f, g)$ and let Q_f be the map: $g \rightarrow J(f, g)$. To obtain the existence of the robust control problem, we prove that P_g is convex and lower semicontinuous for all $g \in \mathcal{K}_2$, and Q_f is concave and upper semicontinuous for all $f \in \mathcal{K}_1$ and we use the classical minimax theorem in infinite dimensions (see, e.g., [3, 23]).

Firstly we prove, for α and γ sufficiently large, the convexity of the map P_g and the concavity of the map Q_f . In order to prove the convexity, it is sufficient to show that for all $(f_1, f_2) \in \mathcal{K}_1$ we have $(P'_g(f_1) - P'_g(f_2)) \cdot f \geq 0$, where $f = f_1 - f_2$ (because P_g is Gâteaux differentiable). According to the definition of J , we have that

$$\begin{aligned} &(P'_g(f_1) - P'_g(f_2)) \cdot f \\ &= \alpha \|f\|_{L^2(\mathcal{Q})}^2 + a \iint_{\mathcal{Q}} \mathcal{R}((|\varphi_1|^2 - |\varphi_2|^2) \psi_2 \bar{\varphi}_2) dx dt \\ &\quad + a \iint_{\mathcal{Q}} \mathcal{R}((|\varphi_1|^2 - \Lambda) (\varphi_1 - \varphi_2) \bar{\psi}_1) dx dt + b \iint_{\mathcal{Q}} (u_1 - u_2) w_2 dx dt \\ &\quad + a \iint_{\mathcal{Q}} \mathcal{R}((|\varphi_1|^2 - \Lambda) (\psi_1 - \psi_2) \bar{\varphi}_2) dx dt + b \iint_{\mathcal{Q}} (u_1 - u_{\text{obs}}) (w_1 - w_2) dx dt, \end{aligned} \tag{5.5}$$

where $(\varphi_i, u_i) = \mathcal{P}(f_i, g)$, $(\psi_i, w_i) = \mathcal{P}'(f_i, g) \cdot (f, 0)$ (solution of problem (\mathcal{P}_{F_1})), for $i = 1, 2$.

According to Theorem 3.2 and Proposition 5.1, we have

$$\begin{aligned}
& a \iint_{\mathcal{Q}} (\mathcal{R}(|\varphi_1|^2 - |\varphi_2|^2) \psi_2 \bar{\varphi}_2) + \mathcal{R}(|\varphi_1|^2 - \Lambda)(\varphi_1 - \varphi_2) \bar{\psi}_1) dx dt \\
& \quad + b \iint_{\mathcal{Q}} (u_1 - u_2) w_2 dx dt \\
& \leq C \left(\|\varphi_1 - \varphi_2\|_{L^2(\mathcal{Q})} (\|\psi_1\|_{L^2(\mathcal{Q})} + \|\psi_2\|_{L^2(\mathcal{Q})}) + \|u_1 - u_2\|_{L^2(\mathcal{Q})} \|w_2\|_{L^2(\mathcal{Q})} \right) \\
& \leq C_0 \|f\|_{L^2(\mathcal{Q})}^2, \tag{5.6} \\
& a \iint_{\mathcal{Q}} \mathcal{R}(|\varphi_1|^2 - \Lambda)(\psi_1 - \psi_2) \bar{\varphi}_2) dx dt + b \iint_{\mathcal{Q}} (u_1 - u_{\text{obs}})(w_1 - w_2) dx dt \\
& \leq C \left(\|\varphi_1\|_{L^2(\mathcal{Q})}^2 - \Lambda\|_{L^2(\mathcal{Q})} \|\psi_1 - \psi_2\|_{L^2(\mathcal{Q})} + \|u_1 - u_{\text{obs}}\|_{L^2(\mathcal{Q})} \|w_1 - w_2\|_{L^2(\mathcal{Q})} \right) \\
& \leq C_1 \|f\|_{L^2(\mathcal{Q})}^{3/2}.
\end{aligned}$$

From (5.5)-(5.6) we deduce that for $\alpha \geq \alpha_l$ such that $\alpha_l > C_0$ and $(\alpha_l - C_0) \min_{f \in \mathcal{H}_1} \|f\|_{L^2}^{1/2} = C_1$, we have $(P'_g(f_1) - P'_g(f_2)) \cdot f \geq 0$ and then the convexity of P_g . In the same way, we can find γ_l such that for $\gamma \geq \gamma_l$ we have the concavity of Q_f .

We prove now that P_g is lower semicontinuous for all $g \in \mathcal{H}_2$, and Q_f is upper semicontinuous for all $f \in \mathcal{H}_1$. Let f_k be a minimizing sequence of J , that is, $\liminf_k J(f_k, g) = \min_{f \in \mathcal{H}_1} J(f, g)$ ($\forall g \in \mathcal{H}_2$). Then f_k is uniformly bounded in \mathcal{H}_1 and we can extract from f_k a subsequence also denoted by f_k such that $f_k \rightharpoonup f_g$ weakly in \mathcal{H}_1 . By using Proposition 5.2, we have then

$$\mathcal{P}(f_k, g) \longrightarrow (\varphi_g, u_g) \text{ strongly in } \mathcal{L}^2(\mathcal{Q}) \times \mathbf{L}^2(\mathcal{Q}). \tag{5.7}$$

Therefore, since the norm is lower semicontinuous, we have that the map $P_g : f \rightarrow J(f, g)$ is lower semicontinuous for all $g \in \mathcal{H}_2$. By using the same technique we obtain then Q_f is upper semicontinuous for all $f \in \mathcal{H}_1$. \square

In order to obtain the necessary optimality conditions which have been satisfied by the solution of the robust control problem, we introduce the following adjoint problem corresponding to the primal problem (5.1) (we denote by $(\varphi, u) = \mathcal{P}(f, g)$ and $(\phi_1, U_1) = (\varphi + \phi, u + U)$): find $(P, Q) \in \mathcal{C} \times \mathbf{E}_n$ such that $(\forall (q, v) \in \mathcal{H}^1(\Omega) \times V_n \text{ and a.e. } t \in (0, T))$:

$$\begin{aligned}
& -\eta \int_{\Omega} \rho \frac{\partial P}{\partial t} q dx - i\eta \kappa \int_{\Omega} \text{div}(\rho U_1) P q dx - i\mu \int_{\Gamma} \rho P q d\Gamma + \int_{\Omega} \rho \bar{B}(u)(P) B(u)(q) dx \\
& \quad + \int_{\Omega} \rho \overline{B(u)(\phi_1)} Q q dx + \int_{\Omega} \left(-\frac{i}{\kappa} \text{div}(\rho \bar{\phi}_1 Q) + \rho U_1 \bar{\phi}_1 Q \right) q dx \\
& \quad + \int_{\Omega} \rho ((2|\phi_1|^2 - 1)P + \bar{\phi}_1^2 \bar{P}) q dx = a \int_{\Omega} (|\varphi|^2 - \Lambda) \bar{\varphi} q dx,
\end{aligned}$$

$$\begin{aligned}
& - \int_{\Omega} \rho \frac{\partial Q}{\partial t} v dx + \int_{\Omega} \rho \operatorname{curl}(Q) \operatorname{curl}(v) dx + \int_{\Omega} \operatorname{div}(Q) \operatorname{div}(\rho q) dx + \int_{\Omega} \rho Q |\phi_1|^2 v dx \\
& \quad + \int_{\Omega} \rho \mathcal{R}(i\eta\kappa \nabla(\phi_1 P) + PB(u)(\phi_1) + \phi_1 \overline{B}(u)(P)) v dx = b \int_{\Omega} (u - u_{\text{obs}}) v dx, \\
& (P(T), Q(T)) = (0, 0).
\end{aligned} \tag{5.8}$$

Remark 5.4. (i) The adjoint problem (5.8) is a linear system. By reversing sense of time, that is, $t := T - t$, and by applying the same way to obtain the result of Theorem 4.2, we obtain the existence and the uniqueness of (P, Q) .

(ii) The adjoint system (5.8) is the weak formulation of the following problem:

$$\begin{aligned}
& -\eta\rho \frac{\partial P}{\partial t} - i\eta\kappa \operatorname{div}(\rho U_1)P + \overline{B}(u) \cdot (\rho \overline{B}(u)(P)) + \overline{\rho B(u)(\phi_1)} Q - \frac{i}{\kappa} \operatorname{div}(\rho \overline{\phi_1} Q) + \rho U_1 \overline{\phi_1} Q \\
& \quad + \rho((2|\phi_1|^2 - 1)P + \overline{\phi_1}^2 \overline{P}) = a(|\varphi|^2 - \Lambda) \overline{\varphi}, \\
& -\rho \frac{\partial Q}{\partial t} + \operatorname{curl}(\rho \operatorname{curl}(Q)) - \rho \nabla(\operatorname{div}(Q)) + \rho Q |\phi_1|^2 \\
& \quad + \rho \mathcal{R}(i\eta\kappa \nabla(\phi_1 P) + PB(u)(\phi_1) + \phi_1 \overline{B}(u)(P)) = b(u - u_{\text{obs}}),
\end{aligned} \tag{5.9a}$$

subjected to the boundary conditions

$$\frac{1}{\kappa^2} \frac{\partial P}{\partial n} = \mu P, \quad Q \cdot \mathbf{n} = 0, \quad \operatorname{curl}(Q) = 0, \quad \text{in } \Sigma, \tag{5.9b}$$

and the final condition

$$(P(T), Q(T)) = (0, 0). \tag{5.9c}$$

We can now give the first-order optimality conditions for the robust control problem (\mathcal{MP}_1) .

THEOREM 5.5. *Under the assumptions of Theorem 5.3, the optimal solution $(f^*, g^*, u^*, \varphi^*) \in \mathcal{H} \times \mathbf{E}_n \times \mathcal{E}$, such that (f^*, g^*) is defined by (\mathcal{MP}_1) and $(\varphi^*, u^*) = \mathcal{P}(f^*, g^*)$ solution of (5.1), satisfies*

$$\begin{aligned}
& \iint_{\Omega} (\rho \operatorname{curl}(Q^*) + \alpha f^*)(f^* - f) dx dt \geq 0, \\
& \iint_{\Omega} (\rho \operatorname{curl}(Q^*) - \gamma g^*)(g^* - g) dx dt \leq 0, \quad \forall (f, g) \in \mathcal{H},
\end{aligned} \tag{5.10}$$

where (P^*, Q^*) is solution of the adjoint problem (5.8) (corresponding to (φ^*, u^*)).

Proof. The cost function J is a composition of (Fréchet) differentiable maps then J is differentiable and we have $(\forall Y = (\beta_1, \beta_2) \in \mathcal{H})$

$$\begin{aligned} J'(f, g) \cdot Y &= a \iint_{\mathcal{Q}} \mathcal{R}(|\varphi|^2 - \Lambda) \bar{\varphi} \psi \, dx \, dt + b \iint_{\mathcal{Q}} (u - u_{\text{obs}}) w \, dx \, dt \\ &\quad + \alpha \iint_{\mathcal{Q}} f \beta_1 \, dx \, dt - \gamma \iint_{\mathcal{Q}} g \beta_2 \, dx \, dt, \end{aligned} \quad (5.11)$$

where $(\psi, w) = \mathcal{P}'(f, g) \cdot Y$ is solution of problem (\mathcal{P}_{F1}) .

By taking $(q, v) = (P, Q)$ in (\mathcal{P}_{F1}) , using Green's formula and integrating by time, we obtain (according to the initial condition)

$$\begin{aligned} & - \eta \iint_{\mathcal{Q}} \rho \frac{\partial P}{\partial t} \psi \, dx \, dt + \int_{\Omega} \rho P(T) \psi(T) \, dx - i \eta \kappa \iint_{\mathcal{Q}} \operatorname{div}(\rho U_1) P \psi \, dx \, dt \\ & \quad + i \eta \kappa \iint_{\mathcal{Q}} \rho w \nabla(\phi_1 P) \, dx \, dt - i \mu \iint_{\Sigma} \rho P \psi \, d\Gamma \, dt + \iint_{\mathcal{Q}} \rho \bar{B}(u)(P) B(u)(\psi) \, dx \, dt \\ & \quad + \iint_{\mathcal{Q}} \rho \phi_1 w \bar{B}(u)(P) \, dx \, dt + \iint_{\mathcal{Q}} \rho B(u)(\phi_1) w P \, dx \, dt \\ & \quad + \iint_{\mathcal{Q}} \rho ((2|\phi_1|^2 - 1) P \psi + \phi_1^2 P \bar{\psi}) \, dx \, dt = 0, \\ & - \iint_{\mathcal{Q}} \rho \frac{\partial Q}{\partial t} w \, dx \, dt + \int_{\Omega} \rho Q(T) w(T) \, dx + \iint_{\mathcal{Q}} \rho \operatorname{curl}(Q) \operatorname{curl}(w) \, dx \, dt \\ & \quad + \iint_{\mathcal{Q}} \operatorname{div}(Q) \operatorname{div}(\rho w) \, dx \, dt + \iint_{\mathcal{Q}} \rho \mathcal{R}(\bar{\psi} B(u)(\phi_1)) Q \, dx \, dt \\ & \quad + \iint_{\mathcal{Q}} \rho (\mathcal{R}(\bar{\phi}_1 B(u)(\psi)) Q + w |\phi_1|^2 Q) \, dx \, dt = \iint_{\mathcal{Q}} \rho (\beta_1 + \beta_2) \operatorname{curl}(Q) \, dx \, dt. \end{aligned} \quad (5.12)$$

Since (P, Q) is solution of (5.8), we have that

$$\begin{aligned} & i \eta \kappa \iint_{\mathcal{Q}} \rho w \nabla(\phi_1 P) \, dx \, dt + \iint_{\mathcal{Q}} \rho \phi_1 w \bar{B}(u)(P) \, dx \, dt + \iint_{\mathcal{Q}} \rho B(u)(\phi_1) w P \, dx \, dt \\ & \quad + \iint_{\mathcal{Q}} \rho (\phi_1^2 P \bar{\psi} - \bar{\phi}_1^2 \bar{P} \psi) \, dx \, dt - \iint_{\mathcal{Q}} \rho \overline{B(u)(\phi_1)} Q \psi \, dx \, dt \\ & \quad - \iint_{\mathcal{Q}} \left(-\frac{i}{\kappa} \operatorname{div}(\rho \bar{\phi}_1 Q) + \rho U_1 \bar{\phi}_1 Q \right) \psi \, dx \, dt + a \iint_{\mathcal{Q}} (|\varphi|^2 - \Lambda) \bar{\varphi} \psi \, dx \, dt = 0, \\ & \iint_{\mathcal{Q}} \rho \mathcal{R}(\bar{\psi} B(u)(\phi_1) + \bar{\phi}_1 B(u)(\psi)) Q \, dx \, dt + b \iint_{\mathcal{Q}} (u - u_{\text{obs}}) w \, dx \, dt \\ & \quad - \iint_{\mathcal{Q}} \rho \mathcal{R}(i \eta \kappa \nabla(\phi_1 P) + P B(u)(\phi_1) + \phi_1 \bar{B}(u)(P)) w \, dx \, dt \\ & = \iint_{\mathcal{Q}} \rho (\beta_1 + \beta_2) \operatorname{curl}(Q) \, dx \, dt. \end{aligned} \quad (5.13)$$

By adding the real part of the first part of (5.13) and the second part of (5.13), we obtain (since $\Re(\phi_1^2 P \bar{\psi} - \bar{\phi}_1^2 \bar{P} \psi) = 0$ and $\int_{\Omega} (-i/\kappa) \operatorname{div}(\rho \bar{\phi}_1 Q) + \rho U_1 \bar{\phi}_1 Q \psi \, dx = \int_{\Omega} \rho \bar{\phi}_1 B(u)(\psi) Q \, dx$, because $Q \cdot \mathbf{n} = 0$)

$$a \iint_{\mathcal{Q}} \Re(|\phi|^2 - \Lambda) \bar{\phi} \psi \, dx \, dt + b \iint_{\mathcal{Q}} (u - u_{\text{obs}}) w \, dx \, dt = \iint_{\mathcal{Q}} \rho (\beta_1 + \beta_2) \operatorname{curl}(Q) \, dx \, dt. \quad (5.14)$$

According to the expression of $J'(f, g) \cdot Y$, we can deduce that

$$J'(f, g) \cdot Y = \iint_{\mathcal{Q}} (\rho \operatorname{curl}(Q) + \alpha f) \beta_1 \, dx \, dt + \iint_{\mathcal{Q}} (\rho \operatorname{curl}(Q) - \gamma g) \beta_2 \, dx \, dt. \quad (5.15)$$

Since (f^*, g^*) is an optimal solution, we have

$$\begin{aligned} \frac{\partial J}{\partial f}(f^*, g^*) \cdot (f^* - f) &\geq 0, \\ \frac{\partial J}{\partial g}(f^*, g^*) \cdot (g^* - g) &\leq 0, \quad \forall (f, g) \in \mathcal{H}, \end{aligned} \quad (5.16)$$

and then

$$\begin{aligned} \iint_{\mathcal{Q}} (\rho \operatorname{curl}(Q^*) + \alpha f^*) (f^* - f) \, dx \, dt &\geq 0, \\ \iint_{\mathcal{Q}} (\rho \operatorname{curl}(Q^*) - \gamma g^*) (g^* - g) \, dx \, dt &\leq 0, \quad \forall (f, g) \in \mathcal{H}. \end{aligned} \quad (5.17)$$

This completes the proof. \square

5.1.2. Distributed disturbance in the initial condition of the order parameter. In this section, the disturbance is in φ_0 and the control is in h , that is, $\varphi_0 = g (g \in \mathcal{L}^2_{\infty}(\Omega))$, $h = f (f \in L^2(\mathcal{Q}))$. So the function (φ, u) is assumed to be related to the disturbance f and control g through the problem (3.3) (a.e. $t \in (0, T)$)

$$\begin{aligned} \eta \int_{\Omega} \rho \frac{\partial \varphi}{\partial t} q \, dx - i\eta\kappa \int_{\Omega} \operatorname{div}(\rho(u+U)) \varphi q \, dx - i\eta\kappa \int_{\Omega} \operatorname{div}(\rho u) \phi q \, dx - i\mu \int_{\Gamma} \rho \varphi q \, d\Gamma \\ + \int_{\Omega} \rho B(u)(\varphi + \bar{\phi}) \bar{B}(u)(q) \, dx + \int_{\Omega} \rho F(\varphi) q \, dx = \int_{\Omega} \rho B(0)(\phi) \bar{B}(0)(q) \, dx, \\ \int_{\Omega} \rho \frac{\partial u}{\partial t} v \, dx + \int_{\Omega} \rho \operatorname{curl}(u) \operatorname{curl}(v) \, dx + \int_{\Omega} \operatorname{div}(\rho u) \operatorname{div}(v) \, dx \\ + \int_{\Omega} \rho \Re((\bar{\varphi} + \bar{\phi}) B(u)(\varphi + \phi)) v \, dx = \int_{\Omega} \rho \Re(\bar{\phi} B(0)(\phi)) v \, dx + \int_{\Omega} \rho f \operatorname{curl}(v) \, dx, \\ (\varphi(0), u(0)) = (g, u_0). \end{aligned} \quad (5.18)$$

To obtain the regularity of Theorem 4.2, we suppose that $u_0 \in \mathbf{L}^2(\Omega)$. Let $\mathcal{P} : (f, g) \rightarrow (\varphi, u) = \mathcal{P}(f, g)$ be the map: $L^2(\mathcal{Q}) \times \mathcal{L}^2_{\infty}(\Omega) \rightarrow \mathcal{E} \times \mathbf{E}_n$ defined by (5.18) and introduce

the cost function defined by

$$J(f, g) = \frac{a}{4} \|\varphi\|^2 - \Lambda \|_{L^2(\mathcal{Q})}^2 + \frac{b}{2} \|u - u_{\text{obs}}\|_{L^2(\mathcal{Q})}^2 + \frac{\alpha}{2} \|f\|_{L^2(\mathcal{Q})}^2 - \frac{\gamma}{2} \|g\|_{L^2}^2, \quad (5.19)$$

where $\alpha, \beta > 0$, $a, b \geq 0$ and $a + b > 0$. The functions $(u_{\text{obs}}, \Lambda) \in L^2(\mathcal{Q}) \times L^\infty(\mathcal{Q})$ represent the observation.

In this section, we study the following minimax control problem (\mathcal{MP}_2) :

find an admissible control $f^* \in \mathcal{K}_1$ and a disturbance $g^* \in \mathcal{K}_2$ such that (f^*, g^*) is a saddle point of the functional J on \mathcal{K} , subject to system (5.18),

where, $\mathcal{K} = \mathcal{K}_1 \times \mathcal{K}_2$ such that \mathcal{K}_1 and \mathcal{K}_2 are nonempty, closed, convex, bounded subsets of $L^2(\mathcal{Q})$ and $\mathcal{L}_\infty^2(\Omega)$, respectively.

The proof of the following propositions and theorem of existence is very similar to that of Propositions 5.1, 5.2 and of Theorem 5.3. So we omit the details.

PROPOSITION 5.6. *The function \mathcal{P} is continuously Fréchet differentiable from $L^2(\mathcal{Q}) \times \mathcal{L}_\infty^2(\Omega)$ to $\mathcal{E} \times \mathbf{E}_n$ with the derivative $\mathcal{P}'(f, g) : Y = (\beta_1, \beta_2) \rightarrow (\psi, w)$ given by the linear problem (\mathcal{P}_{F2}) (a.e. $t \in (0, T)$)*

$$\begin{aligned} & \eta \int_{\Omega} \rho \frac{\partial \psi}{\partial t} q dx - i\eta\kappa \int_{\Omega} \operatorname{div}(\rho U_1) \psi q dx - i\eta\kappa \int_{\Omega} \operatorname{div}(\rho w) \phi_1 q dx - i\mu \int_{\Gamma} \rho \psi q d\Gamma \\ & + \int_{\Omega} \rho (B(u)(\psi) + \phi_1 w) \bar{B}(u)(q) dx + \int_{\Omega} \rho B(u)(\phi_1) w q dx \\ & + \int_{\Omega} \rho ((2|\phi_1|^2 - 1)\psi + \phi_1^2 \bar{\psi}) q dx = 0, \\ & \int_{\Omega} \rho \frac{\partial w}{\partial t} u dx + \int_{\Omega} \rho \operatorname{curl}(w) \operatorname{curl}(v) dx + \int_{\Omega} \operatorname{div}(\rho w) \operatorname{div}(v) dx + \int_{\Omega} \rho \mathcal{R}(\bar{\psi} B(u)(\phi_1)) v dx \\ & + \int_{\Omega} \rho (\mathcal{R}(\bar{\phi}_1 B(u)(\psi)) + w |\phi_1|^2) v dx = \int_{\Omega} \rho \beta_1 \operatorname{curl}(v) dx, \\ & (\psi(0), w(0)) = (\beta_2, 0), \end{aligned} \quad (5.20)$$

where $(U_1, \phi_1) = (u + U, \varphi + \phi)$.

Moreover the following estimates $(\forall (f_i, g_i) \in L^2(\mathcal{Q}) \times \mathcal{L}_\infty^2(\Omega), i = 1, 2)$ exist:

$$(i) \|\mathcal{P}'(f_1, g_1)\|_{\mathcal{L}(L^2(\mathcal{Q}) \times L^2, \mathcal{E} \times \mathbf{E}_n)} \leq C_e,$$

$$(ii) \|\mathcal{P}'(f_1, g_1)Y - \mathcal{P}'(f_2, g_2)Y\|_{\mathcal{E} \times \mathbf{E}_n}^2 \leq C_e (\|X\|_{L^2(\mathcal{Q}) \times L^2} \|Y\|_{L^2(\mathcal{Q}) \times L^2}^2 + \|X\|_{L^2(\mathcal{Q}) \times L^2} \|Y\|_{L^2(\mathcal{Q}) \times L^2}),$$

where $f = f_1 - f_2$, $g = g_1 - g_2$ and $X = (f, g)$.

PROPOSITION 5.7. *The map \mathcal{P} defined by (5.18) is continuous from the weak topology of $L^2(\mathcal{Q}) \times \mathcal{L}_\infty^2(\Omega)$ to the strong topology of $\mathcal{L}^2(\mathcal{Q}) \times L^2(\mathcal{Q})$.*

THEOREM 5.8. *For α and γ sufficiently large, there exist $(f^*, g^*) \in \mathcal{K}$ and $(u^*, \varphi^*) \in \mathbf{E}_n \times \mathcal{E}$ such that (f^*, g^*) is defined by (\mathcal{MP}_2) and $(\varphi^*, u^*) = \mathcal{P}(f^*, g^*)$ is solution of (5.18).*

Now we establish necessary optimality conditions for the robust control problem (\mathcal{MP}_2) .

THEOREM 5.9. *Under the assumptions of Theorem 5.8, the optimal solution $(f^*, g^*, u^*, \varphi^*) \in \mathcal{H} \times \mathbf{E}_n \times \mathcal{E}$, such that (f^*, g^*) is defined by (\mathcal{MP}_2) and $(\varphi^*, u^*) = \mathcal{P}(f^*, g^*)$ is solution of (5.18), satisfies*

$$\begin{aligned} & \iint_{\mathcal{Q}} (\rho \operatorname{curl}(Q^*) + \alpha f^*)(f^* - f) dx dt \geq 0, \\ & \int_{\Omega} \mathcal{R}((\eta \rho P^*(0) - \gamma \bar{g}^*)(g^* - g)) dx \leq 0, \quad \forall (f, g) \in \mathcal{H}, \end{aligned} \quad (5.21)$$

where (P^*, Q^*) is solution of the adjoint problem (5.8) (corresponding to (φ^*, u^*)).

Proof. The cost function J is a composition of (Fréchet) differentiable maps, then J is differentiable and we have $(\forall Y = (\beta_1, \beta_2) \in \mathcal{H})$

$$\begin{aligned} J'(f, g) \cdot Y &= a \iint_{\mathcal{Q}} \mathcal{R}((|\varphi|^2 - \Lambda)\bar{\varphi}\psi) dx dt + b \iint_{\mathcal{Q}} (u - u_{\text{obs}}) w dx dt + \\ & \alpha \iint_{\mathcal{Q}} f \beta_1 dx dt - \gamma \int_{\Omega} \mathcal{R}(\bar{g}\beta_2) dx, \end{aligned} \quad (5.22)$$

where $(\psi, w) = \mathcal{P}'(f, g) \cdot Y$ is solution of problem (\mathcal{P}_{F_2}) .

By taking $(q, \nu) = (P, Q)$ in (\mathcal{P}_{F_2}) and integrating by time, we obtain (according to the initial condition)

$$\begin{aligned} & -\eta \iint_{\mathcal{Q}} \rho \frac{\partial P}{\partial t} \psi dx dt + \eta \int_{\Omega} \rho P(T) \psi(T) dx - \eta \int_{\Omega} \rho P(0) \beta_2 dx - i\eta \kappa \iint_{\mathcal{Q}} \operatorname{div}(\rho U_1) P \psi dx dt \\ & + i\eta \kappa \iint_{\mathcal{Q}} \rho w \nabla(\phi_1 P) dx dt - i\mu \iint_{\Sigma} \rho \psi P d\Gamma dt + \iint_{\mathcal{Q}} \rho \bar{B}(u)(P) B(u)(\psi) dx dt \\ & + \iint_{\mathcal{Q}} \rho \phi_1 w \bar{B}(u)(P) dx dt + \iint_{\mathcal{Q}} \rho B(u)(\phi_1) w P dx dt \\ & + \iint_{\mathcal{Q}} \rho((2|\varphi|^2 - 1)P\psi + \varphi^2 P \bar{\psi}) dx dt = 0, \\ & - \iint_{\mathcal{Q}} \rho \frac{\partial Q}{\partial t} w dx dt + \int_{\Omega} \rho Q(T) w(T) dx + \iint_{\mathcal{Q}} \rho \operatorname{curl}(Q) \operatorname{curl}(w) dx dt \\ & + \iint_{\mathcal{Q}} \operatorname{div}(Q) \operatorname{div}(\rho w) dx dt + \iint_{\mathcal{Q}} \rho \mathcal{R}(\bar{\psi} B(u)(\phi_1)) Q dx dt \\ & + \iint_{\mathcal{Q}} \rho(\mathcal{R}(\bar{\phi}_1 B(u)(\psi)) Q + w |\phi_1|^2 Q) dx dt = \iint_{\mathcal{Q}} \rho \beta_1 \operatorname{curl}(Q) dx dt. \end{aligned} \quad (5.23)$$

Since (P, Q) is solution of (5.8), we have that

$$\begin{aligned}
& i\eta\kappa \iint_{\mathcal{Q}} \rho w \nabla(\phi_1 P) dx dt + \iint_{\mathcal{Q}} \rho \phi_1 w \bar{B}(u)(P) dx dt + \iint_{\mathcal{Q}} \rho B(u)(\phi_1) w P dx dt \\
& + \iint_{\mathcal{Q}} \rho (\phi_1^2 P \bar{\psi} - \bar{\phi}_1^2 \bar{P} \psi) dx dt - \eta \int_{\Omega} \rho P(0) \beta_2 dx - \iint_{\mathcal{Q}} \rho \overline{B(u)(\phi_1)} Q \psi dx dt \\
& - \iint_{\mathcal{Q}} \left(-\frac{i}{\kappa} \operatorname{div}(\rho \bar{\phi}_1 Q) + \rho U_1 \bar{\phi}_1 Q \right) \psi dx dt + a \iint_{\mathcal{Q}} (|\varphi|^2 - \Lambda) \bar{\varphi} \psi dx dt = 0, \\
& \iint_{\mathcal{Q}} \rho \mathcal{R}(\bar{\psi} B(u)(\phi_1)) Q dx dt + \iint_{\mathcal{Q}} \rho \mathcal{R}(\bar{\phi}_1 B(u)(\psi)) Q dx dt \\
& - \iint_{\mathcal{Q}} \rho \mathcal{R}(i\eta\kappa \nabla(\phi_1 P) + P B(u)(\phi_1) + \phi_1 \bar{B}(u)(P)) w dx dt \\
& + b \iint_{\mathcal{Q}} (u - u_{\text{obs}}) w dx dt = \iint_{\mathcal{Q}} \rho \beta_1 \operatorname{curl}(Q) dx dt.
\end{aligned} \tag{5.24}$$

By adding the real part of the first part of (5.24) and the second part of (5.24), we obtain (by using Green's formula)

$$\begin{aligned}
& -\eta \int_{\Omega} \rho \mathcal{R}(P(0) \beta_2) dx + a \iint_{\mathcal{Q}} \mathcal{R}((|\varphi|^2 - \Lambda) \bar{\varphi} \psi) dx dt + b \iint_{\mathcal{Q}} (u - u_{\text{obs}}) w dx dt \\
& = \iint_{\mathcal{Q}} \rho \beta_1 \operatorname{curl}(Q) dx dt.
\end{aligned} \tag{5.25}$$

According to the expression of $J'(f, g) \cdot Y$, we can deduce that

$$J'(f, g) \cdot Y = \iint_{\mathcal{Q}} (\rho \operatorname{curl}(Q) + \alpha f) \beta_1 dx dt + \int_{\Omega} \mathcal{R}((\eta \rho P(0) - \gamma \bar{g}) \beta_2) dx. \tag{5.26}$$

Since (f^*, g^*) is an optimal solution, we have then

$$\begin{aligned}
& \iint_{\mathcal{Q}} (\rho \operatorname{curl}(Q^*) + \alpha f^*) (f^* - f) dx dt \geq 0, \\
& \int_{\Omega} \mathcal{R}((\eta \rho P^*(0) - \gamma \bar{g}^*) (g^* - g)) dx \leq 0, \quad \forall (f, g) \in \mathcal{H}.
\end{aligned} \tag{5.27}$$

This completes the proof. \square

Remark 5.10. In the case where the distributed disturbance is in the initial condition of the potential, we obtain the same results. In this case, the cost functional is given by

$$J(f, g) = \frac{a}{4} \| |\varphi|^2 - \Lambda \|_{L^2(\mathcal{Q})}^2 + \frac{b}{2} \| u - u_{\text{obs}} \|_{L^2(\mathcal{Q})}^2 + \frac{\alpha}{2} \| f \|_{L^2(\mathcal{Q})}^2 - \frac{\gamma}{2} \| g \|_{L^2}^2, \tag{5.28}$$

where $\alpha, \gamma > 0$, $a, b \geq 0$ and $a + b > 0$. The functions $(u_{\text{obs}}, \Lambda) \in \mathbf{L}^2(\mathcal{Q}) \times L^\infty(\mathcal{Q})$ represent the observation.

We can prove also an existence theorem of the robust control problem and obtain necessary optimality conditions for its solution using the same method.

Let $\mathcal{H} = \mathcal{H}_1 \times \mathcal{H}_2$ such that \mathcal{H}_1 and \mathcal{H}_2 are nonempty, closed, convex, bounded subsets of $L^2(\mathcal{Q})$ and $\mathbf{L}^2(\Omega)$, respectively.

For α and γ sufficiently large, there exists $(f^*, g^*, u, \varphi) \in \mathcal{H} \times \mathbf{E}_n \times \mathcal{E}$ satisfying (a.e. $t \in (0, T)$)

$$\begin{aligned}
& \eta \int_{\Omega} \rho \frac{\partial \varphi}{\partial t} q \, dx - i\eta\kappa \int_{\Omega} \operatorname{div}(\rho(u+U))\varphi q \, dx - i\eta\kappa \int_{\Omega} \operatorname{div}(\rho u)\varphi q \, dx - i\mu \int_{\Gamma} \rho \varphi q \, d\Gamma \\
& \quad + \int_{\Omega} \rho B(u)(\varphi + \phi)\bar{B}(u)(q) \, dx + \int_{\Omega} \rho F(\varphi)q \, dx = \int_{\Omega} \rho B(0)(\phi)\bar{B}(0)(q) \, dx, \\
& \int_{\Omega} \rho \frac{\partial u}{\partial t} v \, dx + \int_{\Omega} \rho \operatorname{curl}(u) \operatorname{curl}(v) \, dx + \int_{\Omega} \operatorname{div}(\rho u) \operatorname{div}(v) \, dx \\
& \quad + \int_{\Omega} \rho \mathcal{R}((\bar{\varphi} + \bar{\phi})B(u)(\varphi + \phi))v \, dx = \int_{\Omega} \rho \mathcal{R}(\bar{\phi}B(0)(\phi))v \, dx + \int_{\Omega} \rho f^* \operatorname{curl}(v) \, dx, \\
& (\varphi(0), u(0)) = (\varphi_0, g^*), \\
& \iint_{\mathcal{Q}} (\rho \operatorname{curl}(Q) + \alpha f^*)(f^* - f) \, dx \, dt \geq 0, \\
& \int_{\Omega} (\rho Q(0) - \gamma g^*)(g^* - g) \, dx \leq 0, \quad \forall (f, g) \in \mathcal{H},
\end{aligned} \tag{5.29}$$

where (P, Q) is solution of (5.8).

5.2. The control in the initial condition of the vector potential. In this section, we formulate the problem in two situations: firstly, the case where the worst disturbance is in the initial condition φ_0 , and secondly the case where the disturbance is in the external magnetic field h .

5.2.1. Distributed disturbance in the initial condition of the order parameter. We suppose that the control is in u_0 , that is, $u_0 = f$ ($f \in \mathbf{L}^2(\Omega)$) and the disturbance is in φ_0 , that is, $\varphi_0 = g$ ($g \in \mathcal{L}^2_{\infty}(\Omega)$). So the function (φ, u) is assumed to be related to the disturbance f and control g through the problem (3.3):

$$\begin{aligned}
& \eta \int_{\Omega} \rho \frac{\partial \varphi}{\partial t} q \, dx - i\eta\kappa \int_{\Omega} \operatorname{div}(\rho(u+U))\varphi q \, dx - i\eta\kappa \int_{\Omega} \operatorname{div}(\rho u)\varphi q \, dx - i\mu \int_{\Gamma} \rho \varphi q \, d\Gamma \\
& \quad + \int_{\Omega} \rho B(u)(\varphi + \phi)\bar{B}(u)(q) \, dx + \int_{\Omega} \rho F(\varphi)q \, dx \\
& = \int_{\Omega} \rho B(0)(\phi)\bar{B}(0)(q) \, dx, \quad \text{a.e. } t \in (0, T), \\
& \int_{\Omega} \rho \frac{\partial u}{\partial t} v \, dx + \int_{\Omega} \rho \operatorname{curl}(u) \operatorname{curl}(v) \, dx + \int_{\Omega} \operatorname{div}(\rho u) \operatorname{div}(v) \, dx \\
& \quad + \int_{\Omega} \rho \mathcal{R}((\bar{\varphi} + \bar{\phi})B(u)(\varphi + \phi))v \, dx \\
& = \int_{\Omega} \rho \mathcal{R}(\bar{\phi}B(0)(\phi))v \, dx + \int_{\Omega} \rho h \operatorname{curl}(v) \, dx, \quad \text{a.e. } t \in (0, T), \\
& (\varphi(0), u(0)) = (g, f).
\end{aligned} \tag{5.30}$$

To obtain the regularity of Theorem 4.2, we suppose that $h \in L^2(\mathcal{Q})$. Let $\mathcal{P} : (f, g) \rightarrow (\varphi, u) = \mathcal{P}(f, g)$ be the map: $\mathbf{L}^2(\Omega) \times \mathcal{L}_\infty^2(\Omega) \rightarrow \mathcal{E} \times \mathbf{E}_n$ defined by (5.30) and the cost function is defined by

$$J(f, g) = \frac{a}{4} \|\varphi\|^2 - \Lambda \|\varphi\|_{L^2(\mathcal{Q})}^2 + \frac{b}{2} \|u - u_{\text{obs}}\|_{L^2(\mathcal{Q})}^2 + \frac{\alpha}{2} \|f\|_{L^2}^2 - \frac{\beta}{2} \|g\|_{L^2}^2, \quad (5.31)$$

where $\alpha, \gamma > 0$, $a, b \geq 0$, and $a + b > 0$. The functions $(u_{\text{obs}}, \Lambda) \in \mathbf{L}^2(\mathcal{Q}) \times L^\infty(\mathcal{Q})$ are given.

In this section, we study the following minimax control problem (\mathcal{MP}_3) :

find an admissible control $f^* \in \mathcal{H}_1$ and a disturbance $g^* \in \mathcal{H}_2$ such that (f^*, g^*) is a saddle point of the functional J on \mathcal{H} , subject to system (5.30),

where $\mathcal{H} = \mathcal{H}_1 \times \mathcal{H}_2$ such that \mathcal{H}_1 and \mathcal{H}_2 are nonempty, closed, convex, bounded subsets of $\mathbf{L}^2(\Omega)$ and $\mathcal{L}_\infty^2(\Omega)$, respectively.

The arguments of Section 5.1 extend directly to the present case without further estimates, so we omit the details. We have then the following results

PROPOSITION 5.11. *The function \mathcal{P} is continuously Fréchet differentiable from $\mathbf{L}^2(\Omega) \times \mathcal{L}_\infty^2(\Omega)$ to $\mathcal{E} \times \mathbf{E}_n$ with the derivative $\mathcal{P}'(f, g) : Y = (\beta_1, \beta_2) \rightarrow (\psi, w)$ given by the linear problem (\mathcal{P}_{F_3}) (a.e. $t \in (0, T)$):*

$$\begin{aligned} & \eta \int_{\Omega} \rho \frac{\partial \psi}{\partial t} q \, dx - i\eta\kappa \int_{\Omega} \operatorname{div}(\rho U_1) \psi q \, dx - i\eta\kappa \int_{\Omega} \operatorname{div}(\rho w) \phi_1 q \, dx - i\mu \int_{\Gamma} \rho \psi q \, d\Gamma \\ & \quad + \int_{\Omega} \rho (B(u)(\psi) + \phi_1 w) \bar{B}(u)(q) \, dx + \int_{\Omega} \rho B(u)(\phi_1) w q \, dx \\ & \quad + \int_{\Omega} \rho ((2|\varphi|^2 - 1)\psi + \varphi^2 \bar{\psi}) q \, dx = 0, \\ & \int_{\Omega} \rho \frac{\partial w}{\partial t} v \, dx + \int_{\Omega} \rho \operatorname{curl}(w) \operatorname{curl}(v) \, dx + \int_{\Omega} \operatorname{div}(\rho w) \operatorname{div}(v) \, dx + \int_{\Omega} \rho \mathcal{R}(\bar{\psi} B(u)(\phi_1)) v \, dx \\ & \quad + \int_{\Omega} \rho (\mathcal{R}(\bar{\phi}_1 B(u)(\psi)) + w |\phi_1|^2) v \, dx = 0, \\ & (\psi(0), w(0)) = (\beta_2, \beta_1), \end{aligned} \quad (5.32)$$

where $(U_1, \phi_1) = (u + U, \varphi + \phi)$.

Moreover the following estimates $(\forall (f_i, g_i) \in \mathbf{L}^2(\Omega) \times \mathcal{L}_\infty^2(\Omega), i = 1, 2)$ exist:

- (i) $\|\mathcal{P}'(f_1, g_1)\|_{\mathcal{L}(\mathbf{L}^2(\Omega) \times \mathcal{L}_\infty^2(\Omega), \mathcal{E} \times \mathbf{E}_n)} \leq C_e$,
- (ii) $\|\mathcal{P}'(f_1, g_1)Y - \mathcal{P}'(f_2, g_2)Y\|_{\mathcal{E} \times \mathbf{E}_n}^2 \leq C_e (\|X\|_{L^2} \|Y\|_{L^2}^2 + \|X\|_{L^2}^2 \|Y\|_{L^2})$,

where $f = f_1 - f_2$, $g = g_1 - g_2$ and $X = (f, g)$.

PROPOSITION 5.12. *The map \mathcal{P} defined by (5.30) is continuous from the weak topology of $\mathbf{L}^2(\Omega) \times \mathcal{L}_\infty^2(\Omega)$ to the strong topology of $\mathcal{L}^2(\mathcal{Q}) \times \mathbf{L}^2(\mathcal{Q})$.*

THEOREM 5.13. *For α and γ sufficiently large (i.e., there exists (α_1, γ_1) such that $\alpha \geq \alpha_1$ and $\gamma \geq \gamma_1$), there exist $(f^*, g^*) \in \mathcal{H}$ and $(\varphi^*, u^*) \in \mathcal{E} \times \mathbf{E}_n$ such that (f^*, g^*) is defined by (\mathcal{MP}_3) and $(\varphi^*, u^*) = \mathcal{P}(f^*, g^*)$ is solution of (5.30).*

Next, we establish necessary optimality conditions for the robust control problem (\mathcal{MP}_3) .

THEOREM 5.14. *Under the assumptions of Theorem 3.1, the optimal solution $(f^*, g^*, u^*, \varphi^*) \in \mathcal{H} \times \mathbf{E}_n \times \mathcal{E}$, such that (f^*, g^*) is defined by (\mathcal{MP}_3) and $(\varphi^*, u^*) = \mathcal{P}(f^*, g^*)$ is solution of (5.30), satisfies*

$$\begin{aligned} \int_{\Omega} (\rho Q^*(0) + \alpha f^*)(f^* - f) dx &\geq 0, \\ \int_{\Omega} \mathcal{R}((\eta \rho P^*(0) - \gamma \overline{g^*})(g^* - g)) dx &\leq 0, \quad \forall (f, g) \in \mathcal{H}, \end{aligned} \quad (5.33)$$

where (P^*, Q^*) is solution of the adjoint problem (5.8) (corresponding to (φ^*, u^*)).

Proof. The cost function J is a composition of (Fréchet) differentiable maps, then J is differentiable and we have $(\forall Y = (\beta_1, \beta_2) \in \mathcal{H})$

$$\begin{aligned} J'(f, g) \cdot Y &= a \iint_{\mathcal{D}} \mathcal{R}(|\varphi|^2 - \Lambda) \overline{\varphi} \psi dx dt + b \iint_{\mathcal{D}} (u - u_{\text{obs}}) w dx dt \\ &+ \alpha \int_{\Omega} f \beta_1 dx - \gamma \int_{\Omega} \mathcal{R}(\overline{g} \beta_2) dx, \end{aligned} \quad (5.34)$$

where $(\psi, w) = \mathcal{P}'(f, g) \cdot Y$ is solution of problem (\mathcal{P}_{F_3}) .

By taking $(q, \nu) = (P, Q)$ in (\mathcal{P}_{F_3}) and integrating by time, we obtain (according to the homogeneous boundary conditions and to the initial condition)

$$\begin{aligned} &- \eta \iint_{\mathcal{D}} \rho \frac{\partial P}{\partial t} \psi dx dt + \int_{\Omega} \eta \rho P(T) \psi(T) dx - \int_{\Omega} \eta \rho P(0) \beta_2 dx - i \eta \kappa \iint_{\mathcal{D}} \text{div}(\rho U_1) P \psi dx dt \\ &+ i \eta \kappa \iint_{\mathcal{D}} \rho w \nabla(\phi_1 P) dx dt - i \mu \iint_{\Sigma} \rho \psi P d\Gamma dt + \iint_{\mathcal{D}} \rho \overline{B}(u)(P) B(u)(\psi) dx dt \\ &+ \iint_{\mathcal{D}} \rho \phi_1 w \overline{B}(u)(P) dx dt + \iint_{\mathcal{D}} \rho B(u)(\phi_1) w P dx dt \\ &+ \iint_{\mathcal{D}} \rho ((2|\phi_1|^2 - 1) P \psi + \phi_1^2 P \overline{\psi}) dx dt = 0, \\ &- \iint_{\mathcal{D}} \rho \frac{\partial Q}{\partial t} w dx dt + \int_{\Omega} \rho Q(T) w(T) dx - \int_{\Omega} \rho Q(0) \beta_1 dx + \iint_{\mathcal{D}} \rho \text{curl}(Q) \text{curl}(w) dx dt \\ &+ \iint_{\mathcal{D}} \text{div}(Q) \text{div}(\rho w) dx dt + \iint_{\mathcal{D}} \rho \mathcal{R}(\overline{\psi} B(u)(\phi_1)) Q dx dt \\ &+ \iint_{\mathcal{D}} \rho (\mathcal{R}(\overline{\phi_1} B(u)(\psi)) Q + w |\phi_1|^2 Q) dx dt = 0. \end{aligned} \quad (5.35)$$

Since (P, Q) is solution of (5.8), we have that

$$\begin{aligned}
& - \int_{\Omega} \eta \rho P(0) \beta_2 dx + i \eta \kappa \iint_{\mathcal{Q}} \rho w \nabla (\phi_1 P) dx dt + \iint_{\mathcal{Q}} \rho \phi_1 w \overline{B}(u)(P) dx dt \\
& \quad + \iint_{\mathcal{Q}} \rho B(u)(\phi_1) w P dx dt + \iint_{\mathcal{Q}} \rho (\phi_1^2 P \overline{\psi} - \overline{\phi_1^2} P \psi) dx dt - \iint_{\mathcal{Q}} \rho \overline{B(u)(\phi_1)} Q \psi dx dt \\
& \quad - \iint_{\mathcal{Q}} \left(-\frac{i}{\kappa} \operatorname{div}(\rho \overline{\phi_1} Q) + \rho U_1 \overline{\phi_1} Q \right) \psi dx dt + a \iint_{\mathcal{Q}} (|\varphi|^2 - \Lambda) \overline{\varphi} \psi dx dt = 0, \\
& - \int_{\Omega} \rho Q(0) \beta_1 dx + \iint_{\mathcal{Q}} \rho \Re(\overline{\psi} B(u)(\phi_1) + \overline{\phi_1} B(u)(\psi)) Q dx dt \\
& \quad - \iint_{\mathcal{Q}} \rho \Re(i \eta \kappa \nabla (\phi_1 P) + P B(u)(\phi_1) + \phi_1 \overline{B}(u)(P)) w dx dt \\
& \quad + b \iint_{\mathcal{Q}} (u - u_{\text{obs}}) w dx dt = 0.
\end{aligned} \tag{5.36}$$

By adding the real part of the first part of (5.36) and the second part of (5.36), we obtain

$$\begin{aligned}
& a \iint_{\mathcal{Q}} \Re((|\varphi|^2 - \Lambda) \overline{\varphi} \psi) dx dt + b \iint_{\mathcal{Q}} (u - u_{\text{obs}}) w dx dt \\
& \quad = \int_{\Omega} \rho Q(0) \beta_1 dx + \int_{\Omega} \eta \rho \Re(P(0) \beta_2) dx.
\end{aligned} \tag{5.37}$$

According to the expression of $J'(f, g) \cdot Y$, we can deduce that

$$J'(f, g) \cdot Y = \int_{\Omega} (\rho Q(0) + \alpha f) \beta_1 dx + \int_{\Omega} \Re((\eta \rho P(0) - \gamma g) \beta_2) dx. \tag{5.38}$$

Since (f^*, g^*) is an optimal solution we have then

$$\begin{aligned}
& \int_{\Omega} (\rho Q^*(0) + \alpha f^*) (f^* - f) dx \geq 0, \\
& \int_{\Omega} \Re((\eta \rho P^*(0) - \gamma g^*) (g^* - g)) dx \leq 0, \quad \forall (f, g) \in \mathcal{H}.
\end{aligned} \tag{5.39}$$

This completes the proof. \square

5.2.2. Distributed disturbance in the external magnetic field. In this section, the disturbance is in h and the control is in u_0 , that is, $u_0 = f$ ($f \in \mathbf{L}^2(\Omega)$), $h = g$ ($g \in L^2(\mathcal{Q})$). So the function (φ, u) is assumed to be related to the disturbance g and control f through the problem (3.3) (a.e. $t \in (0, T)$):

$$\begin{aligned}
& \eta \int_{\Omega} \rho \frac{\partial \varphi}{\partial t} q dx - i \eta \kappa \int_{\Omega} \operatorname{div}(\rho(u + U)) \varphi q dx - i \eta \kappa \int_{\Omega} \operatorname{div}(\rho u) \varphi q dx - i \mu \int_{\Gamma} \rho \varphi q d\Gamma \\
& \quad + \int_{\Omega} \rho B(u)(\varphi + \phi) \overline{B}(u)(q) dx + \int_{\Omega} \rho F(\varphi) q dx = \int_{\Omega} \rho B(0)(\phi) \overline{B}(0)(q) dx,
\end{aligned}$$

$$\begin{aligned}
& \int_{\Omega} \rho \frac{\partial u}{\partial t} v dx + \int_{\Omega} \rho \operatorname{curl}(u) \operatorname{curl}(v) dx + \int_{\Omega} \operatorname{div}(\rho u) \operatorname{div}(v) dx \\
& \quad + \int_{\Omega} \rho \mathcal{R}((\bar{\varphi} + \bar{\phi})B(u)(\varphi + \phi)) v dx = \int_{\Omega} \rho \mathcal{R}(\bar{\phi}B(0)(\phi)) v dx + \int_{\Omega} \rho g \operatorname{curl}(v) dx, \\
& (\varphi(0), u(0)) = (\varphi_0, f).
\end{aligned} \tag{5.40}$$

To obtain the regularity of Theorem 4.2, we suppose that $\varphi_0 \in \mathcal{L}_{\infty}^2(\Omega)$. Let $\mathcal{P} : (f, g) \rightarrow (\varphi, u) = \mathcal{P}(f, g)$ be the map: $\mathbf{L}^2(\Omega) \times L^2(\mathcal{Q}) \rightarrow \mathcal{E} \times \mathbf{E}_n$ defined by (5.40) and the cost function is defined by

$$J(f, g) = \frac{a}{4} \|\varphi\|^2 - \Lambda \|_{L^2(\mathcal{Q})}^2 + \frac{b}{2} \|u - u_{\text{obs}}\|_{L^2(\mathcal{Q})}^2 + \frac{\alpha}{2} \|f\|_{L^2}^2 - \frac{\gamma}{2} \|g\|_{L^2(\mathcal{Q})}^2, \tag{5.41}$$

where $\alpha, \gamma > 0$, $a, b \geq 0$ and $a + b > 0$. The functions $(u_{\text{obs}}, \Lambda) \in \mathbf{L}^2(\mathcal{Q}) \times L^{\infty}(\mathcal{Q})$ are given.

In this section, we study the following minimax control problem $(\mathcal{M}\mathcal{P}_4)$:

find an admissible control $f^* \in \mathcal{K}_1$ and a disturbance $g^* \in \mathcal{K}_2$ such that (f^*, g^*) is a saddle point of the functional J on \mathcal{K} , subject to system (5.40),

where, $\mathcal{K} = \mathcal{K}_1 \times \mathcal{K}_2$ such that \mathcal{K}_1 and \mathcal{K}_2 are nonempty, closed, convex, bounded subsets of $\mathbf{L}^2(\Omega)$ and $L^2(\mathcal{Q})$, respectively.

The proof of the following propositions and theorem of existence is obtained by using similar arguments of Section 5.1, so we omit the details.

PROPOSITION 5.15. *The function \mathcal{P} is continuously Fréchet differentiable from $\mathbf{L}^2(\Omega) \times L^2(\mathcal{Q})$ to $\mathcal{E} \times \mathbf{E}_n$ with the derivative $\mathcal{P}'(f, g) : Y = (\beta_1, \beta_2) \rightarrow (\psi, w)$ given by the linear problem (P_{F_4}) (a.e. $t \in (0, T)$):*

$$\begin{aligned}
& \eta \int_{\Omega} \rho \frac{\partial \psi}{\partial t} q dx - i\eta\kappa \int_{\Omega} \operatorname{div}(\rho U_1) \psi q dx - i\eta\kappa \int_{\Omega} \operatorname{div}(\rho w) \phi_1 q dx - i\mu \int_{\Gamma} \rho \psi q d\Gamma \\
& \quad + \int_{\Omega} \rho (B(u)(\psi) + \phi_1 w) \bar{B}(u)(q) dx + \int_{\Omega} \rho B(u)(\phi_1) w q dx \\
& \quad + \int_{\Omega} \rho ((2|\varphi|^2 - 1)\psi + \varphi^2 \bar{\psi}) q dx = 0, \\
& \int_{\Omega} \rho \frac{\partial w}{\partial t} u dx + \int_{\Omega} \rho \operatorname{curl}(w) \operatorname{curl}(v) dx + \int_{\Omega} \operatorname{div}(\rho w) \operatorname{div}(v) dx + \int_{\Omega} \rho \mathcal{R}(\bar{\psi}B(u)(\phi_1)) v dx \\
& \quad + \int_{\Omega} \rho (\mathcal{R}(\bar{\phi}_1 B(u)(\psi)) + w |\phi_1|^2) v dx = \int_{\Omega} \rho \beta_2 \operatorname{curl}(v) dx, \\
& (\psi(0), w(0)) = (0, \beta_1),
\end{aligned} \tag{5.42}$$

where $(U_1, \phi_1) = (u + U, \varphi + \phi)$.

Moreover the following estimates $(\forall (f_i, g_i) \in \mathbf{L}^2(\Omega) \times \mathbf{L}^2(\mathcal{Q}), i = 1, 2)$ exist:

$$(i) \|\mathcal{P}'(f_1, g_1)\|_{\mathcal{L}(\mathbf{L}^2 \times \mathbf{L}^2(\mathcal{Q}), \mathcal{E} \times \mathbf{E}_n)} \leq C_e,$$

$$(ii) \|\mathcal{P}'(f_1, g_1)Y - \mathcal{P}'(f_2, g_2)Y\|_{\mathcal{E} \times \mathbf{E}_n}^2 \leq C_e(\|X\|_{\mathbf{L}^2 \times \mathbf{L}^2(\mathcal{Q})}\|Y\|_{\mathbf{L}^2 \times \mathbf{L}^2(\mathcal{Q})}^2 + \|X\|_{\mathbf{L}^2 \times \mathbf{L}^2(\mathcal{Q})}\|Y\|_{\mathbf{L}^2 \times \mathbf{L}^2(\mathcal{Q})}),$$

where $f = f_1 - f_2$, $g = g_1 - g_2$ and $X = (f, g)$.

PROPOSITION 5.16. *The map \mathcal{P} defined by (5.40) is continuous from the weak topology of $\mathbf{L}^2(\Omega) \times \mathbf{L}^2(\mathcal{Q})$ to the strong topology of $\mathcal{L}^2(\mathcal{Q}) \times \mathbf{L}^2(\mathcal{Q})$.*

THEOREM 5.17. *For α and γ sufficiently large, there exist $(f^*, g^*) \in \mathcal{H}$ and $(u^*, \varphi^*) \in \mathbf{E}_n \times \mathcal{E}$ such that (f^*, g^*) is defined by $(\mathcal{M}\mathcal{P}_4)$ and $(\varphi^*, u^*) = \mathcal{P}(f^*, g^*)$ is solution of (5.40).*

Next we give necessary optimality conditions for the robust control problem $(\mathcal{M}\mathcal{P}_4)$.

THEOREM 5.18. *Under the assumptions of Theorem 5.8, the optimal solution $(f^*, g^*, u^*, \varphi^*) \in \mathcal{H} \times \mathbf{E}_n \times \mathcal{E}$, such that (f^*, g^*) is defined by $(\mathcal{M}\mathcal{P}_4)$ and $(\varphi^*, u^*) = \mathcal{P}(f^*, g^*)$ is solution of (5.40), satisfies*

$$\begin{aligned} \int_{\Omega} (\rho Q^*(0) + \alpha f^*)(f^* - f) dx &\geq 0, \\ \iint_{\mathcal{Q}} (\rho \operatorname{curl}(Q^*) - \gamma g^*)(g^* - g) dx dt &\leq 0, \quad \forall (f, g) \in \mathcal{H}, \end{aligned} \quad (5.43)$$

where (P^*, Q^*) is solution of the adjoint problem (5.8) (corresponding to (φ^*, u^*)).

Proof. The cost function J is a composition of (Fréchet) differentiable maps, then J is differentiable and we have $(\forall Y = (\beta_1, \beta_2) \in \mathcal{H})$

$$\begin{aligned} J'(f, g) \cdot Y &= a \iint_{\mathcal{Q}} \mathcal{R}(|\phi|^2 - \Lambda) \bar{\varphi} \psi dx dt + b \iint_{\mathcal{Q}} (u - u_{\text{obs}}) w dx dt \\ &+ \alpha \int_{\Omega} f \beta_1 dx - \gamma \iint_{\mathcal{Q}} g \beta_2 dx dt, \end{aligned} \quad (5.44)$$

where $(\psi, w) = \mathcal{P}'(f, g) \cdot Y$ is solution of problem (\mathcal{P}_{F_4}) .

By taking $(q, v) = (P, Q)$ in (\mathcal{P}_{F_4}) , using Green's formula, and integrating by time, we obtain (according to the homogeneous boundary conditions and to the initial condition)

$$\begin{aligned} & - \eta \iint_{\mathcal{Q}} \rho \frac{\partial P}{\partial t} \psi dx dt + \eta \int_{\Omega} \rho P(T) \psi(T) dx \\ & - i\eta\kappa \iint_{\mathcal{Q}} \operatorname{div}(\rho U_1) P \psi dx dt + i\eta\kappa \iint_{\mathcal{Q}} \rho w \nabla(\phi_1 P) dx dt \\ & - i\mu \int_{\Gamma} \rho \psi P d\Gamma + \iint_{\mathcal{Q}} \rho \bar{B}(u)(P) B(u)(\psi) dx dt + \iint_{\mathcal{Q}} \rho \phi_1 w \bar{B}(u)(P) dx dt \\ & + \iint_{\mathcal{Q}} \rho B(u)(\phi_1) w P dx dt + \iint_{\mathcal{Q}} \rho ((2|\phi_1|^2 - 1) P \psi + \phi_1^2 P \bar{\psi}) dx dt = 0, \end{aligned}$$

$$\begin{aligned}
& - \iint_{\mathcal{Q}} \rho \frac{\partial Q}{\partial t} w \, dx \, dt + \int_{\Omega} \rho Q(T) w(T) \, dx - \int_{\Omega} \rho Q(0) \beta_1 \, dx + \iint_{\mathcal{Q}} \rho \operatorname{curl}(Q) \operatorname{curl}(w) \, dx \, dt \\
& + \iint_{\mathcal{Q}} \operatorname{div}(Q) \operatorname{div}(\rho w) \, dx \, dt + \iint_{\mathcal{Q}} \rho \mathcal{R}(\overline{\psi} B(u)(\phi_1)) Q \, dx \, dt \\
& + \iint_{\mathcal{Q}} \rho (\mathcal{R}(\overline{\phi_1} B(u)(\psi)) Q + w |\phi_1|^2 Q) \, dx \, dt = \iint_{\mathcal{Q}} \rho \beta_2 \operatorname{curl}(Q) \, dx \, dt. \quad (5.45)
\end{aligned}$$

Since (P, Q) is solution of (5.8), we have that

$$\begin{aligned}
& i\eta\kappa \iint_{\mathcal{Q}} \rho w \nabla(\phi_1 P) \, dx \, dt + \iint_{\mathcal{Q}} \rho \phi_1 w \overline{B}(u)(P) \, dx \, dt + \iint_{\mathcal{Q}} \rho B(u)(\phi_1) w P \, dx \, dt \\
& + \iint_{\mathcal{Q}} \rho (\phi_1^2 P \overline{\psi} - \overline{\phi_1}^2 \overline{P} \psi) \, dx \, dt - \iint_{\mathcal{Q}} \rho \overline{B}(u)(\phi_1) Q \psi \, dx \, dt \\
& - \iint_{\mathcal{Q}} \left(-\frac{i}{\kappa} \operatorname{div}(\rho \overline{\phi_1} Q) + \rho U_1 \overline{\phi_1} Q \right) \psi \, dx \, dt + a \iint_{\mathcal{Q}} (|\phi|^2 - \Lambda) \overline{\varphi} \psi \, dx \, dt = 0, \\
& - \int_{\Omega} \rho Q(0) \beta_1 \, dx + \iint_{\mathcal{Q}} \rho \mathcal{R}(\overline{\psi} B(u)(\phi_1)) Q \, dx \, dt + \iint_{\mathcal{Q}} \rho \mathcal{R}(\overline{\phi_1} B(u)(\psi)) Q \, dx \, dt \\
& - \iint_{\mathcal{Q}} \rho \mathcal{R}(i\eta\kappa \nabla(\phi_1 P)) \, dx \, dt + P B(u)(\phi_1) + \phi_1 \overline{B}(u)(P) w \, dx \, dt \\
& + b \iint_{\mathcal{Q}} (u - u_{\text{obs}}) w \, dx \, dt = \iint_{\mathcal{Q}} \rho \beta_2 \operatorname{curl}(Q) \, dx \, dt. \quad (5.46)
\end{aligned}$$

By adding the real part of the first part of (5.46) and the second part of (5.46), we obtain

$$\begin{aligned}
& - \int_{\Omega} \rho Q(0) \beta_1 \, dx + a \iint_{\mathcal{Q}} \mathcal{R}((|\phi|^2 - \Lambda) \overline{\varphi} \psi) \, dx \, dt + b \iint_{\mathcal{Q}} (u - u_{\text{obs}}) w \, dx \, dt \\
& = \iint_{\mathcal{Q}} \rho \beta_2 \operatorname{curl}(Q) \, dx \, dt. \quad (5.47)
\end{aligned}$$

According to the expression of $J'(f, g) \cdot Y$, we can deduce that

$$J'(f, g) \cdot Y = \int_{\Omega} (\rho Q(0) + \alpha f) \beta_1 \, dx + \iint_{\mathcal{Q}} (\rho \operatorname{curl}(Q) - \gamma g) \beta_2 \, dx \, dt. \quad (5.48)$$

Since (f^*, g^*) is an optimal solution, we obtain

$$\begin{aligned}
& \int_{\Omega} (\rho Q^*(0) + \alpha f^*) (f^* - f) \, dx \geq 0, \\
& \iint_{\mathcal{Q}} (\rho \operatorname{curl}(Q^*) - \gamma g^*) (g^* - g) \, dx \, dt \leq 0, \quad \forall (f, g) \in \mathcal{H}. \quad (5.49)
\end{aligned}$$

This completes the proof. \square

Remark 5.19. In the case where the distributed disturbance is in the initial condition of the potential, we obtain the same results. In this case the cost functional is given by

$$J(f, g) = \frac{a}{4} \|\varphi\|^2 - \Lambda \| \cdot \|_{L^2(\mathcal{Q})}^2 + \frac{b}{2} \|u - u_{\text{obs}}\|_{L^2(\mathcal{Q})}^2 + \frac{\alpha}{2} \|f\|_{L^2}^2 - \frac{\gamma}{2} \|g\|_{L^2}^2, \quad (5.50)$$

where $\alpha, \gamma > 0$, $a, b \geq 0$ and $a + b > 0$.

We can prove also an existence theorem of the robust control problem and obtain necessary optimality conditions for its solution using the same method. Let $\mathcal{H} = \mathcal{H}_1 \times \mathcal{H}_2$ such that \mathcal{H}_1 and \mathcal{H}_2 are nonempty, closed, convex, bounded subsets of $\mathbf{L}^2(\Omega)$.

For α and γ sufficiently large, there exists (f^*, g^*, φ, u) satisfying

$$\begin{aligned} & \eta \int_{\Omega} \rho \frac{\partial \varphi}{\partial t} q dx - i\eta\kappa \int_{\Omega} \operatorname{div}(\rho(u+U))\varphi q dx - i\eta\kappa \int_{\Omega} \operatorname{div}(\rho u)\phi q dx - i\mu \int_{\Gamma} \rho \varphi P d\Gamma \\ & \quad + \int_{\Omega} \rho B(u)(\varphi + \phi)\bar{B}(u)(q) dx + \int_{\Omega} \rho F(\varphi)q dx \\ & = \int_{\Omega} \rho B(0)(\phi)\bar{B}(0)(q) dx, \quad \text{a.e. } t \in (0, T) \\ & \int_{\Omega} \rho \frac{\partial u}{\partial t} v dx + \int_{\Omega} \rho \operatorname{curl}(u) \operatorname{curl}(v) dx + \int_{\Omega} \operatorname{div}(\rho u) \operatorname{div}(v) dx \\ & \quad + \int_{\Omega} \rho \mathcal{R}((\bar{\varphi} + \bar{\phi})B(u)(\varphi + \phi))v dx \\ & = \int_{\Omega} \rho \mathcal{R}(\bar{\phi}B(0)(\phi))v dx dt + \int_{\Omega} \rho h \operatorname{curl}(v) dx, \quad \text{a.e. } t \in (0, T) \\ & (\varphi(0), u(0)) = (\varphi_0, f^* + g^*), \\ & \int_{\Omega} (\rho Q(0) + \alpha f^*)(f^* - f) dx \geq 0, \\ & \int_{\Omega} (\rho Q(0) - \gamma g^*)(g^* - g) dx \leq 0, \quad \forall (f, g) \in \mathcal{H}, \end{aligned} \quad (5.51)$$

where (P, Q) is solution of (5.8).

6. Conclusion

We have developed a robust control method for the time-dependent complex Ginzburg-Landau vortices in superconductivity. This model contains two unknowns, the vector potential u and an order parameter φ (or a complex phase-field variable) coming from thermodynamics, where φ describes the phase of the underlying superconductivity. The case $|\varphi|$ close to 1 corresponds to the superconducting phase, $|\varphi|$ close to 0 to the normal phase. The existence of weak solution as well as regularity and stability results are established. A robust control problem has been studied for two sets of distributed controls: firstly the control is in the external magnetic field (the disturbance is in the external magnetic field or in the initial condition of the parameter), secondly the control is in the initial condition of the vector potential (the disturbance is in the external magnetic field or in the initial condition of the parameter). Under suitable hypotheses, it is shown that

one has existence of solution to a corresponding robust control problem, and the appropriate necessary conditions for saddle point optimality are obtained. These conditions are corresponding to identify the gradient of the cost functional that is very useful in the numerical resolution of the robust control problem. For example we can combining the optimal necessary conditions obtained in this paper and the gradient-iterative algorithm to solve the robust control problem numerically (at each iteration i , we obtain the numerical approximation (f_i, g_i) of the optimal solution (f, g) by $f_i = f_{i-1} - \theta J'_f(f_{i-1}, g_{i-1})$ and $g_i = g_{i-1} + \omega J'_g(f_{i-1}, g_{i-1})$, where $0 < m \leq \theta$, $\omega \leq M$ are the sequences of step lengths and m, M depending on the second Fréchet derivative of J to ensure the convergence result (see, e.g., Ciarlet [16])).

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