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## ON THE ANTIMAXIMUM PRINCIPLE FOR PARABOLIC PERIODIC PROBLEMS WITH WEIGHT


#### Abstract

We prove that an antimaximum principle holds for the Neumann and Dirichlet periodic parabolic linear problems of second order with a time periodic and essentially bounded weight function. We also prove that an uniform antimaximum principle holds for the one dimensional Neumann problem which extends the corresponding elliptic case.


## 1. Introduction

Let $\Omega$ be a bounded domain in $\mathbb{R}^{N}$ with $C^{2+\gamma}$ boundary, $0<\gamma<1$ and let $\tau>0$. Let $\left\{a_{i, j}(x, t)\right\}_{1 \leq i, j \leq N}\left\{b_{j}(x, t)\right\}_{1 \leq j \leq N}$ and $a_{0}(x, t)$ be $\tau$-periodic functions in $t$ such that $a_{i, j}, b_{j}$ and $a_{0}$ belong to $C^{\gamma, \gamma / 2}(\bar{\Omega} \times \mathbb{R}), a_{i, j}=a_{j, i}$ for $1 \leq i, j \leq N$ and $\sum_{i, j} a_{i, j}(x, t) \xi_{i} \xi_{j} \geq c \sum_{i} \xi_{i}^{2}$ for some $c>0$ and all $(x, t) \in \bar{\Omega} \times \mathbb{R},\left(\xi_{1}, \ldots, \xi_{N}\right) \in \mathbb{R}^{N}$. Let $L$ be the periodic parabolic operator given by

$$
\begin{equation*}
L u=\frac{\partial u}{\partial t}-\sum_{i, j} a_{i, j} \frac{\partial^{2} u}{\partial x_{i} \partial x_{j}}+\sum_{j} b_{j} \frac{\partial u}{\partial x_{j}}+a_{0} u \tag{1}
\end{equation*}
$$

Let $B(u)=0$ denote either the Dirichlet boundary condition $u_{\mid \partial \Omega \times \mathbb{R}}=0$ or the Neumann condition $\partial u / \partial v=0$ along $\partial \Omega \times \mathbb{R}$.

Let us consider the problem

$$
\left(P_{\lambda, h}\right) \quad\left\{\begin{array}{c}
L u=\lambda m u+h \text { in } \Omega \times \mathbb{R}, \\
u \quad \tau-\text { periodic in } t \\
B(u)=0 \text { on } \partial \Omega \times \mathbb{R}
\end{array}\right.
$$

where the weight function $m=m(x, t)$ is a $\tau$ - periodic and essentially bounded function, $h=h(x, t)$ is $\tau$-periodic in $t$ and $h \in L^{p}(\Omega \times(0, \tau))$ for some $p>N+2$.

We say that $\lambda^{*} \in \mathbb{R}$ is a principal eigenvalue for the weight $m$ if $\left(P_{\lambda^{*}, h}\right)$ has a positive solution when $h \equiv 0$. The antimaximum principle can be stated as follows:

DEFINITION 1. We will say that the antimaximum principle (AMP) holds to the right (respectively to the left) of a principal eigenvalue $\lambda^{*}$ iffor each $h \geq 0, h \neq 0$ (with

[^0]$h \in L^{p}(\Omega \times(0, \tau))$ for some $\left.p>N+2\right)$ there exists a $\delta(h)>0$ such that $\left(P_{\lambda, h}\right)$ has a negative solution for each $\lambda \in\left(\lambda^{*}, \lambda^{*}+\delta(h)\right)$ (respectively $\lambda \in\left(\lambda^{*}-\delta(h), \lambda^{*}\right)$ ).

We prove that, depending on $m$, these two possibilities happen and that in some cases the AMP holds left and right of $\lambda^{*}$, similarly to the purely stationary case where all data are independent of $t$ (but in that case the period becomes artificial)

Our results are described by means of the real function $\mu_{m}(\lambda), \lambda \in \mathbb{R}$, defined as the unique $\mu \in \mathbb{R}$ such that the homogeneous problem

$$
\left(P_{\mu}\right) \quad\left\{\begin{array}{c}
L u-\lambda m u=\mu u \text { in } \Omega \times \mathbb{R}, \\
u \quad \tau-\text { periodic in } t \\
B(u)=0 \text { on } \partial \Omega \times \mathbb{R}
\end{array}\right.
$$

has a positive solution.
This function was first studied by Beltramo - Hess in [2] for Hölder continuous weight and Dirichlet boundary condition. They proved that $\mu_{m}$ is a concave and real analytic function, for Neumann the same holds ([8], Lemmas 15.1 and 15.2). A given $\lambda \in \mathbb{R}$ is a principal eigenvalue for the weight $m$ iff $\mu_{m}$ has a zero at $\lambda$

We will prove that if $\mu_{m}$ is non constant and if $\lambda^{*}$ is a principal eigenvalue for the weight $m$ then the AMP holds to the left of $\lambda^{*}$ if $\mu_{m}^{\prime}\left(\lambda^{*}\right)>0$, holds to the right of $\lambda^{*}$ if $\mu_{m}^{\prime}\left(\lambda^{*}\right)<0$ and holds right and left of $\lambda^{*}$ if $\mu_{m}^{\prime}\left(\lambda^{*}\right)=0$. As a consequence of these results we will give (see section 3 , Theorem 1), for the case $a_{0} \geq 0$, conditions on $m$ that describe completely what happens respect to the AMP, near each principal eigenvalue.

The notion of AMP is due to Ph. Clement and L. A. Peletier [3]. They proved an AMP to the right of the first eigenvalue for $m=1$, with all data independent of $t$ and $a_{0}(x) \geq 0$, i.e. the elliptic case. Hess [7] proves the same, in the Dirichlet case, for $m \in C(\bar{\Omega})$. Our aim is to extend these results to periodic parabolic problems covering both cases, Neumann and Dirichlet. In section 2 we give a version of the AMP for a compact family of positive operators adapted to our problem and in section 3 we state the main results.

## 2. Preliminaries

Let $Y$ be an ordered real Banach space with a total positive cone $P_{Y}$ with norm preserving order, i.e. $u, v \in Y, 0<u \leq v$ implies $\|u\| \leq\|v\|$. Let $P_{Y}^{\circ}$ denote the interior of $P_{Y}$ in $Y$. We will assume, from now on, that $P_{Y}^{\circ} \neq \emptyset$. Its dual $Y^{\prime}$ is an ordered Banach space with positive cone

$$
P^{\prime}=\left\{y^{\prime} \in Y^{\prime}:\left\langle y^{\prime}, y\right\rangle \geq 0 \text { for all } y \in P\right\}
$$

For $y^{\prime} \in Y^{\prime}$ we set $y^{\perp}=\left\{y \in Y:\left\langle y^{\prime}, y\right\rangle=0\right\}$ and for $r>0, B_{r}^{Y}(y)$ will denote the open ball in $Y$ centered at $y$ with radius $r$. For $v, w \in Y$ with $v<w$ we put $(v, w)$ and $[u, v]$ for the order intervals $\{y \in Y: v<y<w\}$ and $\{y \in Y: v \leq y \leq w\}$
respectively. $B(Y)$ will denote the space of the bounded linear operators on $Y$ and for $T \in B(Y), T^{*}$ will be denote its adjoint $T^{*}: Y^{\prime} \rightarrow Y^{\prime}$.

Let us recall that if $T$ is a compact and strongly positive operator on $Y$ and if $\rho$ is its spectral radius, then, from Krein - Rutman Theorem, (as stated, e.g., in [1], Theorem 3.1), $\rho$ is a positive algebraically simple eigenvalue with positive eigenvectors associated for $T$ and for its adjoint $T^{*}$.

We will also need the following result due to Crandall - Rabinowitz ([4], Lemma 1.3 ) about perturbation of simple eigenvalues:

LEMMA 1. If $T_{0}$ is a bounded operator on $Y$ and if $r_{0}$ is an algebraically simple eigenvalue for $T_{0}$, then there exists $\delta>0$ such that $\left\|T-T_{0}\right\|<\delta$ implies that there exists a unique $r(T) \in R$ satisfying $\left|r(T)-r_{0}\right|<\delta$ for which $r(T) I-T$ is singular. Moreover, the map $T \rightarrow r(T)$ is analytic and $r(T)$ is an algebraically simple eigenvalue for $T$. Finally, an associated eigenvector $v(T)$ can be chosen such that the map $T \rightarrow v(T)$ is also analytic.

We start with an abstract formulation of the AMP for a compact family of operators. The proof is an adaptation, to our setting, of those in [3] and [7].

LEMMA 2. Let $\left\{T_{\lambda}\right\}_{\lambda \in \Lambda}$ be a compact family of compact and strongly positive operators on $Y$. Denote by $\rho(\lambda)$ the spectral radius of $T_{\lambda}$ and $\sigma(\lambda)$ its spectrum. Then for all $0<u \leq v$ in $Y$ there exists $\delta_{u, v}>0$ in $\mathbb{R}$ such that

$$
\left(\rho(\lambda)-\delta_{u, v}, \rho(\lambda)\right) \cap \sigma(\lambda)=\emptyset \text { and }\left(\theta I-T_{\lambda}\right)^{-1} h<0
$$

uniformly in $h \in[u, v] \subset Y$ and $\theta \in\left(\rho(\lambda)-\delta_{u, v}, \rho(\lambda)\right) \subset \mathbb{R}$.
Proof. We have that
(1) $\rho(\lambda)$ is an algebraically simple eigenvalue for $T_{\lambda}$ with a positive eigenvector $\Phi_{\lambda}$
(2) $T_{\lambda}^{*}$ has an eigenvector $\Psi_{\lambda}$ associated to the eigenvalue $\rho(\lambda)$ such that $\left.\left\langle\Psi_{\lambda}, x\right\rangle\right\rangle$ 0 for all $x \in P-\{0\}$.
(3) $\Phi_{\lambda}$ normalized by $\left\|\Phi_{\lambda}\right\|=1$ and $\Psi_{\lambda}$ normalized by $\left\langle\Psi_{\lambda}, \Phi_{\lambda}\right\rangle=1$ imply that $\left\{\left(\Phi_{\lambda}, \Psi_{\lambda}, \rho(\lambda)\right) \in Y \times Y^{\prime} \times \mathbb{R}\right\}$ is compact.
(4) There exists $r>0$ such that $B_{r}^{Y}(0) \subset\left(-\Phi_{\lambda}, \Phi_{\lambda}\right)$ for all $\lambda \in \Lambda$.
(1) and (2) follow from Krein - Rutman Theorem.

For (3) $\{\rho(\lambda), \lambda \in \Lambda\}$ is compact in $(0, \infty)$ because given $\rho\left(\lambda_{n}\right)$, the sequence $T_{\lambda_{n}}$ has a subsequence (still denoted) $T_{\lambda_{n}} \rightarrow T_{\lambda_{\infty}} \in\left\{T_{\lambda}\right\}_{\lambda \in \Lambda}$ in $B(Y)$. Taking into account (1), Lemma 1 provides an $r>0$ such that $T \in B(Y),\left\|T-T_{\infty}\right\|<r$ imply that $0 \notin\left(\rho\left(\lambda_{\infty}\right)-r, \rho\left(\lambda_{\infty}\right)+r\right) \cap \sigma(T)=\{\rho(T)\}$, so $\rho\left(\lambda_{n}\right) \rightarrow \rho\left(\lambda_{\infty}\right)>0$. This lemma also gives $\left\{\Phi_{\lambda}: \lambda \in \Lambda\right\}$ and $\left\{\Psi_{\lambda}: \lambda \in \Lambda\right\}$ compact in $Y$ and $Y^{\prime}$ respectively.
(4) follows remarking that $\left\{\Phi_{\lambda}, \lambda \in \Lambda\right\}$ has a lower bound $v \leq \Phi_{\lambda}, v \in P_{Y}^{\circ}$. Indeed, $\frac{1}{2} \Phi_{\lambda} \in P_{Y}^{\circ}$ so $w-\frac{1}{2} \Phi_{\lambda}=\frac{1}{2} \Phi_{\lambda}+w-\Phi_{\lambda} \in P_{Y}^{\circ}$ for $w \in B_{r(\lambda)}\left(\Phi_{\lambda}\right)$ with
$r(\lambda)>0$. The open covering $\left\{B_{r(\lambda)}\left(\Phi_{\lambda}\right)\right\}$ of $\left\{\Phi_{\lambda}: \lambda \in \Lambda\right\}$ admits a finite subcovering $\left\{B_{r\left(\lambda_{j}\right)}\left(\Phi_{\lambda_{j}}\right), j=1,2, \ldots, l\right\}$ and it is simple to obtain $r_{j} \in(0,1)$ such that $\Phi_{\lambda_{j}}>$ $r_{j} \frac{1}{2} \Phi_{\lambda_{1}} j=1,2, \ldots, l$, so $v=r \Phi_{\lambda_{1}} \leq \frac{1}{2} \Phi_{\lambda_{j}}<\Phi_{\lambda}$ for all $\lambda \in \Lambda$ and some $j(j$ depending on $\lambda$ ).

We prove now the Lemma for each $\lambda$ and $h \in[u, v]$, i.e. we find $\delta_{u, v}(\lambda)$ and we finish by a compactness argument thanks to (1)-(4).
$\Psi_{\lambda}^{\perp}$ is a closed subspace of $Y$ and then, endowed with the norm induced from $Y$, it is a Banach space. It is clear that $\Psi_{\lambda}^{\perp}$ is $T_{\lambda}$ invariant. and that $T_{\lambda \mid \Psi_{\lambda}}: \Psi_{\lambda}^{\perp} \rightarrow \Psi_{\lambda}^{\perp}$ is a compact operator. Now, $\rho(\lambda)$ is a simple eigenvalue for $T_{\lambda}$ with eigenvector $\Phi_{\lambda}$ and $\Phi_{\lambda} \notin \Psi_{\lambda}^{\perp}$, but $\rho(\lambda)>0$ and $T_{\lambda}$ is a compact operator, thus $\rho(\lambda) \notin \sigma\left(T_{\lambda \mid \Psi_{\lambda}^{\perp}}\right)$.

We have also $Y=\mathbb{R} \Phi_{\lambda} \bigoplus \Psi_{\lambda}^{\perp}$, a direct sum decomposition with bounded projections $P_{\Phi_{\lambda}}, P_{\Psi_{\lambda}^{\perp}}$ given by $P_{\Phi_{\lambda}} y=\left\langle\Psi_{\lambda}, y\right\rangle \Phi_{\lambda}$ and $P_{\Psi_{\lambda}^{\perp}} y=y-\left\langle\Psi_{\lambda}, y\right\rangle \Phi_{\lambda}$ respectively. Let $\widetilde{T}_{\lambda}: Y \rightarrow Y$ be defined by $\widetilde{T}_{\lambda}=T_{\lambda} P_{\Psi_{\lambda}^{\perp}}$ Thus $\widetilde{T}_{\lambda}$ is a compact operator. Moreover, $\rho(\lambda)$ does not belongs to its spectrum. (Indeed, suppose that $\rho(\lambda)$ is an eigenvalue for $\widetilde{T}_{\lambda}$, let $v$ be an associated eigenvector. We write $v=P_{\Phi_{\lambda}} v+P_{\Psi_{\lambda}^{\perp}} v$. Then $\rho(\lambda) v=\widetilde{T}_{\lambda}(v)=T_{\lambda} P_{\Psi_{\lambda}^{\perp}} v$ and so $v \in \Psi_{\lambda}^{\perp}$, but $\rho(\lambda) \notin \sigma\left(T_{\lambda \mid \Psi_{\lambda}^{\perp}}\right)$. Contradiction). Thus, for each $\lambda, \rho(\lambda) I-\widetilde{T}_{\lambda}$ has a bounded inverse.

Hence, from the compactness of the set $\{\rho(\lambda): \lambda \in \Lambda\}$ it follows that there exists $\varepsilon>0$ such that $\theta I-\widetilde{T}_{\lambda}$ has a bounded inverse for $(\theta, \lambda) \in D$ where

$$
D=\{(\theta, \lambda): \lambda \in \Lambda, \rho(\lambda)-\varepsilon \leq \theta \leq \rho(\lambda)+\varepsilon\}
$$

and that $\left\|\left(\theta I-\widetilde{T}_{\lambda}\right)^{-1}\right\|_{B(Y)}$ remains bounded as $(\lambda, \theta)$ runs on $D$.
But $\theta I-T_{\lambda \mid \Psi_{\lambda}^{\perp}}: \Psi_{\lambda}^{\perp} \rightarrow \Psi_{\lambda}^{\perp}$ has a bounded inverse given by $\left(\theta I-T_{\lambda \mid \Psi_{\lambda}^{\perp}}\right)^{-1}=$ $\left(\left(\theta I-\widetilde{T}_{\lambda}\right)^{-1}\right)_{\mid \Psi_{\lambda}^{\perp}}$ and so $\left\|\left(\theta I-T_{\lambda \mid \Psi_{\lambda}^{\perp}}\right)^{-1}\right\|_{B\left(\Psi_{\lambda}^{\perp}\right)}$ remains bounded as $(\lambda, \theta)$ runs on $D$.

For $h \in[u, v]$ we set $w_{\lambda, h}=h-\left\langle\Psi_{\lambda}, h\right\rangle \Phi_{\lambda}$. As $\theta \notin \sigma\left(T_{\lambda}\right)$ we have

$$
\begin{equation*}
\left(\theta I-T_{\lambda}\right)^{-1} h=\frac{\left\langle\Psi_{\lambda}, h\right\rangle}{\theta-\rho(\lambda)}\left[\Phi_{\lambda}+\frac{\theta-\rho(\lambda)}{\left\langle\Psi_{\lambda}, h\right\rangle}\left(\left(\theta I-T_{\lambda}\right)_{\mid \Psi_{\lambda}^{\perp}}\right)^{-1} w_{h, \lambda}\right] \tag{2}
\end{equation*}
$$

and $u-\left\langle\Psi_{\lambda}, v\right\rangle \Phi_{\lambda} \leq w_{h, \lambda} \leq v-\left\langle\Psi_{\lambda}, u\right\rangle \Phi_{\lambda}$ that is $\left\|w_{h, \lambda}\right\| \leq c_{u, v}$ for some constant $c_{u, v}$ independent of $h$. Hence $\left(\left(\theta I-T_{\lambda}\right)_{\mid \Psi_{\lambda}^{\perp}}\right)^{-1} w_{h, \lambda}$ remains bounded in $Y$, uniformly on $(\theta, \lambda) \in D$ and $h \in[u, v]$. Also, $\left\langle\Psi_{\lambda}, h\right\rangle \geq\left\langle\Psi_{\lambda}, u\right\rangle$ and since $\left\{\left\langle\Psi_{\lambda}, u\right\rangle\right\}$ is compact in $(0, \infty)$ it follows that $\left\langle\Psi_{\lambda}, h\right\rangle \geq c$ for some positive constant $c$ and all $\lambda \in \Lambda$ and all $h \in[u, v]$. Thus the lemma follows from (4).

REMARK 1. The conclusion of Lemma 2 holds if (1)-(4) are fulfilled.

We will use the following
Corollary 1. Let $\lambda \rightarrow T_{\lambda}$ be continuous map from $[a, b] \subset \mathbb{R}$ into $B(Y)$. If each $T_{\lambda}$ is a compact and strongly positive operator then the conclusion of Lemma 2 holds.

## 3. The AMP for periodic parabolic problems

For $1 \leq p \leq \infty$, denote $X=L_{\tau}^{p}(\Omega \times \mathbb{R})$ the space of the $\tau$ - periodic functions $u: \Omega \times \mathbb{R} \rightarrow \mathbb{R}$ (i.e. $u(x, t)=u(x, t+\tau)$ a.e. $(x, t) \in \Omega \times \mathbb{R})$ such whose restrictions to $\Omega \times(0, \tau)$ belong to $L^{p}(\Omega \times(0, \tau))$. We write also $C_{\tau, B}^{1+\gamma, \gamma}(\bar{\Omega} \times \mathbb{R})$ for the space of the $\tau$-periodic Hölder continuous functions $u$ on $\bar{\Omega} \times \mathbb{R}$ satisfying the boundary condition $B(u)=0$ and $C_{\tau}(\bar{\Omega} \times \mathbb{R})$ for the space of $\tau$ - periodic continuous functions on $\bar{\Omega} \times \mathbb{R}$. We set

$$
\begin{array}{ll}
Y=C_{\tau, B}^{1+\gamma, \gamma}(\bar{\Omega} \times \mathbb{R}) & \text { if } B(u)=u_{\mid \partial \Omega \times \mathbb{R}} \\
\text { and }  \tag{3}\\
Y=C_{\tau}(\bar{\Omega} \times \mathbb{R}) & \text { if } B(u)=\partial u / \partial \nu
\end{array}
$$

In each case, $X$ and $Y$, equipped with their natural orders and norms are ordered Banach spaces, and in the first one, $Y$ has compact inclusion into $X$ and the cone $P_{Y}$ of the positive elements in $Y$ has non empty interior.

Fix $s_{0}>\left\|a_{0}\right\|_{\infty}$. If $s \in\left(s_{0}, \infty\right)$, the solution operator $S$ of the problem

$$
L u+s u=f \text { on } \Omega \times \mathbb{R}, B(u)=0, u \tau \text { - periodic, } \quad f \in Y
$$

defined by $S f=u$, can be extended to an injective and bounded operator, that we still denote by $S$, from $X$ into $Y$ (see [9], Lemma 3.1). This provides an extension of the original differential operator $L$, which is a closed operator from a dense subspace $D \subset Y$ into $X$ (see [9], p. 12). From now on $L$ will denote this extension of the original differential operator.

If $a \in L_{\tau}^{\infty}(\Omega \times \mathbb{R})$ and $\delta_{1} \leq a+a_{0} \leq \delta_{2}$ for some positive constants $\delta_{1}$ and $\delta_{2}$, then $L+a I: X \rightarrow Y$ has a bounded inverse $(L+a I)^{-1}: X \rightarrow C_{B}^{1+\gamma, \gamma}(\bar{\Omega} \times \mathbb{R}) \subset Y$, i.e.

$$
\begin{equation*}
\left\|(L+a I)^{-1} f\right\|_{C^{1+\gamma, \gamma}(\bar{\Omega} \times \mathbb{R})} \leq c\|f\|_{L_{\tau}^{p}(\Omega \times \mathbb{R})} \tag{4}
\end{equation*}
$$

for some positive constant $c$ and all $f$ ([9], Lemma 3.1). So $(L+a I)^{-1}: X \rightarrow X$ and its restriction $(L+a I)_{\mid Y}^{-1}: Y \rightarrow Y$ are compact operators. Moreover, $(L+a I)_{\mid Y}^{-1}:$ $Y \rightarrow Y$ is a strongly positive operator ([9], Lemma 3.7).

If $\partial a_{i, j} / \partial x_{j} \in C(\bar{\Omega} \times \mathbb{R})$ for $1 \leq i, j \leq N$, we recall that for $f \in L_{\tau}^{p}(\Omega \times \mathbb{R})$,
$(L+a I)^{-1} f$ is a weak solution of the periodic problem

$$
\begin{aligned}
\frac{\partial u}{\partial t}-\operatorname{div}(A \nabla u)+\sum_{j}\left(b_{j}+\sum_{i} \frac{\partial a_{i, j}}{\partial x_{i}}\right) \frac{\partial u}{\partial x_{j}}+\left(a+a_{0}\right) u & =f \text { on } \Omega \times \mathbb{R}, \\
B(u) & =0 \text { on } \partial \Omega \times \mathbb{R} \\
u(x, t) & =u(x, t+\tau)
\end{aligned}
$$

where $A$ is the $N \times N$ matrix whose $i, j$ entry is $a_{i, j}$ (weak solutions defined as, e.g., in [10], taking there, the test functions space adapted to the periodicity and to the respective boundary condition) In fact, this is true for a Hölder continuous $f$ (classical solutions are weak solutions) and then an approximation process, using that $L$ is closed and (4), gives the assertion for a general $f$.

REMARK 2. Let $m \in L_{\tau}^{\infty}(\Omega \times \mathbb{R})$ and let $M: X \rightarrow X$ be the operator multiplication by $m$. Then for each $\lambda \in \mathbb{R}$ there exists a unique $\mu \in \mathbb{R}$ such that the problem $\left(P_{\mu}\right)$ from the introduction has a positive solution.

This is shown for $\lambda>0, a_{0} \geq 0$ in [9] (see Remark 3.9 and Lemma 3.10). A slight modification of the argument used there shows that this is true for $\lambda \in \mathbb{R}, a_{0} \in$ $L_{\tau}^{\infty}(\Omega \times \mathbb{R})$ (start with $L+r$ instead of $L$, with $r \in \mathbb{R}$ large enough). Thus $\mu_{m}(\lambda)$ is well defined for all $\lambda \in \mathbb{R}, \mu_{m}$ is a concave function, $\mu_{m}(\lambda)$ is real analytic in $\lambda$, $\mu_{m}(\lambda)$ is an $M$ simple eigenvalue for $L$ and

$$
\mu_{m}(\lambda)=0 \text { if and only if } \lambda \text { is a principal eigenvalue for the weight } m .
$$

Moreover, the positive solution $u_{\lambda}$ of $\left(P_{\mu_{m}(\lambda)}\right)$ can be chosen real analytic in $\lambda$ (as a map from $\mathbb{R}$ into $Y$ ). As in the case $m$ Hölder continuous we have $\mu_{m}(-\lambda)=$ $\mu_{-m}(\lambda), \lambda \in \mathbb{R}$. We recall also that for the Dirichlet problem with $a_{0} \geq 0$ (and also for the Neumann problem with $a_{0} \geq 0, a_{0} \neq 0$ ) we have $\mu_{m}(0)>0$ (see [8], also [5] and [6]).

Given $\lambda \in \mathbb{R}$, we will say that the maximum principle (in brief MP) holds for $\lambda$ if $\lambda$ is not an eigenvalue for the weight $m$ and if $h \in X$ with $h \geq 0, h \neq 0$ implies that the solution $u_{\lambda}$ of the problem $\left(P_{\lambda, h}\right)$ belongs to $P_{Y}^{\circ}$.

The function $\mu_{m}$ describes what happens, with respect to the MP, at a given $\lambda \in \mathbb{R}$ (for the case $m$ Hölder continuous see [8], Theorem 16.6):

$$
\mu_{m}(\lambda)>0 \text { if and only if } \lambda \text { is not an eigenvalue and } M P \text { holds for } \lambda
$$

Indeed, for $h \in X$ with $h \geq 0, h \neq 0$, for $r \in \mathbb{R}$ large enough such that $-\left\|a_{0}\right\|_{\infty}-$ $\|\lambda m\|_{\infty}+r>0$, problem $\left(P_{\lambda, h}\right)$ is equivalent to $\left(r^{-1} I-S_{\lambda}\right) u=H_{\lambda}$ with $S_{\lambda}=$ $(L+r-\lambda M)^{-1}$ and $H_{\lambda}=r^{-1} S_{\lambda} h$. Now $H_{\lambda}>0$. Also $\mu_{m}(\lambda)>0$ if and only if $\widetilde{\rho}(\lambda)<r^{-1}$, where $\widetilde{\rho}(\lambda)$ is the spectral radius of $S_{\lambda}$ so Krein - Rutman Theorem ensures, for such a $u$, that $\mu_{m}(\lambda)>0$ is equivalent to $u \in P_{Y}^{\circ}$. Moreover, $\mu_{m}(\lambda)>0$ implies also that $\lambda$ is not an eigenvalue for the weight $m$, since, if $\lambda$ would be an eigenvalue with an associated eigenfunction $\Phi$ and if $u_{\lambda}$ is a positive solution of $L u=$
$\lambda m u+\mu_{m}(\lambda) u, B(u)=0$, then, for a suitable constant $c, v=u_{\lambda}+c \Phi$ would be a solution negative somewhere for the problem $L v=\lambda m v+\mu_{m}(\lambda) u_{\lambda}, B(v)=0$.

Next theorem shows that $\mu_{m}$ also describes what happens, with respect to the AMP, near to a principal eigenvalue.

THEOREM 1. Let L be the periodic parabolic operator given by (1) with coefficients satisfying the conditions stated there, let $B(u)=0$ be either the Dirichlet condition or the Neumann condition, consider Ygiven by (3), let $m$ be a function in $L_{\tau}^{\infty}(\Omega \times \mathbb{R})$ and let $\lambda^{*}$ be a principal eigenvalue for the weight $m$. Finally, let $u, v \in L_{\tau}^{p}(\Omega \times \mathbb{R})$ for some $p>N+2$, with $0<u \leq v$. Then
(a) If $\lambda \rightarrow \mu_{m}(\lambda)$ vanishes identically, then for all $\lambda \in R$ and all $h \geq 0, h \neq 0$ in $L_{\tau}^{p}(\Omega \times \mathbb{R})$, problem $\left(P_{\lambda, h}\right)$ has no solution.
(b) If $\mu_{m}^{\prime}\left(\lambda^{*}\right)<0$ (respectively $\mu_{m}^{\prime}\left(\lambda^{*}\right)>0$ ), then the AMP holds to the right of $\lambda^{*}$ (respectivley to the left) and its holds uniformly on $h \in[u, v]$, i.e. there exists $\delta_{u, v}>0$ such that for each $\lambda \in\left(\lambda^{*}, \lambda^{*}+\delta_{u, v}\right)$ (respectively $\lambda \in\left(\lambda^{*}-\delta_{u, v}, \lambda^{*}\right)$ ) and for each $h \in[u, v]$, the solution $u_{\lambda, h}$ of $\left(P_{\lambda, h}\right)$ satisfies $u_{\lambda, h} \in-P_{Y}^{\circ}$.
(c) If $\mu_{m}^{\prime}\left(\lambda^{*}\right)=0$ and if $\mu_{m}$ does not vanishes identically, then the AMP holds uniformly on $h$ for $h \in[u, v]$ right and left of $\lambda^{*}$, i.e. there exists $\delta_{u, v}>0$ such that for $0<\left|\lambda-\lambda^{*}\right|<\delta_{u, v}, h \in[u, v]$, the solution $u_{\lambda, h}$ of $\left(P_{\lambda, h}\right)$ is in $-P_{Y}^{\circ}$.

Proof. Let $M: X \rightarrow X$ be the operator multiplication by $m$. Given a closed interval $I$ around $\lambda^{*}$ we choose $r \in(0, \infty)$ such that $r>\lambda^{*} \mu_{m}^{\prime}\left(\lambda^{*}\right)$ and $r-\|\lambda m\|_{\infty}-\left\|a_{0}\right\|_{\infty}>0$ for all $\lambda \in I$. For a such $r$ and for $\lambda \in I$, let $T_{\lambda}: Y \rightarrow Y$ defined by

$$
T_{\lambda}=(L+r I)^{-1}(\lambda M+r I)
$$

so each $T_{\lambda}$ is a strongly positive and compact operator on $Y$ with a positive spectral radius $\rho(\lambda)$ that is an algebraically simple eigenvalue for $T_{\lambda}$ and $T_{\lambda}^{*}$. Let $\Phi_{\lambda}, \Psi_{\lambda}$ be the corresponding positive eigenvectors normalized by $\left\|\Phi_{\lambda}\right\|=1$ and $\left\langle\Psi_{\lambda}, \Phi_{\lambda}\right\rangle=1$. By Lemma $1, \rho(\lambda)$ is real analytic in $\lambda$ and $\Phi_{\lambda}, \Psi_{\lambda}$ are continuous in $\lambda$. As a consequence of Krein - Rutman, we have that $\rho(\lambda)=1$ iff $\lambda$ is a principal eigenvalue for the weight $m$. So $\rho\left(\lambda^{*}\right)=1$. Since $T_{\lambda}$ is strongly positive we have $\Phi_{\lambda} \in P_{Y}^{\circ}$, so there exists $s>0$ such that $B_{s}^{Y}(0) \subset\left(-\Phi_{\lambda}, \Phi_{\lambda}\right)$ for all $\lambda \in I$. Let $H=(L+r)^{-1} h, U=(L+r)^{-1} u$ and $V=(L+r)^{-1} v$. The problem $L u_{\lambda}=\lambda m u_{\lambda}+h$ on $\Omega \times \mathbb{R}, B\left(u_{\lambda}\right)=0$ on $\partial \Omega \times \mathbb{R}$ is equivalent to

$$
\begin{equation*}
u_{\lambda}=\left(I-T_{\lambda}\right)^{-1} H \tag{5}
\end{equation*}
$$

and $u \leq h \leq v$ implies $U \leq H \leq V$. So, we are in the hypothesis of our Lemma 2 and from its proof we get that $\left\|\left(\left(\rho(\lambda) I-T_{\lambda \mid \Psi_{\lambda}^{\perp}}\right)\right)^{-1}\right\|$ remains bounded for $\lambda$ near to $\lambda^{*}$ and from (2) with $\theta=\rho\left(\lambda^{*}\right)=1$ we obtain

$$
\begin{equation*}
u_{\lambda}=\frac{\left\langle\Psi_{\lambda}, H\right\rangle}{1-\rho(\lambda)}\left[\Phi_{\lambda}+\frac{1-\rho(\lambda)}{\left\langle\Psi_{\lambda}, H\right\rangle}\left(\left(I-T_{\lambda}\right)_{\mid \Psi_{\lambda}^{\perp}}\right)^{-1} w_{H, \lambda}\right] \tag{6}
\end{equation*}
$$

where $w_{H, \lambda}=H-\left\langle\Psi_{\lambda}, H\right\rangle \Phi_{\lambda}$.
$T_{\lambda} \Phi_{\lambda}=\rho(\lambda) \Phi_{\lambda}$ is equivalent to $L \Phi_{\lambda}=\frac{\lambda}{\rho(\lambda)} m \Phi_{\lambda}+r\left(\frac{1}{\rho(\lambda)}-1\right) \Phi_{\lambda}$ and, since $\Phi_{\lambda}>0$ this implies

$$
\begin{equation*}
\mu_{m}\left(\frac{\lambda}{\rho(\lambda)}\right)=r\left(\frac{1}{\rho(\lambda)}-1\right) \tag{7}
\end{equation*}
$$

If $\mu_{m}$ vanishes identically then $\rho(\lambda)=1$ for all $\lambda$ and (5) has no solution for all $h \geq 0$, $h \neq 0$. This gives assertion (a) of the theorem.

For (b) suppose that $\mu_{m}^{\prime}\left(\lambda^{*}\right) \neq 0$. Taking the derivative in (7) at $\lambda=\lambda^{*}$ and recalling that $\rho\left(\lambda^{*}\right)=1$ we obtain

$$
\mu_{m}^{\prime}\left(\lambda^{*}\right)\left(1-\lambda^{*} \rho^{\prime}\left(\lambda^{*}\right)\right)=-r \rho^{\prime}\left(\lambda^{*}\right)
$$

so $\rho^{\prime}\left(\lambda^{*}\right)=\mu_{m}^{\prime}\left(\lambda^{*}\right) /\left(\lambda^{*} \mu_{m}^{\prime}\left(\lambda^{*}\right)-r\right)$. We have chosen $r>\lambda^{*} \mu_{m}^{\prime}\left(\lambda^{*}\right)$, thus $\mu_{m}^{\prime}\left(\lambda^{*}\right)>0$ implies $\rho^{\prime}\left(\lambda^{*}\right)<0$ and $\mu_{m}^{\prime}\left(\lambda^{*}\right)<0$ implies $\rho^{\prime}\left(\lambda^{*}\right)>0$ and then, proceeding as at the end of the proof of Lemma 2, assertion (b) of the theorem follows from (6).

If $\mu_{m}^{\prime}\left(\lambda^{*}\right)=0$, since $\mu_{m}$ is concave and analytic we have $\mu_{m}(\lambda)<0$ for $\lambda \neq \lambda^{*}$. Then $1 / \rho$ has a local maximum at $\lambda^{*}$ and (c) follows from (6) as above.

To formulate conditions on $m$ to fulfill the assumptions of Theorem 1 we recall the quantities

$$
P(m)=\int_{0}^{\tau} e s s \sup _{x \in \Omega} m(x, t) d t, \quad N(m)=\int_{0}^{\tau} e s s i n f_{x \in \Omega} m(x, t) d t
$$

The following two theorems describe completely the possibilities, with respect to the AMP, in Neumann and Dirichlet cases with $a_{0} \geq 0$.

THEOREM 2. Let L be given by (1). Assume that either $B(u)=0$ is the Neumann condition and $a_{0} \geq 0, a_{0} \neq 0$ or that $B(u)=0$ is the Dirichlet condition and $a_{0} \geq 0$. Assume in addition that $\partial a_{i, j} / \partial x_{j} \in C(\bar{\Omega} \times \mathbb{R}), 1 \leq i, j \leq N$. Then
(1) If $P(m)>0(P(m) \leq 0), N(m) \geq 0(N(m)<0)$ then there exists a unique principal eigenvalue $\lambda^{*}$ that is positive (negative) and the AMP holds to the right (to the left) of $\lambda^{*}$
(2) If $P(m)>0, N(m)<0$ then there exist two principal eigenvalues $\lambda_{-1}<0$ and $\lambda_{1}>0$ and the AMP holds to the right of $\lambda_{1}$ and to the left of $\lambda_{-1}$.
(3) If $P(m)=$ then $N(m)=0$ then there are no principal eigenvalues.

Moreover if $u, v \in L_{\tau}^{p}(\Omega \times \mathbb{R})$ satisfy $0<u<v$, then in (1) and (2) the AMP holds uniformly on $h$ for $h \in[u, v]$.

Proof. We consider first the Dirichlet problem. If $P(m)>0$ and $N(m) \geq 0$ then there exists a unique principal eigenvalue $\lambda_{1}$ that is positive ([9]). Since $\mu_{m}(0)>0$, $\mu_{m}\left(\lambda_{1}\right)=0$ and $\mu$ is concave, we have $\mu_{m}^{\prime}\left(\lambda_{1}\right)<0$ and (b) of Theorem 1 applies. If $P(m)>0$ and $N(m)<0$ then there exist two eigenvalues $\lambda_{-1}<0$ and $\lambda_{1}>0$ because $\mu_{-m}(-\lambda)=\mu_{m}(\lambda)$ and in this case ( $\mu_{m}$ is concave) we have $\mu_{m}^{\prime}\left(\lambda_{-1}\right)>0$ and $\mu_{m}^{\prime}\left(\lambda_{1}\right)<0$, so Theorem 1 applies. In each case Theorem 1 gives the required uniformity. The other cases are similar. If $P(m)=N(m)=0$ then $m=m(t)$ and $\mu_{m}(\lambda) \equiv \mu_{m}(0)>0([9])$. So (3) holds. Results in [9] for the Dirichlet problem remain valid for the Neumann condition with $a_{0} \geq 0, a_{0} \neq 0$, so the above proof holds.

REMARK 3. If $a_{0}=0$ in the Neumann problem then $\lambda_{0}=0$ is a principal eigenvalue and $\mu_{m}(0)=0$. To study this case we recall that $(L+1)^{-1 *}$ has a positive eigenvector $\Psi \in X^{\prime} \subset Y^{\prime}$ provided by the Krein - Rutman Theorem and (4). Then $\mu_{m}^{\prime}(0)=-\langle\Psi, m\rangle /\langle\Psi, 1\rangle$ where $\langle\Psi, m\rangle=\int_{\Omega \times(0, \tau)} \psi m$ makes sense because $\Psi \in L^{p^{\prime}}(\Omega \times(0, \tau))$ ([9], remark 3.8). Indeed, let $u_{\lambda}$ be a positive $\tau$ periodic solution of $L u_{\lambda}=\lambda m u_{\lambda}+\mu_{m}(\lambda) u_{\lambda}$ on $\Omega \times \mathbb{R}, B\left(u_{\lambda}\right)=0$ with $u_{\lambda}$ real analytic in $\lambda$ and with $u_{0}=1$. Since $\Psi$ vanishes on the range of $L$ we have $0=\lambda\left\langle\Psi, m u_{\lambda}\right\rangle+\mu_{m}(\lambda)\left\langle\Psi, u_{\lambda}\right\rangle$. Taking the derivative at $\lambda=0$ and using that $u_{0}=1$ we get the above expression for $\mu_{m}^{\prime}(0)$.

THEOREM 3. Let L be given by (1). Assume that $B(u)=0$ is the Neumann condition and that $a_{0}=0$. Assume in addition that $\partial a_{i, j} / \partial x_{j} \in C(\bar{\Omega} \times \mathbb{R}), 1 \leq$ $i, j \leq N$. Let $\Psi$ be as in Remark 3.

Then, if $m$ is not a function of $t$ alone, we have
(1) If $\langle\Psi, m\rangle<0(\langle\Psi, m\rangle>0), P(m) \leq 0(N(m) \geq 0)$, then 0 is the unique principal eigenvalue and the $A M P$ holds to the left (to the right) of 0 .
(2) If $\langle\Psi, m\rangle<0(\langle\Psi, m\rangle>0), P(m)>0(N(m)<0)$, then there exists two principal eigenvalues, 0 and $\lambda^{*}$ wich is positive (negative) and the AMP holds to the left (to the right) of 0 and to the right (to the left) of $\lambda^{*}$.
(3) If $\langle\Psi, m\rangle=0$,hen 0 is the unique principal eigenvalue and the AMP holds left and right of 0 .

If $m=m(t)$ is a function of $t$ alone, then we have
(1') If $\int_{0}^{\tau} m(t) d t=0$ then for all $\lambda \in \mathbb{R}$ the above problem $L u=\lambda m u+h$ has no solution.
(2') If $\int_{0}^{\tau} m(t) d t \neq 0$ and $\langle\Psi, m\rangle>0(\langle\Psi, m\rangle<0)$ then 0 is the unique principal eigenvalue and the AMP holds to the right (to the left) of 0.

Moreover, if $u, v \in X$ satisfy $0<u<v$, then in each case (except (1')) the AMP holds uniformly on $h$ for $h \in[u, v]$.

Proof. Suppose that $m$ is not a function of $t$ alone. If $\langle\Psi, m\rangle<0, P(m) \leq 0$ then 0 is the unique principal eigenvalue and $\mu_{m}^{\prime}(0)>0$. If $\langle\Psi, m\rangle<0, P(m)>0$ then there exist two principal eigenvalues: 0 and some $\lambda_{1}>0$ and since $\mu_{m}$ is concave we have $\mu_{m}^{\prime}(0)>0$ and $\mu_{m}^{\prime}\left(\lambda_{1}\right)<0$. If $\langle\Psi, m\rangle=0$ and if $m$ is not funtion of $t$ alone, then $\mu_{m}$ is not a constant and $\mu_{m}^{\prime}(0)=0$ and 0 is the unique principal eigenvalue. In each case, the theorem follows from Theorem 1. The other cases are similar.

If $m$ is a function of $t$ alone then $\mu_{m}(\lambda)=-\frac{P(m)}{\tau} \lambda$, this implies $\frac{P(m)}{\tau}=$ $\langle\Psi, m\rangle /\langle\Psi, 1\rangle$. If $\int_{0}^{\tau} m(t) d t=0$ then $\mu_{m}=0$ and (a) of Theorem 1 applies. If $\int_{0}^{\tau} m(t) d t \neq 0$ and $\langle\Psi, m\rangle>0$ then $P(m)=N(m)>0$ so 0 is the unique principal eigenvalue and $\mu_{m}^{\prime}(0)<0$, in this case Theorem 1 applies also. The case $\langle\Psi, m\rangle<0$ is similar. The remaining case $\int_{0}^{\tau} m(t) d t \neq 0$ and $\langle\Psi, m\rangle=0$ is impossible because $\frac{P(m)}{\tau}=\langle\Psi, m\rangle /\langle\Psi, 1\rangle$.

For one dimensional Neumann problems, similarly to the elliptic case, a uniform AMP holds.

Theorem 4. Suppose $N=1, \Omega=(\alpha, \beta)$ and the Neumann condition. Let $L$ be given by $L u=u_{t}-a u_{x x}+b u_{x}+a_{0} u$, where $a_{0}, b \in C_{\tau}^{\gamma, \gamma / 2}(\bar{\Omega} \times \mathbb{R}), a_{0} \geq 0$ and with $a \in C_{\tau}^{1}(\bar{\Omega} \times \mathbb{R}), \min _{x \in \bar{\Omega} \times \mathbb{R}} a(x, t)>0$. Then the AMP holds uniformly in $h$ (i.e. holds on an interval independent of $h$ ) in each situation considered in Theorem 3.

Proof. Let $\lambda^{*}$ be a principal eigenvalue for $L u=\lambda m u$. Without loss of generality we can assume that $\|h\|_{p}=1$ and that the AMP holds to the right of $\lambda^{*}$. Denote $M$ the operator multiplication by $m$. Let $I_{\lambda^{*}}$ be a finite closed interval around $\lambda^{*}$ and, for $\lambda \in I_{\lambda^{*}}$, let $T_{\lambda}, \Phi_{\lambda}$ and $\Psi_{\lambda}$ be as in the proof of Theorem 1. Each $\Phi_{\lambda}$ belongs to the interior of the positive cone in $C(\bar{\Omega})$ and $\lambda \rightarrow \Phi_{\lambda}$ is a continuous map from $I_{\lambda^{*}}$ into $C(\bar{\Omega})$, thus there exist positive constants $c_{1}, c_{2}$ such that

$$
\begin{equation*}
c_{1} \leq \Phi_{\lambda}(x) \leq c_{2} \tag{8}
\end{equation*}
$$

for all $\lambda \in I_{\lambda^{*}}$ and $x \in \bar{\Omega}$. As in the proof of Lemma 2 we obtain (6). Taking into account (8) and that $\left\|\left(\left(I-T_{\lambda}\right)_{\mid \Psi_{\lambda}^{\perp}}\right)^{-1}\right\|$ remains bounded for $\lambda$ near $\lambda^{*}$, in order to prove our theorem, it is enough to see that there exist a positive constant $c$ independent of $h$, such that $\left\|\left(I-T_{\lambda \mid \Psi_{\lambda}^{\perp}}\right)^{-1} w_{H, \lambda}\right\|_{\infty} /\left\langle\Psi_{\lambda}, H\right\rangle<c$ for $\lambda \in I_{\lambda^{*}}$. Since

$$
\frac{\left(I-T_{\lambda \mid \Psi_{\lambda}^{\perp}}\right)^{-1} w_{H, \lambda}}{\left\langle\Psi_{\lambda}, H\right\rangle}=\left(I-T_{\lambda \mid \Psi_{\lambda}^{\perp}}\right)^{-1}\left(\frac{H}{\left\langle\Psi_{\lambda}, H\right\rangle}-\Phi_{\lambda}\right)
$$

it suffices to prove that there exists a positive constant $c$ such that

$$
\begin{equation*}
\|H\|_{\infty} \leq c\left\langle\Psi_{\lambda}, H\right\rangle \tag{9}
\end{equation*}
$$

for all $h \geq 0$ with $\|h\|_{p}=1$. To show (9) we proceed by contradiction. If (9) does not holds, we would have for all $j \in \mathbb{N}, h_{j} \in L_{\tau}^{p}(\Omega \times \mathbb{R})$ with $h_{j} \geq 0,\left\|h_{j}\right\|_{p}=1$ and $\lambda_{j} \in I_{\lambda^{*}}$ such that

$$
\begin{equation*}
\left\langle\Psi_{\lambda_{j}},(L+r I)^{-1} h_{j}\right\rangle<\frac{1}{j}\left\|(L+r I)^{-1} h_{j}\right\|_{\infty} \tag{10}
\end{equation*}
$$

Thus $\lim _{j \rightarrow \infty}\left\langle\Psi_{\lambda_{j}},(L+r I)^{-1} h_{j}\right\rangle=0$. We claim that

$$
\begin{equation*}
\left\|(L+r I)^{-1} h\right\|_{\infty} / \min _{\bar{\Omega} \times[0, \tau]}(L+r I)^{-1} h \leq c \tag{11}
\end{equation*}
$$

for some positive constant $c$ and all nonnegative and non zero $h \in L_{\tau}^{p}(\Omega \times \mathbb{R})$. If (11) holds, then

$$
\begin{gathered}
\left\langle\Psi_{\lambda_{j}},(L+r I)^{-1} h_{j}\right\rangle \geq\left\langle\Psi_{\lambda_{j}}, 1\right\rangle_{\bar{\Omega} \times[0, \tau]}(L+r I)^{-1} h_{j} \geq \\
\geq c^{\prime} c\left\|(L+r I)^{-1} h_{j}\right\|_{\infty}\left\langle\Psi_{\lambda_{j}}, \Phi_{\lambda_{j}}\right\rangle=c c^{\prime}\left\|(L+r I)^{-1} h_{j}\right\|_{\infty}
\end{gathered}
$$

for some positive constant $c^{\prime}$ independent of $j$, contradicting (10).
It remains to prove (11) which looks like an elliptic Harnack inequality. We may suppose $\alpha=0$. Extending $u:=(L+r I)^{-1} h$ by parity to $[-\beta, \beta]$ we obtain a function $\widetilde{u}$ with $\widetilde{u}(-\beta, t)=\widetilde{u}(\beta, t)$, so we can assume that $\widetilde{u}$ is $2 \beta$ - periodic in $x$ and $\tau$ periodic in $t . \widetilde{u}$ solves weakly the equation $\widetilde{u}_{t}-\widetilde{a} \widetilde{u}_{x x}+\widetilde{b} \widetilde{u}_{x}+\left(a_{0}+r\right) \widetilde{u}=\widetilde{h}$ in $\mathbb{R} \times \mathbb{R}$ where $\widetilde{a}, \widetilde{a}_{0}, \widetilde{h}$ are extensions to $\mathbb{R} \times \mathbb{R}$ like $\widetilde{u}$, but $b$ is extended to an odd function $\widetilde{b}$ in $(-\beta, \beta)$ then $2 \beta$ periodically. In spite of discontinuities of $\widetilde{b}$, a parabolic Harnack inequality holds for $\widetilde{u}=\widetilde{u}(\widetilde{h})([10]$, Th. 1.1) and using the periodicity of $\widetilde{u}$ in $t$, we obtain (11).

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AMS Subject Classification: 35K20, 35P05, 47N20.
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Lavoro pervenuto in redazione il 03.05.2001 e, in forma definitiva, il 25.02.2002.


[^0]:    *Research partially supported by Agencia Córdoba Ciencia, Conicet and Secyt-UNC.

