# SECOND ORDER VARIATIONS IN VARIATIONAL SEQUENCES 

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#### Abstract

In this note we provide a geometrical characterization of the second order variation of a generalized Lagrangian in the framework of variational sequences.

We define the variational vertical derivative as an operator on the sheaves of the variational sequence and stress its link with the classical concept of variation. The main result is the intrinsic formulation of a theorem which states the relation between the variational vertical derivative of the Euler-Lagrange morphism of a generalized Lagrangian and the Euler-Lagrange morphism of the variational vertical derivative of the Lagrangian itself.


## 1. Introduction

Our framework is the calculus of variations on finite order jets of a fibered manifold. More precisely, we consider the geometrical formulation of this framework in terms of variational sequences introduced by Krupka [12, 13]. As it is well known, in this formulation the variational sequence is defined as a quotient of the de Rham sequence on a finite order jet of a fibered manifold with respect to an intrinsically defined subsequence, the "contact" subsequence. Standard objects of the calculus of variations can be interpreted as sheaf sections and morphisms in the variational sequence, which turns out to be an exact resolution of the constant sheaf $\mathbb{R}$ over the relevant fibered manifold.

We provide a geometrical characterization of the second order variation (see e.g. $[1,2,3,4,6,18]$ ) of a Lagrangian in the framework of finite order variational sequences. We introduce the notion of iterated variation of a section as an $i$-parameter 'deformation' of the section by means of vertical flows and thus define the $i-$ th variation of a morphism which is very simply related with the iterated Lie derivative of the morphism itself. Relying on previous results of us [5] on the representation of the Lie derivative operator in the variational sequence we can then define an operator on the quotient sheaves of the sequence, the variational vertical derivative. We stress some linearity properties of this operator and show

[^0]that it is a functor on the category of variational sequences. Making use of suitable representations of the variational vertical derivative we relate the second order variation of a generalized Lagrangian with the variational Lie derivative of generalized Euler-Lagrange operators associated with the Lagrangian itself.

The Lagrangian characterization of the second variation of a Lagrangian in the framework of jet bundles has been considered in [2, 3, 4, 6]. In particular, in [2, 3, 4] it was shown how to recast (up to divergencies) the system formed by the EulerLagrange equations together with the Jacobi equations for a given Lagrangian as the Euler Lagrange equations for a 'deformed' Lagrangian.

As an outcome of the intrinsic representation of the second order variation of a generalized Lagrangian, the above mentioned invariant decomposition of the second variation is here provided at any jet order and geometrically interpreted as a simple and direct application of the representation of the variational Lie derivative of variational morphisms provided in [5] together with a suitable version of a global decomposition formula of vertical morphisms due to Kolář [7, 9]. A new geometric object, the generalized Jacobi morphism, is finally represented in the variational sequence.

## 2. Variational sequences on Jets of fibered manifolds

Our framework is a fibered manifold $\pi: \boldsymbol{Y} \rightarrow \boldsymbol{X}$, with $\operatorname{dim} \boldsymbol{X}=n$ and $\operatorname{dim} \boldsymbol{Y}=$ $n+m$ (see e.g. [17]).

For $r \geq 0$ we are concerned with the $r$-jet space $J_{r} \boldsymbol{Y}$; in particular, we set $J_{0} \boldsymbol{Y} \equiv$ $\boldsymbol{Y}$. We recall the natural fiberings $\pi_{s}^{r}: J_{r} \boldsymbol{Y} \rightarrow J_{s} \boldsymbol{Y}, r \geq s, \pi^{r}: J_{r} \boldsymbol{Y} \rightarrow \boldsymbol{X}$, and, among these, the affine fiberings $\pi_{r-1}^{r}$. We denote by $V \boldsymbol{Y}$ the vector subbundle of the tangent bundle $T \boldsymbol{Y}$ of vectors on $\boldsymbol{Y}$ which are vertical with respect to the fibering $\pi$.

Charts on $\boldsymbol{Y}$ adapted to $\pi$ are denoted by $\left(x^{\lambda}, y^{i}\right)$. Greek indices $\lambda, \mu, \ldots$ run from 1 to $n$ and they label base coordinates, while Latin indices $i, j, \ldots$ run from 1 to $m$ and label fibre coordinates, unless otherwise specified. We denote by $\left(\partial_{\lambda}, \partial_{i}\right)$ and $\left(d^{\lambda}, d^{i}\right)$ the local bases of vector fields and 1 -forms on $\boldsymbol{Y}$ induced by an adapted chart, respectively.

We denote multi-indices of dimension $n$ by boldface Greek letters such as $\boldsymbol{\alpha}=$ $\left(\alpha_{1}, \ldots, \alpha_{n}\right)$, with $0 \leq \alpha_{\mu}, \mu=1, \ldots, n$; by an abuse of notation, we denote with $\lambda$ the multi-index such that $\alpha_{\mu}=0$, if $\mu \neq \lambda, \alpha_{\mu}=1$, if $\mu=\lambda$. We also set $|\boldsymbol{\alpha}|:=\alpha_{1}+\cdots+\alpha_{n}$ and $\boldsymbol{\alpha}!:=\alpha_{1}!\ldots \alpha_{n}!$. The charts induced on $J_{r} \boldsymbol{Y}$ are denoted by $\left(x^{\lambda}, y_{\boldsymbol{\alpha}}^{i}\right)$, with $0 \leq|\boldsymbol{\alpha}| \leq r$; in particular, we set $y_{\mathbf{0}}^{i} \equiv y^{i}$. The local vector fields and forms of $J_{r} \boldsymbol{Y}$ induced by the above coordinates are denoted by $\left(\partial_{i}^{\boldsymbol{\alpha}}\right)$ and $\left(d_{\boldsymbol{\alpha}}^{i}\right)$, respectively.

For $r \geq 1$, we consider the natural complementary fibered morphisms over the affine fibering $J_{r} \boldsymbol{Y} \rightarrow J_{r-1} \boldsymbol{Y}$ induced by contact maps on jet spaces

$$
\text { Д }: J_{r} \boldsymbol{Y} \underset{\boldsymbol{X}}{\times} T \boldsymbol{X} \rightarrow T J_{r-1} \boldsymbol{Y}, \quad \vartheta: J_{r} \boldsymbol{Y} \underset{J_{r-1} \boldsymbol{Y}}{\times} T J_{r-1} \boldsymbol{Y} \rightarrow V J_{r-1} \boldsymbol{Y},
$$

with coordinate expressions, for $0 \leq|\boldsymbol{\alpha}| \leq r-1$, given by

$$
\text { Д }=d^{\lambda} \otimes \text { Д }_{\lambda}=d^{\lambda} \otimes\left(\partial_{\lambda}+y_{\boldsymbol{\alpha}+\lambda}^{j} \partial_{j}^{\boldsymbol{\alpha}}\right), \quad \vartheta=\vartheta_{\boldsymbol{\alpha}}^{j} \otimes \partial_{j}^{\boldsymbol{\alpha}}=\left(d_{\boldsymbol{\alpha}}^{j}-y_{\boldsymbol{\alpha}+\lambda}^{j} d^{\lambda}\right) \otimes \partial_{j}^{\boldsymbol{\alpha}} .
$$

We have the following natural fibered splitting

$$
\begin{equation*}
J_{r} \boldsymbol{Y} \underset{J_{r-1} \boldsymbol{Y}}{\times} T^{*} J_{r-1} \boldsymbol{Y}=\left(J_{r} \boldsymbol{Y} \underset{J_{r-1} \boldsymbol{Y}}{\times} T^{*} \boldsymbol{X}\right) \oplus \stackrel{\mathcal{C}}{ }_{r-1}[\boldsymbol{Y}] \tag{1}
\end{equation*}
$$

where $\stackrel{*}{\mathcal{C}}_{r-1}[\boldsymbol{Y}]:=\operatorname{im} \vartheta_{r}^{*}$ and the canonical isomorphism $\stackrel{*}{\mathcal{C}}_{r-1}[\boldsymbol{Y}] \simeq J_{r} \boldsymbol{Y}{\underset{J}{r-1}}^{\times} \boldsymbol{Y}$ $V^{*} J_{r-1} \boldsymbol{Y}$ holds true (see $[12,13,15,17]$ ).

The above splitting induces also a decomposition of the exterior differential on $\boldsymbol{Y},\left(\pi_{r}^{r+1}\right)^{*} \circ d=d_{H}+d_{V}$, where $d_{H}$ and $d_{V}$ are called the horizontal and vertical differential, respectively. The action of $d_{H}$ and $d_{V}$ on functions and 1-forms on $J_{r} \boldsymbol{Y}$ uniquely characterizes $d_{H}$ and $d_{V}$ (see, e.g., [17, 21] for more details).

If $f: J_{r} \boldsymbol{Y} \rightarrow \mathbb{R}$ is a function, then we set $D_{\lambda} f:=Д_{\lambda} f, D_{\alpha+\lambda} f:=D_{\lambda} D_{\boldsymbol{\alpha}} f$, where the operator $D_{\lambda}$ is the standard formal derivative.

A projectable vector field on $\boldsymbol{Y}$ is defined to be a pair $(\Xi, \xi)$, where $\Xi: \boldsymbol{Y} \rightarrow$ $T \boldsymbol{Y}$ and $\xi: \boldsymbol{X} \rightarrow T \boldsymbol{X}$ are vector fields and $\Xi$ is a fibered morphism over $\xi$. A projectable vector field $(\Xi, \xi)$ can be conveniently prolonged to a projectable vector field $\left(j_{r} \Xi, \xi\right)$, whose coordinate expression can be found e.g. in [5] and [11, 15, 17]. A vertical vector field on $\boldsymbol{Y}$ is a projectable vector field on $\boldsymbol{Y}$ such that $\xi=0$.
i. For $r \geq 0$, we consider the standard sheaves $\stackrel{p}{\Lambda}_{r}$ of $p$-forms on $J_{r} \boldsymbol{Y}$.
ii. For $0 \leq s \leq r$, we consider the sheaves $\stackrel{p}{\mathcal{H}}_{(r, s)}$ and $\stackrel{p}{\mathcal{H}}_{r}$ of horizontal forms, i.e. of local fibered morphisms over $\pi_{s}^{r}$ and $\pi^{r}$ of the type $\alpha: J_{r} \boldsymbol{Y} \rightarrow \stackrel{p}{\wedge} T^{*} J_{s} \boldsymbol{Y}$ and $\beta: J_{r} \boldsymbol{Y} \rightarrow \stackrel{p}{\Lambda} T^{*} \boldsymbol{X}$, respectively.
iii. For $0 \leq s<r$, we consider the subsheaf $\stackrel{p}{\mathcal{C}}_{(r, s)} \subset \stackrel{p}{\mathcal{H}}_{(r, s)}$ of contact forms, i.e. of sections $\alpha \in \stackrel{p}{\mathcal{H}}_{(r, s)}$ with values into ${ }_{\wedge}^{p}\left(\mathcal{C}_{s}[\boldsymbol{Y}]\right)$. There is a distinguished subsheaf $\stackrel{p}{\mathcal{C}}_{r} \subset \stackrel{p}{\mathcal{C}}_{(r+1, r)}$ of local fibered morphisms $\alpha \in \stackrel{p}{\mathcal{C}}_{(r+1, r)}$ such that $\alpha=\stackrel{p}{\wedge} \operatorname{im} \vartheta_{r+1}^{*}[\boldsymbol{Y}] \mathrm{o} \tilde{\alpha}$, where $\tilde{\alpha}$ is a section of the fibration $J_{r+1} \boldsymbol{Y} \underset{J_{r} \boldsymbol{Y}}{\times} \wedge{ }^{p} V^{*} J_{r} \boldsymbol{Y} \rightarrow J_{r+1} \boldsymbol{Y}$ which projects down onto $J_{r} \boldsymbol{Y}$.

According to [21], the fibered splitting (1) yields naturally the sheaf splitting $\stackrel{p}{\mathcal{H}}_{(r+1, r)}=\bigoplus_{t=0}^{p} \stackrel{p-t}{\mathcal{C}}_{(r+1, r)} \wedge \stackrel{t}{\mathcal{H}}_{r+1}$, which restricts to the inclusion $\stackrel{p}{\Lambda}_{r} \subset \bigoplus_{t=0}^{p}$ ${ }^{p-t}{ }_{r} \wedge \stackrel{t}{\mathcal{H}}_{r+1}^{h}$, where $\stackrel{p}{\mathcal{H}}_{r+1}^{h}:=h\left(\stackrel{p}{\Lambda}_{r}\right)$ for $0<p \leq n$ and $h$ is defined to be the restriction to $\stackrel{p}{\Lambda}_{r}$ of the projection of the above splitting onto the non-trivial summand with the highest value of $t$. We define also the map $v:=\mathrm{id}-h$.

Let $\alpha \in \stackrel{1}{\mathcal{C}}_{r} \wedge \stackrel{n}{\mathcal{H}}_{r+1}^{h}$. Then there is a unique pair of sheaf morphisms

$$
\begin{equation*}
E_{\alpha} \in \stackrel{1}{\mathcal{C}}_{(2 r, 0)} \wedge \stackrel{n}{\mathcal{H}}_{2 r+1}^{h}, \quad F_{\alpha} \in \stackrel{1}{\mathcal{C}}_{(2 r, r)} \wedge \stackrel{n}{\mathcal{H}}_{2 r+1}^{h}, \tag{2}
\end{equation*}
$$

such that $\left(\pi_{r+1}^{2 r+1}\right)^{*} \alpha=E_{\alpha}-F_{\alpha}$, and $F_{\alpha}$ is locally of the form $F_{\alpha}=d_{H} p_{\alpha}$, with $p_{\alpha} \in \stackrel{1}{\mathcal{C}}_{(2 r-1, r-1)} \wedge \stackrel{n-1}{\mathcal{H}}_{2 r}($ see e.g. [9, 11, 21] $)$.

Recall (see [21]) that if $\beta \in \stackrel{1}{\mathcal{C}}_{s} \wedge \stackrel{1}{\mathcal{C}}_{(s, 0)} \wedge \stackrel{n}{\mathcal{H}}_{s}$, then, there is a unique $\tilde{H}_{\beta} \in$ $\stackrel{1}{\mathcal{C}}_{(2 s, s)} \otimes \stackrel{1}{\mathcal{C}}_{(2 s, 0)} \wedge \stackrel{n}{\mathcal{H}}_{2 s}$ such that, for all $\Xi: \boldsymbol{Y} \rightarrow V \boldsymbol{Y}, E_{\widehat{\beta}}=C_{1}^{1}\left(j_{2 s} \Xi \otimes \tilde{H}_{\beta}\right)$, where $\left.\hat{\beta}:=j_{s} \Xi\right\lrcorner \beta, C_{1}^{1}$ stands for tensor contraction and $\lrcorner$ denotes inner product. Then there is a unique pair of sheaf morphisms

$$
\begin{equation*}
H_{\beta} \in \stackrel{1}{\mathcal{C}}_{(2 s, s)} \wedge \stackrel{1}{\mathcal{C}}_{(2 s, 0)} \wedge \stackrel{n}{\mathcal{H}}_{2 s}, \quad G_{\beta} \in \stackrel{2}{\mathcal{C}}_{(2 s, s)} \wedge \stackrel{n}{\mathcal{H}}_{2 s} \tag{3}
\end{equation*}
$$

such that $\pi_{s}^{2 s^{*}} \beta=H_{\beta}-G_{\beta}$ and $H_{\beta}=\frac{1}{2} A\left(\tilde{H}_{\beta}\right)$, where $A$ stands for antisymmetrisation. Moreover, $G_{\beta}$ is locally of the type $G_{\beta}=d_{H} q_{\beta}$, where $q_{\beta} \in \stackrel{2}{\mathcal{C}}_{2 s-1} \wedge \stackrel{n-1}{\mathcal{H}}_{2 s-1}$, hence $[\beta]=\left[H_{\beta}\right]$. Coordinate expressions of the morphisms $E_{\alpha}$ and $H_{\beta}$ can be found in [21].
2.1. Variational sequences. We recall now the theory of variational sequences on finite order jet spaces, as it was developed by Krupka in [12]. By an abuse of notation, denote by $d \operatorname{ker} h$ the sheaf generated by the presheaf $d \operatorname{ker} h$. Set $\stackrel{*}{\Theta}_{r}:=$ ker $h+d$ ker $h$. The following diagram is commutative and its rows and columns are exact:


Definition 2.1. The bottom row of the above diagram is called the $r$-th order variational sequence associated with the fibered manifold $\boldsymbol{Y} \rightarrow \boldsymbol{X}$.

The quotient sheaves in the variational sequence can be conveniently represented as follows [21].

Let $k \leq n$. Then, the sheaf morphism $h$ yields the natural isomorphism

$$
I_{k}: \stackrel{k}{\Lambda}_{r} / \stackrel{k}{\Theta}_{r} \rightarrow \stackrel{k}{\mathcal{H}}_{r+1}^{h}:=\stackrel{k}{\mathcal{V}}_{r}:[\alpha] \mapsto h(\alpha) .
$$

Let $k>n$. Then, the projection $h$ induces a natural sheaf isomorphism

$$
I_{k}:\left(\begin{array}{l}
k \\
\Lambda_{r} / \Theta^{k} \\
r
\end{array}\right) \rightarrow\left(\stackrel{k-n}{\mathcal{C}}_{r} \wedge \stackrel{n}{\mathcal{H}}_{r+1}^{h}\right) / h(\overline{d \operatorname{ker} h}):=\stackrel{\mathcal{V}}{r}_{r}:[\alpha] \mapsto[h(\alpha)] .
$$

We remark that a section $\lambda \in \mathcal{V}_{r}$ is just a Lagrangian of order $(r+1)$ of the standard literature. Furthermore, $\mathcal{E}_{n}(\lambda) \in \stackrel{n}{\mathcal{V}}_{r}{ }_{r}$ coincides with the standard higher order Euler-Lagrange morphism $\mathcal{E}(\lambda)$ associated with $\lambda$.

Making use of the above sheaf isomorphisms and of decomposition formulae (2) and (3), in [5] we proved that the Lie derivative operator with respect to the $r$ th order prolongation $j_{r} \Xi$ of a projectable vector field $(\Xi, \xi)$ can be conveniently represented on the quotient sheaves of the variational sequence in terms of an operator, the variational Lie derivative $\mathcal{L}_{j_{r} \Xi}$, as follows:
if $0 \leq p \leq n-1$ and $\mu \in \stackrel{p}{\mathcal{V}}_{r}$, then

$$
\begin{equation*}
\left.\left.\left.\mathcal{L}_{j_{r} \Xi} \mu=\xi\right\lrcorner d_{H} \mu+j_{r+2} \Xi_{V}\right\lrcorner d_{V} \mu+d_{H}(\xi\lrcorner \mu\right) ; \tag{4}
\end{equation*}
$$

if $p=n$ and $\lambda \in \stackrel{\mathcal{V}}{r}^{r}$, then

$$
\begin{equation*}
\left.\left.\left.\mathcal{L}_{j_{r} \Xi} \lambda=\Xi_{V}\right\lrcorner \mathcal{E}(\lambda)+d_{H}\left(j_{r} \Xi_{V}\right\lrcorner p_{d_{V} \lambda}+\xi\right\lrcorner \lambda\right) ; \tag{5}
\end{equation*}
$$

if $p=n+1$ and $\eta \in \stackrel{n+1}{\mathcal{V}}{ }_{r}$, then

$$
\begin{equation*}
\left.\mathcal{L}_{j_{r} \Xi \eta}=\mathcal{E}\left(j_{r} \Xi_{V}\right\lrcorner \eta\right)+\tilde{H}_{d \eta}\left(j_{2 r+1} \Xi_{V}\right) \tag{6}
\end{equation*}
$$

## 3. Variations

We shall here introduce variations of a morphism as multiparameter deformations showing that this is equivalent to take iterated variational Lie derivatives with respect to vertical vector fields.

We define the $i$-th variation of a section and introduce the $i$-th variation of a morphism of the kind $\alpha: J_{r} \boldsymbol{Y} \rightarrow \stackrel{k}{\wedge} T^{*} J_{r} \boldsymbol{Y}$ along the section $\sigma$ in terms of the pull-back by means of the $r$-th prolongation of the $i$-th variation of a section $\sigma: \boldsymbol{X} \rightarrow \boldsymbol{Y}$.

Definition 3.1. Let $\sigma: \boldsymbol{X} \rightarrow \boldsymbol{Y}$ be a section and $i$ any integer. An $i$-th variation of $\sigma$ is a smooth section $\Gamma_{i}: \boldsymbol{I} \times \boldsymbol{X} \rightarrow \boldsymbol{Y}, \mathbf{0} \in \boldsymbol{I} \subset \mathbf{R}^{\mathbf{i}}$, such that $\Gamma_{i}(\mathbf{0})=\sigma$.

In other words, $\Gamma_{i}$ is a $i$-parameter smooth deformation of $\sigma$.
Let $\Xi_{1}, \ldots, \Xi_{i}$ be vertical vector fields on $\boldsymbol{Y}$ and let $\Gamma_{i}\left(t_{1}, \ldots, t_{i}\right)$ be an $i$-th variation of the section $\sigma$ such that

$$
\begin{aligned}
& \left.\frac{\partial \Gamma_{i}}{\partial t_{1}}\left(t_{1}, 0, \ldots, 0\right)\right|_{t_{1}=0}=\Xi_{1} \circ \sigma \\
& \left.\frac{\partial \Gamma_{i}}{\partial t_{2}}\left(t_{1}, t_{2}, 0, \ldots, 0\right)\right|_{t_{2}=0}=\Xi_{2} \circ \Gamma_{i}\left(t_{1}, 0, \ldots, 0\right) \\
& \cdots, \\
& \left.\frac{\partial \Gamma_{i}}{\partial t_{i}}\left(t_{1}, t_{2}, \ldots, t_{i-1}, t_{i}\right)\right|_{t_{i}=0}=\Xi_{i} \circ \Gamma_{i}\left(t_{1}, t_{2}, \ldots, t_{i-1}, 0\right)
\end{aligned}
$$

In this case we say that $\Gamma_{i}$ is generated by the $i$-tuple $\left(\Xi_{1}, \ldots, \Xi_{i}\right)$.
We have thus the following characterization of $\Gamma_{i}$ as the variation of $\sigma$ by means of vertical flows.

Proposition 3.2. Let $\psi_{t_{k}}^{k}$, with $1 \leq k \leq i$, be the flows generated by the vertical vector fields $\Xi_{k}$. Then for the $\Gamma_{i}$ generated by $\left(\Xi_{1}, \ldots, \Xi_{i}\right)$ we have:

$$
\begin{equation*}
\Gamma_{i}\left(t_{1}, \ldots, t_{i}\right)=\psi_{t_{i}}^{i} \circ \ldots \circ \psi_{t_{1}}^{1} \circ \sigma \tag{7}
\end{equation*}
$$

Definition 3.3. Vertical vector fields on any fiber bundle which 'deform' sections as above are called variation vector fields.
3.1. Variations of morphisms and Lie derivative. We shall define the variation of a generic fibered morphism of the kind $\alpha: J_{r} \boldsymbol{Y} \rightarrow \stackrel{k}{\wedge} T^{*} J_{r} \boldsymbol{Y}$ along the section $\sigma$.

Definition 3.4. Let $\alpha: J_{r} \boldsymbol{Y} \rightarrow \stackrel{k}{\wedge} T^{*} J_{r} \boldsymbol{Y}$ and let $\Gamma_{i}$ be an $i$-th variation of the section $\sigma$. We define the $i$-th variation of the morphism $\alpha$ to be

$$
\begin{equation*}
\delta^{i} \alpha:=\left.\frac{\partial^{i}}{\partial t_{1} \ldots \partial t_{i}}\right|_{t_{1}, \ldots, t_{i}=0}\left(\alpha \circ j_{r} \Gamma_{i}\left(t_{1}, \ldots, t_{i}\right)\right) . \tag{8}
\end{equation*}
$$

If just one variation $\Gamma_{1}$, generated by a single vector field $\Xi_{1}$, is involved we write for simplicity $\delta$ instead of $\delta^{1}$. The following Lemma states the relation between the $i-$ th variation of a morphism and its iterated Lie derivative.

Lemma 3.5. Let $\alpha: J_{r} \boldsymbol{Y} \rightarrow \stackrel{k}{\wedge} T^{*} J_{r} \boldsymbol{Y}$ and $L_{j_{r} \Xi_{k}}$ be the Lie derivative operator with respect to $j_{r} \Xi_{k}$.

Let $\Gamma_{i}$ be the $i$-th variation of the section $\sigma$ by means of the variation vector fields $\Xi_{1}, \ldots, \Xi_{i}$ on $\boldsymbol{Y}$. Then we have

$$
\begin{equation*}
\delta^{i} \alpha=\left(j_{r} \sigma\right)^{*} L_{j_{r} \Xi_{1}} \ldots L_{j_{r} \Xi_{i}} \alpha . \tag{9}
\end{equation*}
$$

Proof. By applying Definition 3.4 and Proposition 3.2 we have

$$
\begin{aligned}
& \delta^{i} \alpha=\left.\frac{\partial^{i}}{\partial t_{1} \ldots \partial t_{i}}\right|_{t_{1} \ldots t_{i}=0}\left[\left(j_{r}\left(\psi_{t_{i}}^{i} \circ \ldots \circ \psi_{t_{1}}^{1} \circ \sigma\right)\right)^{*} \alpha\right]= \\
& \left.\frac{\partial^{i}}{\partial t_{1} \ldots \partial t_{i}}\right|_{t_{1} \ldots t_{i}=0}\left[\left(j_{r} \sigma\right)^{*} \circ\left(j_{r} \psi_{t_{1}}^{1}\right)^{*} \circ \ldots \circ\left(j_{r} \psi_{t_{i}}^{i}\right)^{*} \alpha\right]= \\
& \left.\left(j_{r} \sigma\right)^{*} \frac{\partial^{i}}{\partial t_{1} \ldots \partial t_{i}}\right|_{t_{1} \ldots t_{i}=0}\left[\left(j_{r} \psi_{t_{1}}^{1}\right)^{*} \circ \ldots \circ\left(j_{r} \psi_{t_{i}}^{i}\right)^{*}\right] \alpha= \\
& \left(j_{r} \sigma\right)^{*}\left[\left.\left.\frac{\partial}{\partial t_{1}}\left(j_{r} \psi_{t_{1}}^{1}\right)^{*}\right|_{t_{1}=0} \circ \ldots \circ \frac{\partial}{\partial t_{i}}\left(j_{r} \psi_{t_{i}}^{i}\right)^{*}\right|_{t_{i}=0}\right] \alpha= \\
& \left(j_{r} \sigma\right)^{*} L_{j_{r} \Xi_{1} \ldots L_{j_{r} \Xi_{i}} \alpha,}
\end{aligned}
$$

where $\psi_{t_{k}}^{k}$ are the vertical flows respectively generated by $\Xi_{k}$, for $k=1, \ldots, i$.

From now on, we will restrict to the case of the second order variation of a morphism $\lambda \in \stackrel{n}{\Lambda} \Lambda_{r}$. Higher order variations will be deeply investigated elsewhere.

In the following Lemma we show that the first order variation $\delta^{1} \lambda:=\delta \lambda$ of $\lambda$ is simply related to the vertical differential of $\lambda$.

Lemma 3.6. Let $\sigma$ be a section of $\boldsymbol{Y}, \boldsymbol{\Xi}$ a variation vector field on $\boldsymbol{Y}$ and $\lambda \in \stackrel{n}{\Lambda}_{r}$. Then we have

$$
\begin{equation*}
\left.\delta \lambda=j_{r} \sigma^{*}\left(j_{r} \Xi\right\lrcorner d_{V} \lambda\right) \tag{10}
\end{equation*}
$$

Proof. In fact we have

$$
\begin{aligned}
& \delta \lambda=j_{r} \sigma^{*} \mathcal{L}_{\left.\left.j_{r} \Xi \lambda=j_{r} \sigma^{*}(\Xi\lrcorner \mathcal{E}(\lambda)+d_{H}\left(j_{r} \Xi\right\lrcorner p_{d_{V} \lambda}\right)\right)=}^{\left.\left.\left.j_{r} \sigma^{*}(\Xi\lrcorner \mathcal{E}(\lambda)+\left(j_{r} \Xi\right\lrcorner d_{H} p_{d_{V} \lambda}\right)\right)=j_{r} \sigma^{*}\left(j_{r} \Xi\right\lrcorner d_{V} \lambda\right),}
\end{aligned}
$$

since $\left.\Xi\lrcorner d_{H}=d_{H} \Xi\right\lrcorner$, for any vertical vector field $\Xi$ on $\boldsymbol{Y}$.
For notational convenience, in the sequel we shall denote with a superimposed bar all objects defined on a vertical prolongation $V \boldsymbol{Y}$ and with two bars those defined on the iterated vertical prolongation $V(V \boldsymbol{Y})$; e.g. $\bar{\Xi}$ will denote a variational vector field on $V \boldsymbol{Y}$ and $\bar{\delta}$ the variation operator on $V \boldsymbol{Y}$.

By iterating this result we can characterize the second order variation of $\lambda$ as follows.

Proposition 3.7. Let $\lambda \in\left(\stackrel{n}{\Lambda}_{r}\right)_{\boldsymbol{Y}}$ and $\sigma$ be a section of $\boldsymbol{Y} \rightarrow \boldsymbol{X}$. Let $\Xi_{1}, \Xi_{2}$ be two variation vector fields on $\boldsymbol{Y}$ generating a second order variation $\Gamma_{2}$ of $\sigma$, and let $\bar{\Xi}_{2}$ be a variation vector field on $V \boldsymbol{Y}$, which projects down onto $\Xi_{2}$ and generates a first order variation $\bar{\Gamma}$ of $\sigma$. Moreover, let $\bar{d}, \overline{\mathcal{E}}$ and $\bar{p}$ be the exterior
differential, the Euler-Lagrange morphism and the momentum morphism on $V \boldsymbol{Y}$, respectively. Then we have

$$
\begin{align*}
\delta^{2} \lambda & \left.\left.=\bar{\Xi}_{2}\right\lrcorner \overline{\mathcal{E}}(\delta \lambda)+\bar{d}_{H}\left(j_{r} \bar{\Xi}_{2}\right\lrcorner \bar{p}_{\bar{d}_{V} \delta \lambda}\right)  \tag{11}\\
& \left.\left.=\bar{\delta}\left(\Xi_{1}\right\lrcorner \mathcal{E}(\lambda)+d_{H}\left(j_{r} \Xi_{1}\right\lrcorner p_{d_{V} \lambda}\right)\right) . \tag{12}
\end{align*}
$$

Proof. We apply the above Lemma and the fact that $d_{H} \delta=\delta d_{H}$, which follows directly from the analogous naturality property of the Lie derivative operator.

Remark 3.8. Owing to the linearity properties of $d_{V} \lambda$, by an abuse of notation, we can think of the operator $\delta$ as a linear morphism with respect to the vector bundle structure $J_{r} V \boldsymbol{Y} \rightarrow \boldsymbol{Y}$ of the kind

$$
\begin{equation*}
\delta: J_{r} \boldsymbol{Y} \underset{\boldsymbol{Y}}{\times} J_{r} V \boldsymbol{Y} \rightarrow \wedge^{n} T^{*} \boldsymbol{X} . \tag{13}
\end{equation*}
$$

This property can be obviously iterated for each integer $i$, so that one can analogously define an $i$-linear morphism $\delta^{i}$.

## 4. Variational vertical derivatives

In this Section we restrict our attention to morphisms which are sections of sheaves in the variational sequence. We define the $i$-th order variational vertical derivative of morphisms, by showing that the $i$-th variation operator passes to the quotient in the variational sequence.

Lemma 4.1. Let $\alpha \in\left(\stackrel{n}{V}_{r}\right)_{\boldsymbol{Y}}$. We have

$$
\left[\delta^{i} \alpha\right]=\hat{\delta}^{i}[\alpha]
$$

where we set $\hat{\delta}^{i}:=\mathcal{L}_{\Xi_{i}} \ldots \mathcal{L}_{\Xi_{1}}$.
Proof. In fact, we have

$$
\left[\delta^{i} \alpha\right]=\left[L_{\Xi_{i}} \ldots L_{\Xi_{1}} \alpha\right]=\mathcal{L}_{\Xi_{i}} \ldots \mathcal{L}_{\Xi_{1}}[\alpha]:=\hat{\delta}^{i}[\alpha]
$$

since $\delta^{i}$, as well as $L_{\Xi_{i}}$ (see [5]), preserve the contact structure.

Definition 4.2. We call the operator $\hat{\delta}^{i}$ the $i$-th variational vertical derivative operator.

Theorem 4.3. The functor $\hat{\delta}$ is defined on the category of variational sequences and sends the sequence associated with the fibration $\boldsymbol{Y} \rightarrow \boldsymbol{X}$ into the sequence associated with the fibration $V \boldsymbol{Y} \rightarrow \boldsymbol{X}$.

Proof. It follows from Proposition 3.7, Equation (11), Remark 3.8 and Lemma 4.1.

The above can be summarized in the following diagram.


This enables us to represent variations of morphisms in the variational sequence. Here we shall investigate in detail the case of the second order variation of morphisms in the variational sequence. We show that classical results concerning the second variation can be restated in a very simple way in terms of variational Lie derivatives and global decomposition formulae.

To this aim, let us consider how the representations (5) and (6) of the variational Lie derivative specialize in the case of variation vector fields.

Lemma 4.4. Let $\Xi_{1}, \Xi_{2}$ be two variation vector fields on $\boldsymbol{Y}$ and $\bar{\Xi}_{2}$ be a variation vector field on $V \boldsymbol{Y}$, which projects down onto $\Xi_{2}$. Let $\overline{\mathcal{E}}$ and $\bar{p}$ be, respectively, the Euler-Lagrange morphism and the momentum morphism associated with the Lagrangian $\hat{\delta} \lambda$. Then we have

$$
\begin{gather*}
\left.\left.\mathcal{L}_{j_{r} \Xi_{1}} \lambda=\Xi_{1}\right\lrcorner \mathcal{E}(\lambda)+d_{H}\left(j_{r} \Xi_{1}\right\lrcorner p_{d_{V} \lambda}\right) ;  \tag{14}\\
\left.\left.\left.\mathcal{L}_{j_{r} \Xi_{2}} \mathcal{L}_{j_{r} \Xi_{1}} \lambda=\bar{\Xi}_{2}\right\lrcorner \overline{\mathcal{E}}(\hat{\delta} \lambda)+\bar{d}_{H}\left[j_{r} \bar{\Xi}_{2}\right\lrcorner \bar{p}_{\bar{d}_{V}(\hat{\delta} \lambda)}-\hat{\delta}\left(j_{r} \Xi_{1}\right\lrcorner p_{d_{V} \lambda}\right)\right] . \tag{15}
\end{gather*}
$$

Proof. It follows from Equation (5) and Lemma 4.1.
The above Lemma together with Proposition 3.7 and Equation (11) gives us the following characterization of the second order variation of a generalized Lagrangian in the variational sequence.

Proposition 4.5. Let $\lambda \in\left(\stackrel{n}{\mathcal{V}}_{r}\right)_{\boldsymbol{Y}}, \hat{\delta} \lambda \in\left(\stackrel{n}{\mathcal{V}}_{r}\right)_{V \boldsymbol{Y}}$. We have

$$
\begin{equation*}
\left.\hat{\delta}^{2} \lambda=\overline{\mathcal{E}}\left(\bar{\Xi}_{2}\right\lrcorner \hat{\delta} \lambda\right)+\tilde{H}_{d \hat{\delta} \lambda}\left(\bar{\Xi}_{2}\right), \tag{16}
\end{equation*}
$$

where $\tilde{H}_{\text {d } \hat{\delta} \lambda}$ is the unique morphism belonging to $\stackrel{1}{\mathcal{C}}_{(2 r, r)} \otimes \stackrel{1}{\mathcal{C}}_{(2 r, 0)} \wedge \stackrel{n}{\mathcal{H}}_{2 r}$ such that, for all $\Xi_{1}: \boldsymbol{Y} \rightarrow V \boldsymbol{Y}, E_{\left.j_{r} \Xi\right\rfloor d \hat{\delta} \lambda}=C_{1}^{1}\left(j_{2 r} \Xi_{1} \otimes \tilde{H}_{d \hat{\delta} \lambda}\right)$, and $C_{1}^{1}$ stands for tensor contraction.

Proof. In fact we have

$$
\begin{equation*}
\hat{\delta}^{2} \lambda=\mathcal{L}_{j_{r} \bar{\Xi}_{2}} \hat{\delta} \lambda, \tag{17}
\end{equation*}
$$

so that the assertion follows from a straightforward application of the representation provided by Equation (6).

This result provides an equivalent intrinsic interpretation of the second variation of the action functional evaluated in [4], equation (2.14), by just setting $\bar{\Xi}_{2}=\rho$ and $\Xi_{1}=\eta$.

The following is an application of an abstract result due to Kolář [10], concerning a global decomposition formula for vertical morphisms.

Lemma 4.6. Let $\hat{\mu}: J_{s} \boldsymbol{Y} \rightarrow \stackrel{*}{\mathcal{C}}_{k}[V \boldsymbol{Y}] \wedge \stackrel{p}{\wedge} T^{*} \boldsymbol{X}$, with $0 \leq p \leq n$ and let $\bar{d}_{H} \hat{\mu}=0$. Then we have $\hat{\mu}=E_{\hat{\mu}}+F_{\hat{\mu}}$, where

$$
\begin{equation*}
E_{\hat{\mu}}: J_{2 s+k} V \boldsymbol{Y} \rightarrow \stackrel{*}{\mathcal{C}}_{0}[V \boldsymbol{Y}] \wedge \stackrel{p}{\wedge} T^{*} \boldsymbol{X}, \tag{18}
\end{equation*}
$$

and locally, $F_{\hat{\mu}}=\bar{d}_{H} M_{\hat{\mu}}$, with $M_{\hat{\mu}}: J_{2 s+k-1} V \boldsymbol{Y} \rightarrow \stackrel{*}{\mathcal{C}}_{k-1}[V \boldsymbol{Y}] \wedge \wedge^{p-1} \Lambda^{*} \boldsymbol{X}$.
Proof. Following e.g. [9, 10, 21], the global morphisms $E_{\hat{\mu}}$ and $D_{H} M_{\hat{\mu}}$ can be evaluated by means of a backwards procedure. Hereafter the canonical isomorphism $J_{r} V \boldsymbol{Y} \simeq V J_{r} \boldsymbol{Y}$ is obviously understood.

Now, it is very easy to see from Remark 3.8 and by means of a simple calculation that the following holds true.

Lemma 4.7. Let $\chi(\lambda):=\tilde{H}_{d \hat{\delta} \lambda}$. We have $\chi(\lambda): J_{2 r} \boldsymbol{Y} \rightarrow \stackrel{\mathcal{C}}{ }_{r}[V \boldsymbol{Y}] \wedge\left(n^{n} T^{*} \boldsymbol{X}\right)$ and $\bar{d}_{H} \chi(\lambda)=0$.

Thus, as a straightforward application of Lemma 4.6, we obtain our main result which consists in a suitable geometrical interpretation of the second variation of a generalized Lagrangian and provides a new characterization of the Jacobi morphism in the framework of variational sequences.

Theorem 4.8. Let $\chi(\lambda)$ be as in the above Lemma. Then we have

$$
\chi(\lambda)=E_{\chi(\lambda)}+F_{\chi(\lambda)}
$$

where

$$
\begin{equation*}
E_{\chi(\lambda)}: J_{4 r} \boldsymbol{Y} \rightarrow \stackrel{*}{\mathcal{C}}_{0}[V \boldsymbol{Y}] \wedge\left(\stackrel{n}{\wedge} T^{*} \boldsymbol{X}\right) \tag{19}
\end{equation*}
$$

and locally, $F_{\chi(\lambda)}=\bar{d}_{H} M_{\chi(\lambda)}$, with $M_{\chi(\lambda)}: J_{4 r-1} \boldsymbol{Y} \rightarrow \stackrel{*}{\mathcal{C}}_{r-1}[V \boldsymbol{Y}] \wedge \wedge^{n-1} \Lambda^{*} \boldsymbol{X}$.
Remark 4.9. If the coordinate expression of $\chi(\lambda)$ is given by

$$
\chi(\lambda)=\chi_{i}^{\boldsymbol{\alpha}} \boldsymbol{\vartheta}_{\boldsymbol{\alpha}}^{i} \wedge \omega,
$$

where $\boldsymbol{\vartheta}_{\boldsymbol{\alpha}}^{i}$ are contact forms on $J_{k} V \boldsymbol{Y}$, the corresponding coordinate expressions of $E_{\chi(\lambda)}$ and $M_{\chi(\lambda)}$ are respectively given by

$$
\begin{aligned}
E_{\chi(\lambda)} & =E_{i} \boldsymbol{\vartheta}^{i} \wedge \omega \\
M_{\chi(\lambda)} & =M_{i}^{\boldsymbol{\alpha}+\lambda} \boldsymbol{\vartheta}_{\boldsymbol{\alpha}}^{i} \wedge \omega_{\lambda} .
\end{aligned}
$$

We have, in particular

$$
E_{\chi(\lambda)}=(-1)^{|\boldsymbol{\beta}|} D_{\boldsymbol{\beta}} \chi_{i}^{\boldsymbol{\beta}} \vartheta^{i} \wedge \omega
$$

with $0 \leq|\boldsymbol{\beta}| \leq k$.
Definition 4.10. We call the morphism $\mathcal{J}(\lambda):=E_{\chi(\lambda)}$ the generalized Jacobi morphism associated with the Lagrangian $\lambda$.

Again, this result can be compared with equation (2.19) of [4].
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## References

[1] C. Caratheodory: Calculus of variations and partial differential equations of the first order, Chelsea Publ. Co., New York, 1982, p. 262.
[2] B. Casciaro, M. Francaviglia: Covariant second variation for first order Lagrangians on fibered manifolds. I: generalized Jacobi fields, Rend. Matem. Univ. Roma VII (16) (1996) 233-264.
[3] B. Casciaro, M. Francaviglia: A New Variational Characterization of Jacobi Fields along Geodesics, Ann. Mat. Pura e Appl. CLXXII (IV) (1997) 219-228.
[4] B. Casciaro, M. Francaviglia, V. Tapia: On the Variational Characterization of Generalized Jacobi Equations, Proc. Diff. Geom. and its Appl. (Brno, 1995); J. Janyška, I. Koláŕ, J. Slovák eds., Masaryk University (Brno, 1996) 353-372.
[5] M. Francaviglia, M. Palese, R. Vitolo: Symmetries in Finite Order Variational Sequences, to appear in Czech. Math. Journ..
[6] H. Goldschmidt, S. Sternberg: The Hamilton-Cartan Formalism in the Calculus of Variations, Ann. Inst. Fourier, Grenoble 23 (1) (1973) 203-267.
[7] I. KoláŘ: Lie Derivatives and Higher Order Lagrangians, Proc. Diff. Geom. and its Appl. (Nové Město na Moravě, 1980); O. Kowalski ed., Univerzita Karlova (Praha, 1981) 117-123.
[8] I. Koláǩ: On the second tangent bundle and generalized Lie derivatives, Tensor, N.S. 38 (1) (1982) 98-102.
[9] I. Kolář: A Geometrical Version of the Higher Order Hamilton Formalism in Fibred Manifolds, J. Geom. Phys. 1 (2) (1984) 127-137.
[10] I. Kolář: Some Geometrical Aspects of the Higher Order Variational Calculus, Geom. Meth. in Phys., Proc. Diff. Geom. and its Appl., (Nové Město na Moravě, 1983); D. Krupka ed., J. E. Purkyně University (Brno, 1984) 155-166
[11] D. Krupka: Some Geometric Aspects of Variational Problems in Fibred Manifolds, Folia Fac. Sci. Nat. UJEP Brunensis 14, J. E. Purkyně Univ. (Brno, 1973) 1-65.
[12] D. Krupka: Variational Sequences on Finite Order Jet Spaces, Proc. Diff. Geom. and its Appl. (Brno, 1989); J. Janyška, D. Krupka eds., World Scientific (Singapore, 1990) 236-254.
[13] D. Krupka: Topics in the Calculus of Variations: Finite Order Variational Sequences, Proc. Diff. Geom. and its Appl. (Opava, 1993) 473-495.
[14] D. Krupka, A. Trautman: General Invariance of Lagrangian Structures, Bull. Acad. Polon. Sci., Math. Astr. Phys. 22 (1974) (2) 207-211.
[15] L. Mangiarotti, M. Modugno: Fibered Spaces, Jet Spaces and Connections for Field Theories, in Proc. Int. Meet. on Geom. and Phys., Pitagora Editrice (Bologna, 1983) 135165.
[16] M. Palese: Geometric Foundations of the Calculus of Variations. Variational Sequences, Symmetries and Jacobi Morphisms. Ph.D. Thesis, University of Torino (2000).
[17] D.J. Saunders: The Geometry of Jet Bundles, Cambridge Univ. Press (Cambridge, 1989).
[18] H. Rund: The Hamilton-Jacobi theory in the calculus of variations, D. Van Nostrand Company LTD (London 1966) p. 127.
[19] A. Trautman: A Metaphysical Remark on Variational Principles, Acta Phys. Pol. B XX (1996) 1-9.
[20] R. Vitolo: On Different Geometric Formulations of Lagrangian Formalism, Diff. Geom. and its Appl. 10 (1999) 225-255.
[21] R. Vitolo: Finite Order Lagrangian Bicomplexes, Math. Proc. Cambridge Phil. Soc. 125 (1) (1999) 321-333.

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