

ONE REMARK ON VARIATIONAL PROPERTIES OF GEODESICS IN PSEUDORIEMANNIAN AND GENERALIZED FINSLER SPACES

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Abstract. A new variational property of geodesics in (pseudo-)Riemannian and Finsler spaces has been found.

1. Introduction

Let us assume an *n*-dimensional **Finsler space** F_n with local coordinates $x \equiv (x^1, \ldots, x^n)$ on the underlying manifold M_n , and a (positive definite) metric form with local expression

$$\mathrm{d}s^2 = g_{ij}(x, \dot{x}) \mathrm{d}x^i \mathrm{d}x^j \,. \tag{1}$$

Here $g_{ij}(x^1, \ldots, x^n, \dot{x}^1, \ldots, \dot{x}^n)$ are components of the metric tensor, and (x, \dot{x}) denote adapted local coordinates on the tangent bundle TM, i.e., $(\dot{x}^1, \ldots, \dot{x}^n)$ are coordinates of the "tangent vector" \dot{x} at x. Metric depends on "positions" and "velocities" in general.

In the Finsler space F_n there exists a (fundamental) function $F(x, \dot{x})$ which is homogeneous of the second degree in \dot{x}^i and satisfies

$$g_{ij}(x,\dot{x}) = rac{\partial^2 F(x,\dot{x})}{\partial \dot{x}^i \partial \dot{x}^j} \cdot$$

Particularly, the equality

$$F(x, \dot{x}) = g_{ij}(x, \dot{x}) \mathrm{d}x^i \mathrm{d}x^j$$

holds [3]. As it is well known, in the particular case when components of the metric tensor depend only on position coordinates (i.e., are independent of "velocity coordinates" \dot{x}) the Finsler space F_n turns out to be a **Riemannian space** V_n .

2. Pseudo-Riemannian and (Generalized) Finslerian Spaces

In what follows, the signature of the (non-degenerate) metric form is supposed to be arbitrary (we no more restrict ourselves onto positive definite metrics only) so that we can write

$$ds^2 = eg_{ij}(x, \dot{x}) dx^i dx^j, \qquad e = \pm 1$$
(2)

and the sign is determined in such a way that $ds^2 \ge 0$.

In short, we will call such metrics and spaces **Finslerian metrics** and **Finsler spaces** again, or *Riemannian*, respectively (more usually, they are called pseudo-Riemannian, or semi-Riemannian).

The arc length of a curve γ , given by parametrization $x^i = x^i(t)$, is given in a Finsler or Riemannian space (in our sense) by the integral

$$s = \int_{t_0}^{t_1} \sqrt{eg_{ij}(x(t), \dot{x}(t))\dot{x}^i(t)\dot{x}^j(t)} \,\mathrm{d}t, \qquad \dot{x}^i(t) = \frac{\mathrm{d}x^i(t)}{\mathrm{d}t} \,. \tag{3}$$

It is well known [3], that this integral is stationary in a Finsler space if and only if its extremals are **geodesic curves** determined by the equations

$$\ddot{x}^h + 2G^h(x, \dot{x}) = \varrho(t)\dot{x}^h \tag{4}$$

where $\rho(t)$ is a function, g^{ij} are components of the matrix inverse to (g_{ij}) , and

$$G^{h} = \frac{1}{2}g^{ij} \left(\frac{\partial^{2}F(x,\dot{x})}{\partial \dot{x}^{j}\partial x^{k}} \dot{x}^{k} - \frac{\partial F(x,\dot{x})}{\partial \dot{x}^{j}} \right)$$

are components of the Berwald connection. Let us emphasize that extremals of the integral of length are independent of reparametrization of geodesics. In Riemannian spaces, [2, 3], the components read

$$G^h = \frac{1}{2} \Gamma^h_{ij}(x) \, \dot{x}^i \dot{x}^j$$

where Γ_{ij}^h are the Christoffels of second type.

Many authors define a **geodesic** in V_n as an extremal curve of the integral

$$I = \int_{t_0}^{t_1} g_{ij}(x) \dot{x}^i \dot{x}^j \,\mathrm{d}t.$$
 (5)

Extremals of this variational problem are those geodesics which satisfy the equations (4) with $\rho(t) \equiv 0$.

Analogous situation is in Finsler spaces (in our generalized sense). Extremal curves of the integral (5) are determined together with their parameter, which is used to be called *canonical*. Note that particularly, arc length in V_n or F_n , respectively, is always canonical.

3. Generalized Variational Problem of Geodesics

In a Riemannian or in a Finsler space (in a more general sense explained above) consider the following more general **variational problem**

$$I = \int_{t_0}^{t_1} f(e \, g_{ij}(x, \dot{x}) \dot{x}^i \dot{x}^j) \,\mathrm{d}\tau$$
(6)

where e takes the values ± 1 , and $f(\tau)$ is a differentiable real-valued function (at least of class two) defined on some open domain $D \subset \mathbb{R}$ which is regular on D in the sense that $f'(\tau) \neq 0$ for all $\tau \in D$.

As an immediate consequence of the Euler-Lagrange equations for the Lagrange function $\mathcal{L} = f(e g_{ij} \dot{x}^i \dot{x}^j)$, it can be checked that the extremals satisfy the equations

$$\ddot{x}^h + 2G^h(x, \dot{x}) = -\frac{\mathrm{d}}{\mathrm{d}t} (\ln |f'(eg_{\alpha\beta} \dot{x}^\alpha \dot{x}^\beta)|) \dot{x}^h.$$
(7)

We can prove the following theorem.

Theorem 1. In (generalized) Finsler or Riemannian spaces, respectively, geodesic lines parameterized by a canonical parameter, which satisfy the condition

$$eg_{\alpha\beta}\dot{x}^{\alpha}\dot{x}^{\beta} = k \in D$$

are extremals of the integral (6).

Theorem 2. Consider (all) extremals of the integral (6) in a Finsler space (or in a Riemannian space, respectively). All curves arising under all possible regular reparameterizations of extremal curves belong to extremals, too, if and only if the function f takes the form $f(x) \equiv \alpha \sqrt{x}$ where α is some non-zero constant.

Theorem 3. All possible extremals of the integral (6) are just those geodesics which figure in Theorem 1 and Theorem 2. More precisely, in the particular case $f(x) \equiv \alpha \sqrt{x}, 0 \neq \alpha = \text{const}$, they are represented by all unparameterized geodesics (i.e., geodesics under all possible regular reparameterizations), while for all other functions f, satisfying the above assumptions of the problem (6), extremals are represented just by canonically parameterized geodesics only.

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